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# Improved Sobolev Embedding Theorems for Vector-valued Functions

*To the memory of our friend Rentaro AGEMI*

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## Abstract

The aim of this paper is to give an extension of the improved Sobolev embedding theorem for single-valued functions to the case of vector-valued functions which is involved with the three-dimensional massless Dirac operator together with the three- or two-dimensional Weyl–Dirac (or Pauli) operator, the Cauchy–Riemann operator and also the four-dimensional Euclidian Dirac operator.

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## 1 Introduction and Results

The improved Sobolev embedding theorem is the following inequality: For  $1 \leq p < q < \infty$ , there exists a positive constant  $C$  only depending on  $p$  and  $q$  (and  $n$ ) such that

$$\|\psi\|_q \leq C \|\nabla \psi\|_p^{p/q} \|\psi\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)} \quad (1.1)$$

for every  $\mathbb{C}$ -valued function  $\psi$  on  $\mathbb{R}^n$  which satisfies  $\nabla \psi \in L^p(\mathbb{R}^n)$  and belongs to the Banach space  $B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^n)$ , where  $\nabla = (\partial_1, \dots, \partial_n)$ ,  $\partial_j = \partial/\partial x_j$ ,  $i = 1, 2, \dots, n$ . Here with  $a < 0$ ,  $B_{\infty,\infty}^a(\mathbb{R}^n)$  stands for the homogeneous Besov space of indices  $(a, \infty, \infty)$  with norm

$$\|\psi\|_{B_{\infty,\infty}^a} := \sup_{t>0} t^{-a/2} \|e^{t\Delta} \psi\|_{\infty} \quad (1.2)$$

(e.g. [T, Sect.2.5.2, pp.190–192]). Here  $e^{t\Delta}$  stands for the heat semigroup acting on the  $\mathbb{C}$ -valued functions  $\psi$  on  $\mathbb{R}^n$ , where  $\Delta$  is the Laplacian in  $\mathbb{R}^n$ , and  $\|e^{t\Delta} \psi\|_{\infty} := \sup_x |(e^{t\Delta} \psi)(x)|$ . This was shown by Cohen et al. [CDPX] (cf. Cohen et al. [CMO]) and Ledoux [Le]. In fact, (1.1) is a very general inequality which covers not only

the classical Sobolev inequalities  $\|\psi\|_q \leq C\|\nabla\psi\|_p$  with  $\frac{1}{q} = \frac{1}{p} - \frac{1}{n}$ ,  $1 \leq p < n$ , for every function  $\psi$  vanishing at infinity in a certain mild sense, but also the Gagliardo–Nirenberg inequalities

$$\|\psi\|_q \leq C\|\nabla\psi\|_p^{p/q}\|\psi\|_r^{1-(p/q)}, \quad \frac{1}{q} = \frac{1}{p} - \frac{r}{qn}. \quad (1.3)$$

In all the inequalities the functions  $\psi$  are supposed to be *single-valued* functions. In this work we will show an inequality like (1.1) for the case where the  $\psi$  are *vector-valued* functions. Of course, inequality (1.1) holds also if one replaces single-valued functions  $\psi$  by vector-valued functions  $f$ , understanding their semi-norm  $\|\nabla f\|_p$  on the right-hand side of (1.1) in the sense of (1.11) as below. But what we want to have is an inequality in the situation where the semi-norm concerned with the first-order derivatives is related to the massless Dirac operator

$$\alpha \cdot p = \alpha_1 p_1 + \alpha_2 p_2 + \alpha_3 p_3 = \alpha \cdot (-i\nabla) = -i(\alpha_1 \partial_1 + \alpha_2 \partial_2 + \alpha_3 \partial_3), \quad (1.4)$$

therefore, acting on  $\mathbb{C}^4$ -valued functions  $f(x) = {}^t(f_1(x), f_2(x), f_3(x), f_4(x))$  defined in special 3-dimensional space  $\mathbb{R}^3$ , though not in general  $\mathbb{R}^n$ . In (1.4),  $\alpha := (\alpha_1, \alpha_2, \alpha_3)$  is the triple of the  $4 \times 4$  Dirac matrices which satisfy the anti-commutation relation  $\alpha\alpha_k + \alpha_k\alpha_j = 2\delta_{jk}I_4$ ,  $j, k = 1, 2, 3$ , where  $I_4$  is the  $4 \times 4$ -identity matrix. We are concerned mainly with what are usually called “Dirac matrices”:

$$\alpha_j = \begin{pmatrix} 0_2 & \sigma_j \\ \sigma_j & 0_2 \end{pmatrix} \quad (j = 1, 2, 3) \quad (1.5)$$

with the  $2 \times 2$  zero matrix  $0_2$  and the triple of  $2 \times 2$  Pauli matrices

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (1.6)$$

In the beginning let us confirm the notations to be used about norms for vectors and functions. First of all, the absolute value of a number  $c := a + ib \in \mathbb{C}$  is denoted, as usual, by  $|c| := \sqrt{a^2 + b^2}$ . Next, we shall use the standard notations of the  $\ell^p$  and  $\ell^\infty$  norm for an  $m$ -vector  $a = {}^t(a_1, a_2, \dots, a_m) \in \mathbb{C}^m$ :

$$\begin{aligned} |a|_{\ell^p} &:= \left( \sum_{k=1}^m |a_k|^p \right)^{1/p} = (|a_1|^p + |a_2|^p + \dots + |a_m|^p)^{1/p}, \quad 1 \leq p < \infty, \\ |a|_{\ell^\infty} &:= \bigvee_{k=1}^m |a_k| = |a_1| \vee |a_2| \vee \dots \vee |a_m|, \end{aligned} \quad (1.7)$$

where  $b_1 \vee b_2 \vee \dots \vee b_m$  denotes  $\max\{b_1, b_2, \dots, b_m\}$ . The  $L^p$  and  $L^\infty$  norms for a  $\mathbb{C}^m$ -valued function  $f(x) = {}^t(f_1(x), f_2(x), \dots, f_m(x))$  are given, respectively, by

$$\|f\|_p = \left( \int |f(x)|_{\ell^p}^p dx \right)^{1/p}, \quad 1 \leq p < \infty, \quad (1.8)$$

In [IS] we considered the case  $m = 4$  and introduced the semi-norm

$$\begin{aligned} \|(\alpha \cdot p)f\|_p &= \left( \int |(\alpha \cdot p)f(x)|_{\ell^p}^p dx \right)^{1/p}, \quad 1 \leq p < \infty, \\ |(\alpha \cdot p)f(x)|_{\ell^p}^p &= \left| \sum_{j=1}^3 \alpha_j p_j f(x) \right|_{\ell^p}^p = \sum_{k=1}^4 \left| \left( \sum_{j=1}^3 \alpha_j p_j f \right)_k(x) \right|^p = \sum_{k=1}^4 \left| \left( \sum_{j=1}^3 \alpha_j \partial_j f \right)_k(x) \right|^p. \end{aligned} \quad (1.9)$$

for  $f(x) = {}^t(f_1(x), f_2(x), f_3(x), f_4(x))$  defined on  $\mathbb{R}^3$ . The Banach spaces obtained as completion in the norm  $\|f\|_{\alpha \cdot p, 1, p} := (\|f\|_p^p + \|(\alpha \cdot p)f\|_p^p)^{1/p}$  of the linear space  $C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$  and the linear space  $\{f \in C^\infty(\mathbb{R}^3; \mathbb{C}^4); f, (\alpha \cdot p)f \in L^p(\mathbb{R}^3; \mathbb{C}^4), j = 1, 2, 3\}$  were denoted in [IS] by  $\mathbb{H}_0^{1,p}(\mathbb{R}^3)$  and  $\mathbb{H}^{1,p}(\mathbb{R}^3)$ , respectively. However, in the present paper we denote them by  $H_{\alpha \cdot p, 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4)$  and  $H_{\alpha \cdot p}^{1,p}(\mathbb{R}^3; \mathbb{C}^4)$ , respectively.

Note that

$$\|(\alpha \cdot p)f\|_p \leq 3^{1-(1/p)} \|\nabla f\|_p, \quad (1.10)$$

where

$$\|\nabla f\|_p \equiv \left( \int |\nabla f(x)|_{\ell^p}^p dx \right)^{1/p}, \quad |\nabla f(x)|_{\ell^p}^p := \sum_{j=1}^3 |\partial_j f(x)|_{\ell^p}^p = \sum_{j=1}^3 \sum_{k=1}^4 |\partial_j f_k(x)|^p. \quad (1.11)$$

A proof of (1.10) only uses that  $\|\sum_{j=1}^m \psi_j\|_p \leq m^{1-(1/p)} (\sum_{j=1}^m \|\psi_j\|_p^p)^{1/p}$  for single-valued functions  $\psi_j$ ,  $j = 1, 2, \dots, m$ , an inequality following from Hölder's inequality.

As is the case for the Sobolev spaces of single-valued functions, so does coincidence hold for our Dirac–Sobolev spaces of vector-valued functions:  $H_{\alpha \cdot p, 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4) = H_{\alpha \cdot p}^{1,p}(\mathbb{R}^3; \mathbb{C}^4) = W_{\alpha \cdot p}^{1,p}(\mathbb{R}^3; \mathbb{C}^4)$ , where the last space is the Banach space of all  $f \in L^p(\mathbb{R}^3; \mathbb{C}^4)$  such that  $(\alpha \cdot p)f$  belongs to  $L^p(\mathbb{R}^3; \mathbb{C}^4)$ . It is shown in [IS] that, for  $1 < p < \infty$ ,  $H_{\alpha \cdot p, 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4)$  coincides with  $H_0^{1,p}(\mathbb{R}^3; \mathbb{C}^4)$ , the completion of  $C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$  in the norm  $\|f\|_{1,p} := (\|f\|_p^p + \|\nabla f\|_p^p)^{1/p}$ , while for  $p = 1$  the latter is a proper subspace of the former.

With  $a < 0$ , let  $B_{\infty, \infty}^a(\mathbb{R}^n; \mathbb{C}^4)$  be the homogeneous Besov space for  $\mathbb{C}^4$ -valued functions  $f(x)$  on  $\mathbb{R}^n$  of indices  $(a, \infty, \infty)$  with norm

$$\|f\|_{B_{\infty, \infty}^a} := \sup_{t>0} t^{-a/2} \|P_t f\|_\infty. \quad (1.12)$$

Here  $P_t := e^{t\Delta} I_4 = e^{t\Delta} I_4$  ( $I_4 : 4 \times 4$ -identity matrix) stands for the heat semigroup acting on the  $\mathbb{C}^4$ -valued functions  $f$  on  $\mathbb{R}^n$ , where  $\Delta$  is the Laplacian in  $\mathbb{R}^n$ ,  $e^{t\Delta}$  being the heat semigroup acting on the  $\mathbb{C}$ -valued functions on  $\mathbb{R}^n$ , and  $\|P_t f\|_\infty := \sup_x |P_t f(x)|_{\ell^\infty} = \sup_x \vee_{k=1}^4 |e^{t\Delta} f_k(x)|$ .

With the notations above concerning vector-valued functions, it is easy to see the following *trivial* version of (1.1) for  $\mathbb{C}^4$ -valued functions  $f$  holding : For  $1 \leq p < q < \infty$ , there exists a positive constant  $C$  such that

$$\|f\|_q \leq C \|\nabla f\|_p^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)} \quad (1.13)$$

for every  $\mathbb{C}^4$ -valued function  $f \in B_{\infty, \infty}^a(\mathbb{R}^n; \mathbb{C}^4)$  which satisfies  $\|\nabla f\|_p < \infty$ , therefore, in particular, for every  $f$  in the Sobolev space  $H_0^{1,p}(\mathbb{R}^n; \mathbb{C}^n) = H(\mathbb{R}^n; \mathbb{C}^4) = W^{1,p}(\mathbb{R}^n; \mathbb{C}^n)$  as well as in  $B_{\infty, \infty}^a(\mathbb{R}^n; \mathbb{C}^n)$ .

Then the first attempt to get a version of (1.1) for vector-valued functions in our sense was done in the paper [BES] where the authors showed, replacing the  $L^q$  norm of  $f$  on its left-hand side by the weak  $L^q$  norm of  $f$ , the following inequality, which they called *Dirac–Sobolev inequality* : For  $1 \leq p < q < \infty$ , there exists a constant  $C > 0$  such that

$$\|f\|_{q, \infty} \leq C \|(\alpha \cdot p)f\|_p^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)} \quad (1.14)$$

for every  $f \in B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$  which satisfies  $(\alpha \cdot p)f \in L^p(\mathbb{R}^3; \mathbb{C}^4)$ , therefore, in particular, for every  $f \in H_{\alpha \cdot p, 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4) \cap B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$ . As a result, this  $f$  belongs to the weak  $L^q$  space with the weak  $L^q$  norm defined by

$$\|f\|_{q,\infty} := \left[ \sup_{u>0} u^q |\{ |f|_{\ell^\infty} \geq u \}| \right]^{1/q}, \quad (1.15)$$

where  $|\{ |f|_{\ell^\infty} \geq u \}| = \int \chi_{\{|f|_{\ell^\infty} \geq u\}}(x) dx$  is the measure of the set  $\{|f|_{\ell^\infty} \geq u\}$  on which  $u \leq |f(x)|_{\ell^\infty} := \bigvee_{k=1}^4 |f_k(x)|$ ,  $dx$  being the Lebesgue measure on  $\mathbb{R}^3$ , and  $\chi_E(x)$  stands for the characteristic function of a subset  $E$  of  $\mathbb{R}^3$ .

Now one may ask oneself whether or not, for any  $1 \leq p < q < \infty$ , inequality (1.14) can hold valid, if replacing the weak  $L^q$  norm of  $f$  on the left-hand side by its strong  $L^q$  one as in the vector-valued version (1.13) of the original (1.1) but equipping on the right-hand side with *either* the first-order-derivative semi-norm  $\|(\alpha \cdot p)f\|_p$  as in (1.14) *or* some other one related to the massless Dirac operator  $\alpha \cdot p$ . In particular, we ask whether or not there exists a positive constant  $C$  such that

$$\|f\|_q \leq C \|(\alpha \cdot p)f\|_p^{p/q} \|f\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)} \quad (1.16)$$

for every  $f \in B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$  which satisfies  $(\alpha \cdot p)f \in L^p(\mathbb{R}^3; \mathbb{C}^4)$ . However, this replacement does not work so well; indeed (1.16) cannot hold for  $p = 1$ , although it holds for  $1 < p < q < \infty$ . A counterexample for this is essentially found in Balinsky–Evans–Umeda [BEU], which we will refer to in Section 2 below. This suggests us that in order to get an inequality like (1.16) with the strong  $L^q$  norm of  $f$  kept on the left-hand side, we have to replace the semi-norm  $\|(\alpha \cdot p)f\|_p$  on the right-hand side by a somewhat stronger one. This leads us to introduce a third semi-norm  $M_{\alpha \cdot p; p}(f)$  concerned with  $L^p$ -norm of the first-order derivatives of functions  $f = {}^t(f_1, f_2, f_3, f_4)$  in the space  $C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$ . Noting that the massless Dirac operator (1.4) can be rewritten, based on the representations (1.5) of the Dirac matrices  $\alpha_j$ ,  $j = 1, 2, 3$ , as

$$\alpha \cdot p = \begin{pmatrix} 0 & 0 & p_3 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & -p_3 \\ p_3 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & -p_3 & 0 & 0 \end{pmatrix}, \quad (1.17)$$

decompose it into the sum of its two parts:

$$\begin{aligned} \alpha \cdot p &= (\alpha \cdot p)P_{13} + (\alpha \cdot p)P_{24} \\ &= \begin{pmatrix} 0 & 0 & p_3 & 0 \\ 0 & 0 & p_1 + i p_2 & 0 \\ p_3 & 0 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & 0 & -p_3 \\ 0 & p_1 - i p_2 & 0 & 0 \\ 0 & -p_3 & 0 & 0 \end{pmatrix} \end{aligned} \quad (1.18)$$

where  $P_{13} := \text{diag}(1, 0, 1, 0)$  and  $P_{24} := \text{diag}(0, 1, 0, 1)$  are two projection matrices acting on the space  $\mathbb{C}^4$  of four-vectors, which satisfies that  $P_{13} + P_{24} = I_4$ , and define

$$M_{\alpha \cdot p; p}(f) := \left[ \|(\alpha \cdot p)P_{13}f\|_p^p + \|(\alpha \cdot p)P_{24}f\|_p^p \right]^{1/p}. \quad (1.19)$$

*At first sight, this introduction of the semi-norm  $M_{\alpha \cdot p; p}(f)$  here may appear to be artificial but we shall see soon that the semi-norm turns out to be rather intrinsic.*

Let us see how this semi-norm  $M_{\alpha \cdot p; p}(f)$  in (1.19) is related to the other semi-norms,  $\|(\alpha \cdot p)f\|_p$  and  $\|\nabla f\|_p$ . We have from (1.17)

$$(\alpha \cdot p)f = \begin{pmatrix} p_3 f_3 + (p_1 - i p_2) f_4 \\ (p_1 + i p_2) f_3 - p_3 f_4 \\ p_3 f_1 + (p_1 - i p_2) f_2 \\ (p_1 + i p_2) f_1 - p_3 f_2 \end{pmatrix},$$

so that, recalling the definition (1.9) of the  $\ell^p$  norm, we have

$$\begin{aligned} \|(\alpha \cdot p)f\|_{\ell^p}^p &= |p_3 f_3 + (p_1 - i p_2) f_4|^p + |(p_1 + i p_2) f_3 - p_3 f_4|^p \\ &\quad + |p_3 f_1 + (p_1 - i p_2) f_2|^p + |(p_1 + i p_2) f_1 - p_3 f_2|^p \\ &= |(p_1 + i p_2) f_1 - p_3 f_2|^p + |(p_1 - i p_2) f_2 + p_3 f_1|^p \\ &\quad + |(p_1 + i p_2) f_3 - p_3 f_4|^p + |(p_1 - i p_2) f_4 + p_3 f_3|^p, \end{aligned}$$

where we have rearranged the four terms, when passing through the second equality. Hence

$$\begin{aligned} \|(\alpha \cdot p)f\|_p^p &= \|(\partial_1 + i \partial_2) f_1 - \partial_3 f_2\|_p^p + \|(\partial_1 - i \partial_2) f_2 + \partial_3 f_1\|_p^p \\ &\quad + \|(\partial_1 + i \partial_2) f_3 - \partial_3 f_4\|_p^p + \|(\partial_1 - i \partial_2) f_4 + \partial_3 f_3\|_p^p. \end{aligned} \quad (1.20)$$

Then one can calculate the right-hand side of (1.19) to get

$$\begin{aligned} M_{\alpha \cdot p; p}(f)^p &= [\| \partial_3 f_3 \|_p^p + \|(\partial_1 + i \partial_2) f_3\|_p^p] + [\| \partial_3 f_1 \|_p^p + \|(\partial_1 + i \partial_2) f_1\|_p^p] \\ &\quad + [\|(\partial_1 - i \partial_2) f_4\|_p^p + \| \partial_3 f_3 \|_p^p] + [\|(\partial_1 - i \partial_2) f_2\|_p^p + \| \partial_3 f_2 \|_p^p] \\ &= (\|(\partial_1 + i \partial_2) f_1\|_p^p + \| \partial_3 f_1 \|_p^p) + (\|(\partial_1 - i \partial_2) f_2\|_p^p + \| \partial_3 f_2 \|_p^p) \\ &\quad + (\|(\partial_1 + i \partial_2) f_3\|_p^p + \| \partial_3 f_3 \|_p^p) + (\|(\partial_1 - i \partial_2) f_4\|_p^p + \| \partial_3 f_4 \|_p^p). \end{aligned} \quad (1.21)$$

We can compare (1.20) and (1.21) and recall (1.10) to show with aid of Hölder's inequality that for  $1 \leq p < \infty$ ,

$$2^{-(1-(1/p))} \|(\alpha \cdot p)f\|_p \leq M_{\alpha \cdot p; p}(f) \leq 2^{1-(1/p)} \|\nabla f\|_p, \quad (1.22)$$

so that the semi-norm  $M_{\alpha \cdot p; p}(f)$  is an intermediate one in strength lying between the other two first-order-derivative semi-norms  $\|(\alpha \cdot p)f\|_p$  and  $\|\nabla f\|_p$ . We shall denote by  $H_{M_{\alpha \cdot p; p}, 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4)$  the Banach space obtained as completion in the norm  $\|f\|_{M_{\alpha \cdot p; p}, 1, p} := (\|f\|_p^p + M_{\alpha \cdot p; p}(f)^p)^{1/p}$  of the space  $C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$ . From (1.22) we see the following inclusion relation among the three Banach spaces:

$$H_0^{1,p}(\mathbb{R}^3; \mathbb{C}^4) \subseteq H_{M_{\alpha \cdot p}, 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4) \subseteq H_{(\alpha \cdot p), 0}^{1,p}(\mathbb{R}^3; \mathbb{C}^4). \quad (1.23)$$

Now we are going to see a significant character of the semi-norm  $M_{\alpha \cdot p; p}(f)$  introduced in (1.19), by considering the other decompositions of the Dirac operator  $\alpha \cdot p$  in (1.17) than the one (1.18). In fact, there are a few other decompositions:

$M_\alpha^{(1)}$

$$\begin{aligned} \alpha \cdot p &= (\alpha \cdot p)P_{14} + (\alpha \cdot p)P_{23} \\ &\equiv \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & 0 & -p_3 \\ p_3 & 0 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & p_3 & 0 \\ 0 & 0 & p_1 + i p_2 & 0 \\ 0 & p_1 - i p_2 & 0 & 0 \\ 0 & -p_3 & 0 & 0 \end{pmatrix}, \end{aligned} \quad (1.24)$$

where  $P_{14} := \text{diag}(1, 0, 0, 1)$  and  $P_{23} := \text{diag}(0, 1, 1, 0)$  are two projection matrices acting on the space  $\mathbb{C}^4$  of four-vectors, so that  $P_{14} + P_{23} = I_4$ . Note that both the operators  $(\alpha \cdot p)P_{14}$  and  $(\alpha \cdot p)P_{23}$  on the right are selfadjoint, i.e.  $((\alpha \cdot p)P_{14})^* = (\alpha \cdot p)P_{14}$ ,  $((\alpha \cdot p)P_{23})^* = (\alpha \cdot p)P_{23}$ .

$M_\alpha^{(2)}$

$$\begin{aligned} \alpha \cdot p &= \begin{pmatrix} 0_2 & \sigma_1 p_1 + \sigma_2 p_2 \\ \sigma_3 p_3 & 0_2 \end{pmatrix} + \begin{pmatrix} 0_2 & \sigma_3 p_3 \\ \sigma_1 p_1 + \sigma_2 p_2 & 0_2 \end{pmatrix} \\ &\equiv \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & 0 \\ p_3 & 0 & 0 & 0 \\ 0 & -p_3 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & p_3 & 0 \\ 0 & 0 & 0 & -p_3 \\ 0 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} \\ &=: (\alpha \cdot p)_1 + (\alpha \cdot p)_2, \end{aligned} \tag{1.25}$$

where note that  $(\alpha \cdot p)_2$  is the adjoint of  $(\alpha \cdot p)_1$  as operators, say, in  $L^2(\mathbf{R}^3; \mathbb{C}^4)$ , i.e.  $(\alpha \cdot p)_2 = (\alpha \cdot p)_1^*$ .

$M_\alpha^{(3)}$

$$\begin{aligned} \alpha \cdot p &= \begin{pmatrix} 0 & \sigma_1 p_1 + \sigma_2 p_2 \\ \sigma_1 p_1 + \sigma_2 p_2 & 0 \end{pmatrix} + \begin{pmatrix} 0 & \sigma_3 p_3 \\ \sigma_3 p_3 & 0 \end{pmatrix} \\ &\equiv \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & 0 \\ 0 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & p_3 & 0 \\ 0 & 0 & 0 & -p_3 \\ p_3 & 0 & 0 & 0 \\ 0 & -p_3 & 0 & 0 \end{pmatrix} \\ &=: (\alpha \cdot p)_3 + (\alpha \cdot p)_4, \end{aligned} \tag{1.26}$$

where note that both the operators  $(\alpha \cdot p)_3$  and  $(\alpha \cdot p)_4$  on the right are selfadjoint.

Then we can see in the following proposition that the semi-norm  $M_{\alpha \cdot p; p}(f)$  of  $f \in C_0^\infty(\mathbf{R}^3; \mathbb{C}^4)$  defined by (1.19), though with the rather artificial decomposition (1.18) dependent on the pair  $(P_{13}, P_{24})$  of projection matrices, turns out to be meaningful enough to have some universal character.

**Proposition 1.0** *The semi-norm  $M_{\alpha \cdot p; p}(f)$  in (1.19) coincides with the ones to be defined with the decompositions (1.24), (1.25) and (1.26):*

$$M_{\alpha \cdot p; p}^{(1)}(f) := [\|(\alpha \cdot p)P_{14}f\|_p^p + \|(\alpha \cdot p)P_{23}f\|_p^p]^{1/p}; \tag{1.27a}$$

$$M_{\alpha \cdot p; p}^{(2)}(f) := [\|(\alpha \cdot p)_1 f\|_p^p + \|(\alpha \cdot p)_2 f\|_p^p]^{1/p} = [\|(\alpha \cdot p)_1 f\|_p^p + \|(\alpha \cdot p)_1^* f\|_p^p]^{1/p}; \tag{1.27b}$$

$$M_{\alpha \cdot p; p}^{(3)}(f) := [\|(\alpha \cdot p)_3 f\|_p^p + \|(\alpha \cdot p)_4 f\|_p^p]^{1/p}. \tag{1.27c}$$

More generally, in fact, every decomposition of  $\alpha \cdot p$  into its two parts,  $\alpha \cdot p = (\alpha \cdot p)_5 + (\alpha \cdot p)_6$ , such that each row of both the matrices  $(\alpha \cdot p)_5$  and  $(\alpha \cdot p)_6$  contains only one nonzero entry, defines the semi-norm  $M_{\alpha \cdot p; p}(f)$  which has the expression (1.21).

*Proof.* In fact, direct calculation of the right-hand sides of (1.27a), (1.27b) and (1.27c) in view of (1.24), (1.25) and (1.26) yields nothing but a rearrangement of the

last member of the expression (1.21) of  $M_{\alpha \cdot p; p}(f)$ . The assertion for the more general case is evident.  $\square$

We note that (1.27a) says that our semi-norm  $M_{\alpha \cdot p; p}(f)$ , which is defined in (1.19) by using the pair  $(P_{13}, P_{24})$  of projection matrices can be defined by using another pair  $(P_{14}, P_{23})$ . However, among all three possible pairs of projection matrices,  $(P_{13}, P_{24})$ ,  $(P_{14}, P_{23})$ ,  $(P_{12}, P_{34})$ , whose sum becomes the identity matrix  $I_4$ , the decomposition  $\alpha \cdot p = (\alpha \cdot p)P_{12} + (\alpha \cdot p)P_{34}$  to be defined with the remaining last pair consisting of  $P_{12} = \text{diag}(1, 1, 0, 0)$  and  $P_{34} = \text{diag}(0, 0, 1, 1)$ , is not fit for our semi-norm  $M_{\alpha \cdot p; p}(f)$ , since this decomposition does not satisfy the condition for the more general case in Proposition 1.0. In Section 6, we shall come back to this decomposition to discuss the issue.

The main result of this work is the following theorem.

**Theorem 1.1.** (with 3-dimensional massless Dirac operator) (i) For  $1 \leq p < q < \infty$ , a  $\mathbb{C}^4$ -valued function  $f = {}^t(f_1, f_2, f_3, f_4)$  belongs to  $L^q(\mathbb{R}^3; \mathbb{C}^4)$ , if  $f$  belongs to  $B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$  and satisfies  $M_{\alpha \cdot p; p}(f) < \infty$ , and further, there exists a positive constant  $C$  such that

$$\|f\|_q \leq C M_{\alpha \cdot p; p}(f)^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}. \quad (1.28)$$

Therefore this holds, in particular, for every  $f \in H_{M_{\alpha \cdot p}, 0}^{1, p}(\mathbb{R}^3; \mathbb{C}^4) \cap B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$ .

(ii) For  $\infty > p > 1$ , the three semi-norms  $\|(\alpha \cdot p)f\|_p$ ,  $M_{\alpha \cdot p; p}(f)$  and  $\|\nabla f\|_p$  are equivalent, so that the corresponding three Banach spaces in (1.23) coincide with one another:

$$H_0^{1, p}(\mathbb{R}^3; \mathbb{C}^4) = H_{M_{\alpha \cdot p}, 0}^{1, p}(\mathbb{R}^3; \mathbb{C}^4) = H_{(\alpha \cdot p), 0}^{1, p}(\mathbb{R}^3; \mathbb{C}^4). \quad (1.29)$$

Therefore assertion (i) turns out: For  $1 < p < q < \infty$ , there exists a positive constant  $C$  such that

$$\|f\|_q \leq C \|(\alpha \cdot p)f\|_p^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.30)$$

for every  $f \in B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$  whose semi-norm  $\|(\alpha \cdot p)f\|_p$ ,  $M_{\alpha \cdot p; p}(f)$  or  $\|\nabla f\|_p$  is finite. Therefore this holds, in particular, for every  $f$  in the above space (1.29) which belongs to  $B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)$ . (1.30) is equivalent to the vector-valued version (1.13) of (1.1) with  $n = 3$ .

Similarly we can also show the following five results in related different situations.

First, replacing the Dirac operator  $\alpha \cdot p$  in Theorem 1.1 by the 3-dimensional Weyl-Dirac (or Pauli) operator

$$\sigma \cdot p := \sigma_1 p_1 + \sigma_2 p_2 + \sigma_3 p_3 = \begin{pmatrix} p_3 & p_1 - i p_2 \\ p_1 + i p_2 & -p_3 \end{pmatrix} \quad (1.31)$$

acting on  $\mathbb{C}^2$ -valued  $C^\infty$  function  $h := {}^t(h_1, h_2)$  on  $\mathbb{R}^3$ , where the  $\sigma_j$ ,  $j = 1, 2, 3$ , are the Pauli matrices in (1.6), we have exactly the same result. For  $h := {}^t(h_1, h_2)$  whose four first-order derivatives  $(\partial_1 + i\partial_2)h_1$ ,  $\partial_3 h_1$ ,  $(\partial_1 - i\partial_2)h_2$  and  $\partial_3 h_2$  are  $p$ -th power integrable in  $\mathbb{R}^3$ , consider the semi-norm

$$\begin{aligned} M_{\sigma \cdot p; p}(h) &:= [\|(\sigma \cdot p)P_1 h\|_p^p + \|(\sigma \cdot p)P_2 h\|_p^p]^{1/p} \\ &= [\|(\partial_1 + i\partial_2)h_1\|_p^p + \|\partial_3 h_1\|_p^p + \|(\partial_1 - i\partial_2)h_2\|_p^p + \|\partial_3 h_2\|_p^p]^{1/p}, \end{aligned} \quad (1.32)$$



decomposing  $\sigma \cdot p$  into the sum of its two parts:

$$\sigma \cdot p = (\sigma \cdot p)P_1 + (\sigma \cdot p)P_2 = \begin{pmatrix} p_3 & 0 \\ p_1 + i p_2 & 0 \end{pmatrix} + \begin{pmatrix} 0 & p_1 - i p_2 \\ 0 & -p_3 \end{pmatrix}.$$

Here  $P_1 = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$  and  $P_2 = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$  are two projection matrices acting on the  $\mathbb{C}^2$  of two-vectors and note that  $P_1 + P_2 = I_2$  ( $2 \times 2$ -identity matrix). By the same argument as before around Proposition 1.0 for  $\alpha \cdot p$ , it is also seen that this semi-norm  $M_{\sigma \cdot p; p}(h)$  defined by (1.32) with the decomposition (1.31) of  $\sigma \cdot p$  coincides with the one to be defined with another decomposition:

$$\sigma \cdot p = \begin{pmatrix} 0 & p_1 - i p_2 \\ p_1 + i p_2 & 0 \end{pmatrix} + \begin{pmatrix} p_3 & 0 \\ 0 & -p_3 \end{pmatrix} =: (\sigma \cdot p)_1 + (\sigma \cdot p)_2,$$

i.e.  $M_{\sigma \cdot p; p}(h) = [\|(\sigma \cdot p)_1 h\|_p^p + \|(\sigma \cdot p)_2 h\|_p^p]^{1/p}$ .

The Banach spaces obtained as completions of  $C_0^\infty(\mathbb{R}^3; \mathbb{C}^2)$  by the norms  $\|h\|_{M_{\sigma \cdot p, 1, p}} := (\|h\|_p^p + M_{\sigma \cdot p; p}(h)^p)^{1/p}$  and  $\|h\|_{\sigma \cdot p, 1, p} := (\|h\|_p^p + \|(\sigma \cdot p)h\|_p^p)^{1/p}$  are denoted by  $H_{M_{\sigma \cdot p, 0}}^{1, p}(\mathbb{R}^3; \mathbb{C}^2)$ ,  $H_{(\sigma \cdot p), 0}^{1, p}(\mathbb{R}^3; \mathbb{C}^2)$ , respectively.

**Corollary 1.2.** (with 3-dimensional Weyl–Dirac operator) (i) *For  $1 \leq p < q < \infty$ , a  $\mathbb{C}^2$ -valued functions  $h = {}^t(h_1, h_2)$  belongs to  $L^q(\mathbb{R}^3; \mathbb{C}^2)$ , if  $h$  belongs to  $B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^2)$  and satisfies  $M_{\sigma \cdot p; p}(h) < \infty$ , and further, there exists a positive constant  $C$  such that*

$$\|h\|_q \leq C M_{\sigma \cdot p; p}(h)^{p/q} \|h\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}. \quad (1.33)$$

*Therefore this holds, in particular, for  $h \in H_{M_{\sigma \cdot p, 0}}^{1, p}(\mathbb{R}^3; \mathbb{C}^2) \cap B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^2)$ .*

(ii) *For  $\infty > p > 1$ , the three semi-norms  $\|(\sigma \cdot p)h\|_p$ ,  $M_{\sigma \cdot p; p}(h)$  and  $\|\nabla h\|_p$  are equivalent, so that the corresponding three Banach spaces coincide with one another:*

$$H_0^{1, p}(\mathbb{R}^3; \mathbb{C}^2) = H_{M_{\sigma \cdot p, 0}}^{1, p}(\mathbb{R}^3; \mathbb{C}^2) = H_{\sigma \cdot p, 0}^{1, p}(\mathbb{R}^3; \mathbb{C}^2). \quad (1.34)$$

*Therefore assertion (i) turns out: For  $1 < p < q < \infty$ , there exists a positive constant  $C$  such that*

$$\|h\|_q \leq C \|(\sigma \cdot p)h\|_p^{p/q} \|h\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.35)$$

*for every  $h \in B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^2)$  whose semi-norm  $\|(\sigma \cdot p)h\|_p$ ,  $M_{\sigma \cdot p; p}(h)$  or  $\|\nabla h\|_p$  is finite. Therefore this holds, in particular, for every  $f$  in the space (1.34) which belongs to  $B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^2)$ . (1.35) is equivalent to the vector-valued version (1.13) of (1.1) with  $n = 3$ .*

Second, for  $\mathbb{C}$ -valued  $C^\infty$  functions  $\psi$  whose two first-order derivatives  $(\partial_1 - i\partial_2)\psi$  and  $\partial_3\psi$  are  $p$ -th power integrable in  $\mathbb{R}^3$ , consider the semi-norm

$$M_{(\partial_1 - i\partial_2) \vee \partial_3; p}(\psi) := [\|(\partial_1 - i\partial_2)\psi\|_p^p + \|\partial_3\psi\|_p^p]^{1/p}. \quad (1.36)$$

The Banach space obtained as completion of  $C_0^\infty(\mathbb{R}^3)$  by the norm  $\|\psi\|_{M_{(\partial_1 - i\partial_2) \vee \partial_3, 1, p}} := (\|\psi\|_p^p + M_{(\partial_1 - i\partial_2) \vee \partial_3; p}(\psi)^p)^{1/p}$  is denoted by  $H_{M_{(\partial_1 - i\partial_2) \vee \partial_3, 0}}^{1, p}(\mathbb{R}^3)$ .

**Corollary 1.3.** (i) For  $1 \leq p < q < \infty$ , a function  $\psi$  belongs to  $L^q(\mathbb{R}^3)$ , if  $\psi$  belongs to  $B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3)$  and satisfies  $M_{(\partial_1-i\partial_2)\vee\partial_3;p}(\psi) < \infty$ , and further, there exists a positive constant  $C$  such that

$$\|\psi\|_q \leq C M_{(\partial_1-i\partial_2)\vee\partial_3;p}(\psi)^{p/q} \|\psi\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)}. \quad (1.37)$$

Therefore, in particular, for every  $\psi \in H_{M_{(\partial_1-i\partial_2)\vee\partial_3,0}}^{1,p}(\mathbb{R}^3) \cap B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3)$ .

(ii) For  $\infty > p > 1$ , the two semi-norms  $M_{(\partial_1-i\partial_2)\vee\partial_3;p}(\psi)$  and  $\|\nabla f\|_p$  are equivalent, so that the corresponding two Banach spaces coincide with each other:

$$H_0^{1,p}(\mathbb{R}^3; \mathbb{C}^2) = H_{M_{(\partial_1-i\partial_2)\vee\partial_3,0}}^{1,p}(\mathbb{R}^3; \mathbb{C}^2). \quad (1.38)$$

Therefore assertion (i) turns out: For  $1 < p < q < \infty$ , there exists a positive constant  $C$  such that

$$\|\psi\|_q \leq C M_{(\partial_1-i\partial_2)\vee\partial_3;p}(\psi)^{p/q} \|\psi\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.39)$$

for every  $f \in B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3)$  whose semi-norm  $M_{(\partial_1-i\partial_2)\vee\partial_3;p}(\psi)$  or  $\|\nabla f\|_p$  is finite. Therefore this holds, in particular, for every  $f$  in the space (1.38) which belongs to  $B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^3)$ . (1.39) is equivalent to the vector-valued version (1.13) of (1.1) with  $n = 2$ .

Third, we shall consider the *two-dimensional* Weyl–Dirac (or Pauli) operators made from two of the three Pauli matrices (1.6). There are the following three:

$$(\sigma \cdot p)^{(a)} f := (\sigma_1 p_1 + \sigma_2 p_2) f = \begin{pmatrix} 0 & p_1 - i p_2 \\ p_1 + i p_2 & 0 \end{pmatrix} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}, \quad (1.40a)$$

$$(\sigma \cdot p)^{(b)} f := (\sigma_3 p_1 + \sigma_1 p_2) f = \begin{pmatrix} p_1 & p_2 \\ p_2 & -p_1 \end{pmatrix} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}, \quad (1.40b)$$

$$(\sigma \cdot p)^{(c)} f := (\sigma_3 p_1 + \sigma_2 p_2) f = \begin{pmatrix} p_1 & -i p_2 \\ i p_2 & -p_1 \end{pmatrix} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}, \quad (1.40c)$$

for  $f := {}^t(f_1, f_2)$ . As we shall see later in Lemma 5.1, these three operators  $(\sigma \cdot p)^{(a)}$ ,  $(\sigma \cdot p)^{(b)}$ ,  $(\sigma \cdot p)^{(c)}$  are unitarily equivalent, so that the three semi-norms  $\|(\sigma \cdot p)^{(a)} f\|_p$ ,  $\|(\sigma \cdot p)^{(b)} f\|_p$ ,  $\|(\sigma \cdot p)^{(c)} f\|_p$  are equivalent. Therefore we write any of these three operators as  $(\sigma \cdot p)^{(2)}$  so as to distinguish it from the *three-dimensional* Weyl–Dirac (or Pauli) operator  $\sigma \cdot p$  in (1.31), and any of these semi-norms as  $\|(\sigma \cdot p)^{(2)} f\|_p$  to consider the norm  $\|f\|_{(\sigma \cdot p)^{(2)}, 1, p} := (\|f\|_p^p + \|(\sigma \cdot p)^{(2)} f\|_p^p)^{1/p}$ . What can be shown just in the same way as in [IS] is that the Banach space  $H_{(\sigma \cdot p)^{(2)}, 0}^{1,p}(\mathbb{R}^2; \mathbb{C}^2)$  obtained as completion of  $C_0^\infty(\mathbb{R}^2; \mathbb{C}^2)$  in this norm coincides for  $1 < p < \infty$  with the Sobolev spaces  $H_0^{1,p}(\mathbb{R}^2; \mathbb{C}^2) = H^{1,p}(\mathbb{R}^2; \mathbb{C}^2)$ , but is for  $p = 1$  strictly larger. Differing from Corollary 1.2 for 3-dimensional case, the following theorem for 2-dimensional case gives a *true* extension of inequality (1.1) for single-valued functions to the case for vector-valued functions.

**Theorem 1.4.** (with 2-dimensional Weyl–Dirac (or Pauli) operator) For  $1 \leq p < q < \infty$  there exists a positive constant  $C$  such that

$$\|f\|_q \leq C \|(\sigma \cdot p)^{(2)} f\|_p^{p/q} \|f\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)} \quad (1.41)$$

for every  $f \in B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^2; \mathbb{C}^2)$  which satisfies  $\|(\sigma \cdot p)^{(2)} f\|_p < \infty$ . Therefore this holds, in particular, for every  $f \in H_{(\sigma \cdot p)^{(2)},0}^{1,p}(\mathbb{R}^2; \mathbb{C}^2) \cap B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^2; \mathbb{C}^2)$ .

Forth, from Corollary 1.3 or Theorem 1.4 we can get the following inequality involved with the Cauchy–Riemann operator  $\frac{1}{2}(\partial_1 + i\partial_2)$  in  $\mathbb{R}^2$ .

**Corollary 1.5.** (with Cauchy–Riemann operator) *For  $1 \leq p < q < \infty$ , there exists a positive constant  $C$  such that*

$$\|\psi\|_q \leq C \|(\partial_1 + i\partial_2)\psi\|_p^{p/q} \|\psi\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)} \quad (1.42)$$

for every  $\psi \in B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^2)$  which satisfies  $\|(\partial_1 + i\partial_2)\psi\|_p < \infty$ .

Finally, we are going to consider the four-dimensional Euclidian Dirac operator

$$\beta \cdot p = \sum_{k=1}^4 \beta_k \cdot p_k = -i \sum_{k=1}^4 \beta_k \partial_k, \quad (1.43)$$

with  $p = (p_1, p_2, p_3, p_4)$ ,  $p_k = -i\partial_k$ ,  $k = 1, 2, 3, 4$ , which acts on  $\mathbb{C}^4$ -valued functions  $f(x) = {}^t(f_1(x), f_2(x), f_3(x), f_4(x))$  defined in 4-dimensional Euclidian space-time  $\mathbb{R}^4$ . Here we are using the symbol  $\beta$  for a quadruple  $\beta := (\beta_1, \beta_2, \beta_3, \beta_4)$  of the Dirac matrices which are  $4 \times 4$  Hermitian matrices satisfying the anti-commutation relation  $\beta_j \beta_k + \beta_k \beta_j = \delta_{jk} I_4$ ,  $j, k = 1, 2, 3, 4$ . As the first three of it, we take here, with the same triple of Pauli matrices as in (1.6),

$$\beta_j := \alpha_j = \begin{pmatrix} 0_2 & \sigma_j \\ \sigma_j & 0_2 \end{pmatrix} \quad (j = 1, 2, 3), \quad (1.44)$$

and, as the fourth  $\beta_4$ , we adopt

$$\beta_4 := \alpha_5 = \begin{pmatrix} 0_2 & -iI_2 \\ iI_2 & 0_2 \end{pmatrix}, \quad (1.45)$$

but not the usual  $\alpha_4$  given by

$$\alpha_4 = \begin{pmatrix} I_2 & 0_2 \\ 0_2 & -I_2 \end{pmatrix}.$$

The  $\alpha_4$  is often written as “ $\beta$ ”, but of course, different from our  $\beta$  on the left-hand side of (1.41) above (e.g. [BeSa, p.48]). For this, see e.g. [W] where  $\alpha_5$  is given as in (1.45) and read in [ItZ, p.693] as  $\alpha_5 := i\gamma^5\gamma^0 = \alpha_1\alpha_2\alpha_3\alpha_4$  (see also [G]). Note that as the five  $\alpha_k$ ,  $k = 1, 2, 3, 4, 5$ , are mutually anti-commuting, Hermitian matrices satisfying  $\alpha_j \alpha_k + \alpha_k \alpha_j = 2\delta_{jk} I_4$ ,  $j, k = 1, 2, 3, 4$ , so are the four  $\beta_k$ ,  $k = 1, 2, 3, 4$ . (Here  $\delta_{jk}$  is the usual Kronecker delta, one when the indices are the same, otherwise one.) Therefore  $\beta \cdot p = \sum_{k=1}^4 \beta_k \cdot p_k$  is a selfadjoint operator in  $L^2(\mathbf{R}^4; \mathbb{C}^4)$  as well as  $\sum_{k=1}^4 \alpha_k p_k$ .

Then similarly to the 3-dimensional case before (see around (1.18)), we consider the semi-norm  $M_{\beta \cdot p; p}(f)$  as well as the semi-norm  $\|(\beta \cdot p)f\|$  concerning the first-order derivatives of functions of functions  $f = {}^t(f_1, f_2, f_3, f_4)$  in the space  $C_0^\infty(\mathbb{R}^4; \mathbb{C}^4)$ . To define  $M_{\beta \cdot p; p}(f)$ , note first that the 4-dimensional Euclidian Dirac operator (1.43) can

be rewritten, based on the representation (1.44) together with (1.45) for the matrices  $\beta_k$ ,  $k = 1, 2, 3, 4$ , as

$$\beta \cdot p = \begin{pmatrix} 0 & 0 & p_3 - i p_4 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & -(p_3 + i p_4) \\ p_3 + i p_4 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix}. \quad (1.46)$$

Then decompose it into the sum of its two parts:

$$\begin{aligned} \beta \cdot p &= (\beta \cdot p)P_{13} + (\beta \cdot p)P_{24} \\ &= \begin{pmatrix} 0 & 0 & p_3 - i p_4 & 0 \\ 0 & 0 & p_1 + i p_2 & 0 \\ p_3 + i p_4 & 0 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & 0 & -(p_3 + i p_4) \\ 0 & p_1 - i p_2 & 0 & 0 \\ 0 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix}, \end{aligned} \quad (1.47)$$

where  $P_{13} := \text{diag}(1, 0, 1, 0)$  and  $P_{24} := \text{diag}(0, 1, 0, 1)$  are the same two projection matrices acting on the space  $\mathbb{C}^4$  of four-vectors as before, and define

$$M_{\beta \cdot p; p}(f) := [\|(\beta \cdot p)P_{13}f\|_p^p + \|(\beta \cdot p)P_{24}f\|_p^p]^{1/p}. \quad (1.48)$$

Let us see how this semi-norm  $M_{\beta \cdot p; p}(f)$  in (1.48) is related to the other semi-norms  $\|(\beta \cdot p)f\|_p$  and  $\|\nabla f\|_p$ . However, we should note here that the latter  $\|\nabla f\|_p$  differs from (1.11), since in the present case we have the 4-dimensional gradient  $\nabla = (\partial_1, \partial_2, \partial_2, \partial_2)$ , so that  $|\nabla f(x)|_{\ell^p}^p := \sum_{j=1}^4 |\partial_j f(x)|_{\ell^p}^p = \sum_{j=1}^4 \sum_{k=1}^4 |\partial_j f_k(x)|^p$ .

Then

$$(\beta \cdot p)f = \begin{pmatrix} (p_3 - i p_4)f_3 + (p_1 - i p_2)f_4 \\ (p_1 + i p_2)f_3 - (p_3 + i p_4)f_4 \\ (p_3 + i p_4)f_1 + p_1 - i p_2)f_2 \\ (p_1 + i p_2)f_1 - (p_3 - i p_4)f_2 \end{pmatrix},$$

so that, recalling the definition of the  $\ell^p$  norm in (1.9), we have

$$\begin{aligned} |(\beta \cdot p)f|_{\ell^p}^p &= |(p_3 - i p_4)f_3 + (p_1 - i p_2)f_4|^p + |(p_1 + i p_2)f_3 - (p_3 + i p_4)f_4|^p \\ &\quad + |(p_3 + i p_4)f_1 + (p_1 - i p_2)f_2|^p + |(p_1 + i p_2)f_1 - (p_3 - i p_4)f_2|^p \\ &= |(p_1 + i p_2)f_1 - (p_3 - i p_4)f_2|^p + |(p_1 - i p_2)f_2 + (p_3 + i p_4)f_1|^p \\ &\quad + |(p_1 + i p_2)f_3 - (p_3 + i p_4)f_4|^p + |(p_1 - i p_2)f_4 + (p_3 - i p_4)f_3|^p, \end{aligned}$$

where we have rearranged the four terms, when passing through the second equality. Hence

$$\begin{aligned} \|(\beta \cdot p)f\|_p^p &= \|(\partial_1 + i\partial_2)f_1 - (\partial_3 - i\partial_4)f_2\|_p^p + \|(\partial_1 - i\partial_2)f_2 + (\partial_3 + i\partial_4)f_1\|_p^p \\ &\quad + \|(\partial_1 + i\partial_2)f_3 - (\partial_3 + i\partial_4)f_4\|_p^p + \|(\partial_1 - i\partial_2)f_4 + (\partial_3 - i\partial_4)f_3\|_p^p. \end{aligned} \quad (1.49)$$

Then one can calculate the right-hand side of (1.48) to get

$$\begin{aligned} &M_{\beta \cdot p; p}(f)^p \\ &= \|(\beta \cdot p)P_{13}f\|_p^p + \|(\beta \cdot p)P_{24}f\|_p^p \\ &= (\|(\partial_1 + i\partial_2)f_1\|_p^p + \|(\partial_3 + i\partial_4)f_1\|_p^p) + (\|(\partial_1 - i\partial_2)f_2\|_p^p + \|(\partial_3 - i\partial_4)f_2\|_p^p) \\ &\quad + (\|(\partial_1 + i\partial_2)f_3\|_p^p + \|(\partial_3 - i\partial_4)f_3\|_p^p) + (\|(\partial_1 - i\partial_2)f_4\|_p^p + \|(\partial_3 + i\partial_4)f_4\|_p^p). \end{aligned} \quad (1.50)$$

Similarly to the 3-dimensional case before (see (1.22), (1.23)), for the semi-norms (1.49) and (1.48)/(1.50) we have with  $1 \leq p < \infty$ ,

$$2^{-(1-(1/p))} \|(\beta \cdot p)f\|_p \leq M_{\beta \cdot p; p}(f) \leq 2^{1-(1/p)} \|\nabla f\|_p. \quad (1.51)$$

The Banach space  $H_{(\beta \cdot p), 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4) / H_{M_{\beta \cdot p}, 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4)$  is defined as completion of the space  $C_0^\infty(\mathbb{R}^4; \mathbb{C}^4)$  in the norm  $\|f\|_{(\beta \cdot p), 1, p} := (\|f\|_p^p + \|(\beta \cdot p)f\|_p^p)^{1/p} / \|f\|_{M_{\beta \cdot p}, 1, p} := (\|f\|_p^p + M_{\beta \cdot p; p}(f)^p)^{1/p}$ . From (1.51) we see the following inclusion relation :

$$H_0^{1,p}(\mathbb{R}^4; \mathbb{C}^4) \subseteq H_{M_{\beta \cdot p}, 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4) \subseteq H_{(\beta \cdot p), 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4). \quad (1.52)$$

Now we note the semi-norm  $M_{\beta \cdot p; p}(f)$  has a significant character as that of  $M_{\alpha \cdot p; p}(f)$  in Proposition 1.0, by considering other decompositions of the Euclidian Dirac operator  $\beta \cdot p$  in (1.46), than (1.47), into the sum of its two parts:

$$\underline{M_{\beta}^{(1)}}$$

$$\begin{aligned} \beta \cdot p &= (\beta \cdot p)P_{14} + (\beta \cdot p)P_{23} \\ &\equiv \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & 0 & -(p_3 + i p_4) \\ p_3 + i p_4 & 0 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & p_3 - i p_4 & 0 \\ 0 & 0 & p_1 + i p_2 & 0 \\ 0 & p_1 - i p_2 & 0 & 0 \\ 0 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix}, \end{aligned} \quad (1.53)$$

where  $P_{14} := \text{diag}(1, 0, 0, 1)$  and  $P_{23} := \text{diag}(0, 1, 1, 0)$  are the same two projection matrices acting on the space  $\mathbb{C}^4$  of four-vectors as before, and note that both the operators  $(\beta \cdot p)P_{14}$  and  $(\beta \cdot p)P_{23}$  on the right are selfadjoint, i.e.  $((\beta \cdot p)P_{14})^* = (\beta \cdot p)P_{14}$ ,  $((\beta \cdot p)P_{23})^* = (\beta \cdot p)P_{23}$ .

$$\underline{M_{\beta}^{(2)}}$$

$$\begin{aligned} \beta \cdot p &= \begin{pmatrix} 0_2 & \sigma_1 p_1 + \sigma_2 p_2 \\ \sigma_3 p_3 + i I_2 p_4 & 0_2 \end{pmatrix} + \begin{pmatrix} 0_2 & \sigma_3 p_3 - i I_2 p_4 \\ \sigma_1 p_1 + \sigma_2 p_2 & 0_2 \end{pmatrix} \\ &\equiv \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & 0 \\ p_3 - i p_4 & 0 & 0 & 0 \\ 0 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix} \\ &\quad + \begin{pmatrix} 0 & 0 & p_3 - i p_4 & 0 \\ 0 & 0 & 0 & -(p_3 + i p_4) \\ 0 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} \\ &=: (\beta \cdot p)_1 + (\beta \cdot p)_2, \end{aligned} \quad (1.54)$$

where note that  $(\beta \cdot p)_2$  is the adjoint of  $(\beta \cdot p)_1$  as operators, say, in  $L^2(\mathbf{R}^3; \mathbb{C}^4)$ , i.e.  $(\beta \cdot p)_2 = (\beta \cdot p)_1^*$ .

$$\underline{M_{\beta}^{(3)}}$$

$$\begin{aligned}
\beta \cdot p &= \begin{pmatrix} 0 & \sigma_1 p_1 + \sigma_2 p_2 \\ \sigma_1 p_1 + \sigma_2 p_2 & 0 \end{pmatrix} + \begin{pmatrix} 0 & \sigma_3 p_3 - i I_2 p_4 \\ \sigma_3 p_3 + i I_2 p_4 & 0 \end{pmatrix} \\
&\equiv \begin{pmatrix} 0 & 0 & 0 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & 0 \\ 0 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} \\
&\quad + \begin{pmatrix} 0 & 0 & p_3 - i p_4 & 0 \\ 0 & 0 & 0 & -(p_3 + i p_4) \\ p_3 + i p_4 & 0 & 0 & 0 \\ 0 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix} \\
&=: (\beta \cdot p)_3 + (\beta \cdot p)_4, \tag{1.55}
\end{aligned}$$

where note that both the operators  $(\beta \cdot p)_3$  and  $(\beta \cdot p)_4$  on the right are selfadjoint.

Then we can confirm, in the same way as in Proposition 1.0 for  $M_{\alpha \cdot p}(f)$  with  $\alpha \cdot p$ , that the semi-norm  $M_{\alpha \cdot p}(f)$  of  $f$  defined by (1.48) with the rather artificial decomposition (1.47) turns out to be equal to the ones to be defined with the other decompositions (1.53), (1.54) and (1.55), taking account of the expression (1.50) for  $M_{\beta \cdot p}(f)$ :

$$M_{\beta \cdot p}^{(1)}(f) := [\|(\alpha \cdot p)P_{14}f\|_p^p + \|(\alpha \cdot p)P_{23}f\|_p^p]^{1/p}; \tag{1.56a}$$

$$M_{\beta \cdot p}^{(2)}(f) := [\|(\beta \cdot p)_1 f\|_p^p + \|(\alpha \cdot p)_2 f\|_p^p]^{1/p} = [\|(\beta \cdot p)_1 f\|_p^p + \|(\beta \cdot p)_1^* f\|_p^p]^{1/p}; \tag{1.56b}$$

$$M_{\beta \cdot p}^{(3)}(f) := [\|(\beta \cdot p)_3 f\|_p^p + \|(\beta \cdot p)_4 f\|_p^p]^{1/p}. \tag{1.56c}$$

Further, more generally, every decomposition of  $\beta \cdot p$  into its two parts,  $\beta \cdot p = (\beta \cdot p)_5 + (\beta \cdot p)_6$ , such that each row of both the matrices  $(\beta \cdot p)_5$  and  $(\beta \cdot p)_6$  contains only one nonzero entry, defines the semi-norm  $M_{\beta \cdot p}(f)$  which has the expression (1.50). However, as mentioned for the operator  $\alpha \cdot p$  after Proposition 1.0, the decomposition  $\beta \cdot p = (\beta \cdot p)P_{12} + (\beta \cdot p)P_{34}$  is not fit for the semi-norm  $M_{\beta \cdot p}(f)$ , to which we will come back in Section 6 to discuss the issue.

**Theorem 1.6.** (with 4-dimensional Euclidian Dirac operator). (i) For  $1 \leq p < q < \infty$ , a  $\mathbb{C}^4$ -valued function  $f = {}^t(f_1, f_2, f_3, f_4)$  belongs to  $L^q(\mathbb{R}^4; \mathbb{C}^4)$ , if  $f$  belongs to  $B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^4; \mathbb{C}^4)$  and satisfies  $M_{\beta \cdot p}(f) < \infty$ , and further, there exists a positive constant  $C$  such that

$$\|f\|_q \leq C M_{\beta \cdot p}(f)^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}. \tag{1.57}$$

Therefore this holds, in particular, for every  $f \in H_{M_{\beta \cdot p}, 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4) \cap B_{\infty, \infty}^{p/(p-q)}(\mathbb{R}^4; \mathbb{C}^4)$ .

(ii) For  $\infty > p > 1$ , the three semi-norms  $\|(\beta \cdot p)f\|_p$ ,  $M_{\beta \cdot p}(f)$  and  $\|\nabla f\|_p$  are equivalent, so that the corresponding three Banach spaces (1.52) coincide with one another:

$$H_0^{1,p}(\mathbb{R}^4; \mathbb{C}^4) = H_{M_{\beta \cdot p}, 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4) = H_{(\beta \cdot p), 0}^{1,p}(\mathbb{R}^4; \mathbb{C}^4). \tag{1.58}$$

Therefore assertion (i) turns out: For  $1 < p < q < \infty$ , there exists a positive constant  $C$  such that

$$\|f\|_q \leq C \|(\beta \cdot p)f\|_p^{p/q} \|f\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.59)$$

for every  $f \in B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^4; \mathbb{C}^4)$  whose semi-norm  $\|(\beta \cdot p)f\|_p$ ,  $M_{\beta \cdot p; p}(f)$  or  $\|\nabla f\|_p$  is finite. Therefore this holds, in particular, for every  $f$  in the above space (1.58) which belongs to  $B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^4; \mathbb{C}^4)$ . (1.59) is equivalent to the vector-valued version (1.13) of (1.1) with  $n = 4$ .

We note here that the 4-dimensional Euclidian Dirac operator  $\sum_{k=1}^4 \beta_k p_k$  in (1.47) turns, if  $\beta_4 p_4 = -i\beta_4 \partial_4$  is removed from it, the 3-dimensional massless Dirac operator  $\sum_{j=1}^3 \alpha_j p_j$  in (1.17), which reduces Theorem 1.6 to Theorem 1.1.

Finally, as is the case for Sobolev spaces of single-valued functions, it is seen for the two spaces of vector-valued functions which we introduced in (1.23) and (1.52) that each of them coincides with the following two spaces:

$$\begin{aligned} H_{M_{\alpha \cdot p; 0}}^{1,p}(\mathbb{R}^3; \mathbb{C}^4) &= H_{M_{\alpha \cdot p; p}}^{1,p}(\mathbb{R}^3; \mathbb{C}^4) \\ &= \{f \in L^p(\mathbb{R}^3; \mathbb{C}^4); (\alpha \cdot p)P_{13}f, (\alpha \cdot p)P_{24}f \in L^p(\mathbb{R}^3; \mathbb{C}^4)\} \\ &= \{f \in L^p(\mathbb{R}^3; \mathbb{C}^4); (\alpha \cdot p)_1 f, (\alpha \cdot p)_2 f \in L^p(\mathbb{R}^3; \mathbb{C}^4)\}; \\ H_{M_{\beta \cdot p; 0}}^{1,p}(\mathbb{R}^4; \mathbb{C}^4) &= H_{M_{\beta \cdot p; p}(f)}^{1,p}(\mathbb{R}^4; \mathbb{C}^4) \\ &= \{f \in L^p(\mathbb{R}^4; \mathbb{C}^4); (\beta \cdot p)P_{13}f, (\beta \cdot p)P_{24}f \in L^p(\mathbb{R}^4; \mathbb{C}^4)\} \\ &= \{f \in L^p(\mathbb{R}^4; \mathbb{C}^4); (\beta \cdot p)_1 f, (\beta \cdot p)_2 f \in L^p(\mathbb{R}^4; \mathbb{C}^4)\}. \end{aligned}$$

In each of these two formulas, the second space is the Banach space obtained as completion with respect to the norm  $\|f\|_{M_{\alpha \cdot p; 1,p}}$  [resp.  $\|f\|_{M_{\beta \cdot p; 1,p}}$ ] of the linear space of all  $f \in C^\infty(\mathbb{R}^3; \mathbb{C}^4) \cap L^p(\mathbb{R}^3; \mathbb{C}^4)$  [resp.  $C^\infty(\mathbb{R}^4; \mathbb{C}^4) \cap L^p(\mathbb{R}^4; \mathbb{C}^4)$ ]. In the third and fourth spaces the first-order derivatives are taken in the distribution sense.

The proof of the improved Sobolev inequality (1.1) for single-valued functions in [CDPX] and [CMO] was based on wavelet analysis, while Ledoux [Le] made a different approach by a direct semigroup argument. We do our proof, modifying the method used by Ledoux so as to be able to apply to vector-valued functions.

The plan of this paper is as follows. Section 2 collects remarks to the results, stated in Section 1, for vector-valued functions to compare them with the improved Sobolev inequality (1.1) and the Dirac–Sobolev inequality (1.14) obtained in [BES]. Section 3 gives examples where the simple-minded, vector-valued version (1.16) connected not only with the three-dimensional massless Dirac operator but also with the four-dimensional Euclidian Dirac operator fails to hold for  $p = 1$ . In Section 4, we give proof of Theorem 1.1, and in Section 5, proofs of all the other five Corollaries 1.2, 1.3, Theorem 1.4, Corollary 1.5, Theorem 1.6. In Section 6 we make concluding comments on the first-order-derivative semi-norm connected with the Dirac operators which we have introduced in Section 1. It is defined at first with a rather artificial decomposition of the Dirac operator into two parts, but later turns out to be meaningful enough to have universal character. The final Section 7 briefly summarizes all our results to exhibit their significance and difference from the case of single-valued functions.

## 2 Remarks

1°. Theorem 1.1 (i) (ii): We compare our inequality (1.28) with (1.16)/(1.30), the trivial version (1.13) and the first vector-valued one (1.14) of inequality (1.1) shown in [BES].

To do so, first we collect the results of equivalence and non-equivalence among the three first-order-derivative semi-norms  $\|\nabla f\|_p$  in (1.11),  $\|(\alpha \cdot p)f\|_p$  in (1.9) and  $M_{\alpha \cdot p; p}(f)$  in (1.18), which are under relation (1.22). When  $1 < p < \infty$ , these three are all equivalent, which we shall see in the proof of Theorem 1.1 (ii) in Section 3 below, but different when  $p = 1$ . In this case  $p = 1$ , we showed non-equivalence between  $\|\nabla f\|_1$  and  $\|(\alpha \cdot p)f\|_1$  in [IS, Theorem 1.3 (iii)]. Non-equivalence between  $\|(\alpha \cdot p)f\|_1$  and  $M_{\alpha \cdot p; 1}(f)$  can be seen in view of their respective explicit expressions (1.20) and (1.21), and that between  $\|\nabla f\|_1$  and  $M_{\alpha \cdot p; 1}(f)$  in view of their respective definition (1.11) and explicit expression (1.21), both from the fact that (2.2) below cannot hold. In particular, the two inclusions in (1.23) are strict.

Next we going to observe the difference and coincidence among inequalities (1.28), (1.16)/(1.30), (1.13) and (1.14). For  $1 < p < \infty$ , the first three, i.e. (1.28), (1.16)/(1.30) and (1.13), are equivalent, and strictly sharper than and hence an improvement of the last one, (1.14). The former is because of equivalence of the three first-order-derivative semi-norms concerned as just seen above, and the latter because the  $L^q$  norm  $\|f\|_q$  on the left of (1.28) is stronger than the weak  $L^q$  norm  $\|f\|_{q, \infty}$  on the left of (1.14). For  $p = 1$ , (1.16)/(1.28) does not hold in general, and (1.28) is sharper than (1.13), because the semi-norm  $M_{\alpha \cdot p; 1}(f)$  on the right of (1.28) is weaker than the semi-norm  $\|\nabla f\|_1$  on the right of (1.13). In the case  $p = 1$ , however, two inequalities (1.28) and (1.14) cannot be compared so as to say which of them is sharper, because  $M_{\alpha \cdot p; 1}(f)$  on the right of (1.28) is not weaker than  $\|(\alpha \cdot p)f\|_1$  on the right of (1.14), though  $\|f\|_q$  on the left of (1.28) is stronger than  $\|f\|_{q, \infty}$  on the left of (1.14). As a result, (1.28) for  $p = 1$  is a new inequality for vector-valued version of (1.1).

2°. Corollary 1.2 (i) (ii): The same remark as 1° above applies to the case for the 3-dimensional Weyl–Dirac (or Pauli) operator  $\sigma \cdot p$  in place of the Dirac operator  $\alpha \cdot p$ .

3°. Corollary 1.3 (i) (ii): For  $p = 1$ , the semi-norm  $M_{(\partial - i\partial_2) \vee \partial_3; 1}(\psi)$  in (1.36) is bounded by the semi-norm  $\|\nabla \psi\|_1$ , i.e.

$$M_{(\partial - i\partial_2) \vee \partial_3; 1}(\psi) \leq \|\nabla \psi\|_1, \quad (2.1)$$

but not reversely (See [St, pp.59–60, III, Propositions 3, 4, and p.48, 6.1] and [IS, Lemma 4.3]). Therefore the Banach space  $H_{M_{(\partial - i\partial_2) \vee \partial_3; 0}}^{1, p}(\mathbb{R}^3; \mathbb{C}^2)$  obtained as completion of  $C_0^\infty(\mathbb{R}^3)$  with respect to the norm  $\|\psi\|_{M_{(\partial - i\partial_2) \vee \partial_3; 1, p}} := \|\psi\|_1 + M_{(\partial - i\partial_2) \vee \partial_3; p}(\psi)$  is strictly larger than the space  $H_0^{1, 1}(\mathbb{R}^3)$ . Therefore for  $p = 1$ , Corollary 1.3 (i) gives a slightly more general result than (1.1) of Ledoux [Le] though only in the case  $n = 3$ . However, for  $1 < p < \infty$ , it is nothing but his result though our result only concerns the case  $n = 3$ , since the semi-norm  $M_{(\partial - i\partial_2) \vee \partial_3; p}(\psi)$  is equivalent to the semi-norm  $\|\nabla \psi\|_p$ . In this sense, therefore our inequality (1.37) for  $\mathbb{C}$ -valued functions  $\psi$  is more general, though only for  $n = 3$ . Here it should be noted that it holds that for  $1 < p < \infty$ ,

$$\|\partial_1 \psi\|_p + \|\partial_2 \psi\|_p \leq C_p \|(\partial_1 - i\partial_2)\psi\|_p, \quad (2.2)$$



for all  $\psi \in C_0^\infty(\mathbb{R}^2)$  with a positive constant  $C_p$ , but cannot for  $p = 1$  (cf. [St, pp.59–60, III, Propositions 3, 4, and p.48, 6.1] and [IS, Lemma 4.3]). Therefore (2.2) implies that for  $1 < p < \infty$ ,

$$\begin{aligned} \|\nabla\psi\|_p &\equiv \left(\sum_{j=1}^3 \|\partial_j\psi\|_p^p\right)^{1/p} \leq (C_p^{p/(p-1)} + 1)^{(p-1)/p} [\|(\partial_1 - i\partial_2)\psi\|_p^p + \|\partial_3\psi\|_p^p]^{1/p} \\ &\equiv (C_p^{p/(p-1)} + 1)^{(p-1)/p} M_{(\partial_1 - i\partial_2) \vee \partial_3; p}(\psi), \end{aligned} \quad (2.3)$$

so that the two semi-norms  $M_{(\partial_1 - i\partial_2) \vee \partial_3; p}(\psi)$  and  $\|\nabla\psi\|_p$  are equivalent.

4°. Corollary 1.5: By analogous discussion made in Remark 3<sup>0</sup> to Corollary 1.3, (1.42) is also more general than (1.1) with  $n = 2$  for  $p = 1$ , but equivalent to it for  $\infty > p > 1$ .

5°. Theorem 1.6 and again Theorem 1.1: It can be seen that these two theorems hold also for some different representations of the 3-dimensional massless Dirac operator and 4-dimensional Euclidian Dirac operator than (1.17) and (1.46).

In fact, consider first the 4-dimensional Euclidian Dirac operators. Let  $\beta' = (\beta'_1, \beta'_2, \beta'_3, \beta'_4)$  be another quadruple of anti-commuting, Hermitian  $4 \times 4$ -matrices satisfying  $\beta'_j \beta'_k + \beta'_k \beta'_j = 2\delta_{jk} I_4$ ,  $j, k = 1, 2, 3, 4$ . Then Theorem 1.6 holds for the Euclidian Dirac operator  $\beta' \cdot p = \sum_{k=1}^4 \beta'_k p_k$  with corresponding projections  $P'_{13}, P'_{24}$ . Indeed, by the ‘fundamental theorem’ in [P, p.8] or [G, p.190], there exists a non-singular  $4 \times 4$ -matrix  $S$  such that  $\beta'_k = S\beta_k S^{-1}$  for  $k = 1, 2, 3, 4$ . So  $S$  is a similarity transformation which maps  $\mathbb{C}^4$  one-to-one onto  $\mathbb{C}^4$ , and in fact can be take to be a unitary matrix, because the  $\beta_k$  and  $\beta'_k$  are Hermitian. Then

$$\beta \cdot p = S^{-1}(\beta' \cdot p)S, \quad (\beta \cdot p)P_{13} = S^{-1}(\beta' \cdot p)P'_{13}S, \quad (\beta \cdot p)P_{24} = S^{-1}(\beta' \cdot p)P'_{24}S,$$

where  $P'_{13} := SP_{13}S^{-1}$  and  $P'_{24} := SP_{24}S^{-1}$  are projection matrices acting on  $\mathbb{C}^4$  such that  $P'_{13} + P'_{24} = I_4$ . It implies equivalence of the related semi-norms concerning  $\beta' \cdot p$  and  $\beta \cdot p$  in the following sense:

$$\begin{aligned} (\|S^{-1}\|_{\ell^p \rightarrow \ell^p})^{-1} \|(\beta \cdot p)f\|_p &\leq \|(\beta' \cdot p)(Sf)\|_p \leq \|S\|_{\ell^p \rightarrow \ell^p} \|(\beta \cdot p)f\|_p, \\ (\|S^{-1}\|_{\ell^p \rightarrow \ell^p})^{-1} \|(\beta \cdot p)P_{13}f\|_p &\leq \|(\beta' \cdot p)P'_{13}(Sf)\|_p \leq \|S\|_{\ell^p \rightarrow \ell^p} \|(\beta \cdot p)P_{13}f\|_p, \\ (\|S^{-1}\|_{\ell^p \rightarrow \ell^p})^{-1} \|(\beta \cdot p)P_{24}f\|_p &\leq \|(\beta' \cdot p)P'_{24}(Sf)\|_p \leq \|S\|_{\ell^p \rightarrow \ell^p} \|(\beta \cdot p)P_{24}f\|_p, \end{aligned}$$

with  $1 \leq p < \infty$ , where  $f = {}^t(f_1, f_2, f_3, f_4)$ , which yields equivalence of the semi-norms  $M_{\beta' \cdot p; p}(Sf)$  and  $M_{\beta \cdot p; p}(f)$ :

$$C_p^{-1} M_{\beta \cdot p; p}(f) \leq M_{\beta' \cdot p; p}(Sf) \leq C_p M_{\beta \cdot p; p}(f)$$

with a positive constant  $C_p$  depending on  $p$ . In particular, all this holds also for the 4-dimensional Euclidian Dirac operator  $\sum_{k=1}^4 \alpha_k p_k$ .

Though above we have dealt only the case corresponding to decomposition (1.47) of  $\beta \cdot p$ , the same is true for the cases corresponding to the other decompositions (1.53), (1.54) or (1.55).

Next, for Theorem 1.1, the same is valid, if one may consider, for  $\alpha' = (\alpha'_1, \alpha'_2, \alpha'_3)$  another triple of anti-commuting, Hermitian  $4 \times 4$ -matrices satisfying  $\alpha'_j \alpha'_k + \alpha'_k \alpha'_j = 2\delta_{jk} I_4$ ,  $j, k = 1, 2, 3$ , the Dirac operator  $\alpha' \cdot p = \sum_{j=1}^3 \alpha'_j p_j$  together with the corresponding projection matrices  $P'_{13}, P'_{24}$  to introduce the related semi-norms.

### 3 Counterexamples for $p = 1$

Inequalities of the type (1.16), i.e. (1.30) of Theorem 1.1 for the three-dimensional massless Dirac operator  $\alpha \cdot p$ , (1.35) of Corollary 1.2 with 3-dimensional Weyl–Dirac (or Pauli) operator  $\sigma \cdot p$ , (1.59) of Theorem 1.6 with 4-dimensional Euclidian Dirac operator  $\beta \cdot p$ , do not in general hold for  $p = 1$ , although they do for  $1 < p < \infty$ . This is why, for  $p = 1$ , we had to introduce the intermediate first-order-derivative semi-norms  $M_{\alpha \cdot p; p}(f)$  in (1.19),  $M_{\sigma \cdot p; p}(h)$  in (1.32),  $M_{\beta \cdot p; p}(f)$  in (1.48). Here, before going further, we keep Theorem 1.4 in mind that *nevertheless it holds for all  $1 \leq p < \infty$  with the 2-dimensional Weyl–Dirac (or Pauli) operator  $(\sigma \cdot p)^{(2)}$* , i.e. (1.40abc).

In this section, following the idea in the recent paper [BEU] for the 3-dimensional Weyl–Dirac (or Pauli) operator, we construct counterexamples not only for (1.30) with  $\alpha \cdot p$  but also for (1.59) with  $\beta \cdot p$ , though the construction for both is only slightly different. To the latter, as a matter of fact, we will come back in Section 6 to make some important comments on the semi-norms concerned.

In [BEU], they observed, for the 3-dimensional Weyl–Dirac (or Pauli) operator  $\sigma \cdot p$ , that, for  $1 < p < 3$  with  $q = \frac{3p}{3-p}$ , the following inequality:

$$\|h\|_q \leq C(p) \|(\sigma \cdot p)h\|_p \quad (3.1)$$

holds for all  $h \in C_0^\infty(\mathbb{R}^3; \mathbb{C}^2)$  with a positive constant  $C(p)$  depending on  $p$ . This is a consequence from the usual Sobolev inequality together with the fact that, for  $1 < p < \infty$ , the two semi-norms  $\|(\sigma \cdot p)h\|_p$  and  $\|\nabla h\|_p$  are equivalent (cf. [IS] and Lemma 3.2 of the present paper where analogous results are given for the Dirac operator  $\alpha \cdot p$  instead of Weyl–Dirac (or Pauli)  $\sigma \cdot p$ ). They showed also that (3.1) is untrue when  $p = 1$ , by using a zero mode for an appropriate Weyl–Dirac (or Pauli) operator constructed by Loss–Yau [LoY] to make a sequence  $\{h_n\} \subset C_0^\infty(\mathbb{R}^3; \mathbb{C}^2)$  such that  $\{\|(\sigma \cdot p)h_n\|_1\}$  is uniformly bounded for all over  $n$ , but that  $\|h_n\|_{3/2} \geq (\text{positive constant}) \cdot (\log n)^{2/3}$ , concluding invalidity of (3.1) for  $p = 1$ . As a result, this sequence will turn out to violate (1.35) in Corollary 1.2.

We will modify their argument so as to apply to our cases of Theorems 1.1 and 1.6 to construct an example. First we consider the case for three-dimensional massless Dirac operator  $\alpha \cdot p$  and next for 4-dimensional Euclidian Dirac operator  $\beta \cdot p$ .

*An example for (1.30) of Theorem 1.1 with  $p = 1$  to fail to hold.*

So with  $x \in \mathbb{R}^3$  and  $|x| = (x_1^2 + x_2^2 + x_3^2)^{1/2}$ , let

$$\begin{aligned} e(x) &:= \frac{1}{(1 + |x|^2)^{3/2}} (I_4 + i\alpha \cdot x) \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \\ &= \frac{1}{(1 + |x|^2)^{3/2}} \begin{pmatrix} 1 & 0 & ix_3 & ix_1 + x_2 \\ 0 & 1 & ix_1 - x_2 & -ix_3 \\ ix_3 & ix_1 + x_2 & 1 & 0 \\ ix_1 - x_2 & -ix_3 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \\ &= \frac{1}{(1 + |x|^2)^{3/2}} \begin{pmatrix} 1 \\ 0 \\ ix_3 \\ ix_1 - x_2 \end{pmatrix}, \end{aligned} \quad (3.2)$$

where  $I_4$  is the  $4 \times 4$ -identity matrix. Then we can see  $e(x)$  satisfies the following equation

$$(\alpha \cdot p)e(x) = \frac{3}{1 + |x|^2}e(x), \quad (3.3)$$

and inequalities:

$$\begin{aligned} |e(x)|_{\ell^\infty} &= \frac{1 \vee |ix_3| \vee |ix_1 - x_2|}{(1 + |x|^2)^{3/2}} = \frac{1 \vee |x_3| \vee (x_1^2 + x_2^2)^{1/2}}{(1 + |x|^2)^{3/2}} \\ &\leq \frac{(1 + x_1^2 + x_2^2 + x_3^2)^{1/2}}{(1 + |x|^2)^{3/2}} = \frac{1}{1 + |x|^2}, \end{aligned} \quad (3.4)$$

$$\begin{aligned} |e(x)|_{\ell^q}^q &= \frac{1 + |ix_3|^q + |ix_1 - x_2|^q}{(1 + |x|^2)^{3q/2}} = \frac{1 + (x_3^2)^{q/2} + (x_1^2 + x_2^2)^{q/2}}{(1 + |x|^2)^{3q/2}} \\ &\geq \frac{(1 + |x|^2)^{q/2}}{(1 + |x|^2)^{3q/2}} = \left( \frac{1}{1 + |x|^2} \right)^q \quad (1 \leq q \leq 2), \end{aligned} \quad (3.5)$$

where (3.5) is due to that  $a^{q/2} + b^{q/2} \geq (a + b)^{q/2}$  for  $a \geq 0$ ,  $b \geq 0$  and  $1 \leq q \leq 2$ .

For each positive integer  $n$ , put  $f_n(x) = \rho_n(|x|)e(x)$ , where  $\rho_n(r)$  is a nonnegative cutoff function in  $C_0^\infty(\mathbb{R})$  such that  $\rho_n(r) = 1$  ( $r \leq n$ );  $= 0$  ( $r \geq n + 2$ ), and further  $|\rho'_n(r)| \equiv |(d/dr)\rho_n(r)| \leq 1$  for all  $r \geq 0$ . Then it is evident that  $f_n$  belongs to  $C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$ .

We are going to see that inequality (1.16)/(1.30) does not hold with any constant  $C > 0$  for  $p = 1$ ,  $q = \frac{3}{2}$  and hence  $\frac{p}{q} = \frac{2}{3}$ . Indeed, there exists no constant  $C$  such that, for all  $n$ ,

$$\|f_n\|_{3/2} \leq C \|(\alpha \cdot p)f_n\|_1^{2/3} \|f_n\|_{B_{\infty,\infty}^{-2}}^{1/3}. \quad (3.6)$$

First, we show that the sequence  $\{(\alpha \cdot p)f_n\}_{n=1}^\infty$  is uniformly bounded in  $L^1$ . Indeed, since

$$\begin{aligned} (\alpha \cdot p)f_n(x) &= \rho_n(|x|)(\alpha \cdot p)e(x) + ((\alpha \cdot p)\rho_n(|x|))e(x) \\ &= \rho_n(|x|) \frac{3}{1 + |x|^2}e(x) - i\rho'_n(|x|) \frac{\alpha \cdot x}{|x|}e(x) \\ &= \frac{3\rho_n(|x|)}{(1 + |x|^2)^{5/2}} \begin{pmatrix} 1 \\ 0 \\ ix_3 \\ ix_1 - x_2 \end{pmatrix} + \frac{\rho'_n(|x|)}{|x|(1 + |x|^2)^{3/2}} \begin{pmatrix} |x|^2 \\ 0 \\ -ix_3 \\ -ix_1 + x_2 \end{pmatrix}, \end{aligned}$$

we can estimate the  $L^1$  norm of  $(\alpha \cdot p)f_n$ , noting  $\rho'_n(|x|) = 0$  for  $|x| \leq n$  and  $|x| \geq n + 2$

and using polar coordinates, to get

$$\begin{aligned}
\|(\alpha \cdot p)f_n\|_1 &\leq \int_{\{x \in \mathbb{R}^3; |x| \leq n+2\}} \frac{3(1 + |ix_3| + |ix_1 - x_2|)}{(1 + |x|^2)^{5/2}} dx \\
&\quad + \int_{\{x \in \mathbb{R}^3; n \leq |x| \leq n+2\}} \frac{|x|^2 + |-ix_3| + |-ix_1 + x_2|}{|x|(1 + |x|^2)^{3/2}} dx \\
&\leq \int_{|x| \leq n+2} \frac{3(\sqrt{2}|x| + 1)}{(1 + |x|^2)^{5/2}} dx + \int_{n \leq |x| \leq n+2} \frac{|x|^2 + \sqrt{2}|x|}{|x|(1 + |x|^2)^{3/2}} dx \\
&= \int_0^{n+2} \frac{3(\sqrt{2}r + 1)4\pi r^2 dr}{(1 + r^2)^{5/2}} + \int_n^{n+2} \frac{(r + \sqrt{2})4\pi r^2 dr}{(1 + r^2)^{3/2}} \\
&\leq 12\pi \int_0^{n+2} \frac{2}{1 + r^2} dr + 4\pi \int_n^{n+2} 2dr \\
&= 24\pi \tan^{-1}(n+2) + 16\pi \leq 24\pi \cdot \frac{\pi}{2} + 16\pi, \tag{3.7}
\end{aligned}$$

where we have used that  $(r + \sqrt{2})r^2 \leq 2(1 + r^2)^{3/2}$  and  $(\sqrt{2}r + 1)r^2 \leq 2(1 + r^2)^{3/2}$  for all  $r \geq 0$ . Thus we have shown the sequence  $\{\|(\alpha \cdot p)f_n\|_1\}$  is uniformly bounded.

Next, we study how  $\{f_n\}_{n=1}^\infty$  behaves in the norm of  $B_{\infty,\infty}^{-2}(\mathbb{R}^3; \mathbb{C}^4)$  for large  $n$ . In fact, we shall show

$$\|f_n\|_{B_{\infty,\infty}^{-2}} = O(\log n). \tag{3.8}$$

Here note that  $\frac{p}{p-q} = -\frac{1}{\frac{3}{2}-1} = -2$ . Indeed, we have with (3.4)

$$\begin{aligned}
\|f_n\|_{B_{\infty,\infty}^{-2}} &= \sup_{t>0} t \|P_t f_n\|_\infty = \sup_{t>0} t \sup_x \int \frac{1}{(4\pi t)^{3/2}} e^{-\frac{|x-y|^2}{4t}} \rho_n(|y|) |e(y)|_{\ell^\infty} dy \\
&\leq \frac{2}{(4\pi)^{3/2}} \sup_{t>0} \sup_x \int \left( \frac{|x-y|^2}{4t} \right)^{1/2} e^{-\frac{|x-y|^2}{4t}} \frac{\rho_n(|y|)}{|x-y|(1 + |y|^2)} dy \\
&\leq \frac{2}{(4\pi)^{3/2}} (2e)^{-1/2} \sup_x \int_{|y| \leq n+2} \frac{1}{|x-y|(1 + |y|^2)} dy,
\end{aligned}$$

where the last inequality is due to the fact that  $s^{1/2}e^{-s} \leq (2e)^{-1/2}$  for all  $s > 0$ . Then we use polar coordinates to get

$$\begin{aligned}
\|f_n\|_{B_{\infty,\infty}^{-2}} &\leq \frac{1}{(4\pi)^{3/2}} \left(\frac{2}{e}\right)^{1/2} \sup_x \int_0^{n+2} \frac{r^2}{1 + r^2} dr \int_0^\pi \frac{2\pi \sin \theta d\theta}{(|x|^2 + r^2 - 2|x|r \cos \theta)^{1/2}} \\
&= \frac{2\pi}{(4\pi)^{3/2}} \left(\frac{2}{e}\right)^{1/2} \sup_x \int_0^{n+2} \frac{r^2 dr}{1 + r^2} \left[ \frac{(|x|^2 + r^2 - 2|x|r \cos \theta)^{1/2}}{|x|r} \right]_{\theta=0}^{\theta=\pi} \\
&= \frac{1}{2(2\pi e)^{1/2}} \sup_x \frac{1}{|x|} \int_0^{n+2} \frac{r[ (|x| + r) - |x| - r ]}{1 + r^2} dr \\
&= \left[ \sup_{|x| \geq n+2} \vee \sup_{|x| \leq n+2} \right] \frac{1}{2(2\pi e)^{1/2}} \frac{1}{|x|} \int_0^{n+2} \frac{r[ (|x| + r) - |x| - r ]}{1 + r^2} dr \\
&=: V_{\alpha,1} \vee V_{\alpha,2}.
\end{aligned}$$

Then we can conclude (3.8) above, noting

$$\begin{aligned}
2(2\pi e)^{1/2}V_{\alpha,1} &= \sup_{|x| \geq n+2} \frac{1}{|x|} \int_0^{n+2} \frac{2r^2}{1+r^2} dr \leq 2, \\
2(2\pi e)^{1/2}V_{\alpha,2} &= \sup_{|x| \leq n+2} \frac{1}{|x|} \left[ \int_0^{|x|} \frac{2r^2}{1+r^2} dr + \int_{|x|}^{n+2} \frac{2|x|r}{1+r^2} dr \right] \\
&\leq 2 + \log(1 + (n+2)^2) = O(\log n).
\end{aligned}$$

Thus, by (3.8) and since, as already seen above, the sequence  $\{ \|(\alpha \cdot p)f_n\|_1 \}$  is uniformly bounded, we see the sequence  $\{ \|(\alpha \cdot p)f_n\|_1^{2/3} \|f_n\|_{\mathbb{B}_{\infty,\infty}^{-2}}^{1/3} \}$  on the right-hand side of (3.6) is of order  $O((\log n)^{1/3})$ , while, for the left-hand side, we have by (3.5) with  $q = \frac{3}{2}$

$$\begin{aligned}
\|f_n\|_{3/2} &\geq \left( \int_{|x| \leq n} \frac{1}{(1+|x|^2)^{3/2}} dx \right)^{2/3} = \left( \int_0^n \frac{4\pi r^2}{(1+r^2)^{3/2}} dr \right)^{2/3} \\
&\geq \left( \int_1^n \frac{4\pi r^2}{(1+r^2)^{3/2}} dr \right)^{2/3} \geq \left( \int_1^n \frac{4\pi}{r} dr \right)^{2/3} \geq (4\pi)^{2/3} (\log n)^{2/3}. \quad (3.9)
\end{aligned}$$

This means that inequality (3.6) or (1.16)/(1.30) with  $p = 1$ ,  $q = \frac{3}{2}$  does not hold.

*An example for (1.59) of Theorem 1.6 with  $p = 1$  to fail to hold.*

This case is with  $x \in \mathbb{R}^4$  and  $|x| = (x_1^2 + x_2^2 + x_3^2 + x_4^2)^{1/2}$ . We can use the same arguments as above to construct a sequence  $\{f_n\}$  in  $C_0^\infty(\mathbb{R}^4; \mathbb{C}^4)$  such that (1.59) fails to hold for any fixed constant  $C$ , starting, instead of (3.2), from the following function

$$\begin{aligned}
\hat{e}(x) &:= \frac{1}{(1+|x|^2)^2} (I_4 + i\beta \cdot x) \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \\
&= \frac{1}{(1+|x|^2)^2} \begin{pmatrix} 1 & 0 & ix_3 + x_4 & ix_1 + x_2 \\ 0 & 1 & ix_1 - x_2 & -ix_3 + x_4 \\ ix_3 - x_4 & ix_1 + x_2 & 1 & 0 \\ ix_1 - x_2 & -ix_3 - x_4 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \\
&= \frac{1}{(1+|x|^2)^2} \begin{pmatrix} 1 \\ 0 \\ ix_3 - x_4 \\ ix_1 - x_2 \end{pmatrix}. \quad (3.10)
\end{aligned}$$

It can be seen that  $\hat{e}(x)$  satisfies the following equation

$$(\beta \cdot p)\hat{e}(x) = \frac{4}{1+|x|^2} \hat{e}(x), \quad (3.11)$$

and inequalities:

$$|\hat{e}(x)|_{\ell^\infty} = \frac{1 \vee |ix_3 - x_4| \vee |ix_1 - x_2|}{(1+|x|^2)^2} \leq \frac{1}{(1+|x|^2)^{3/2}}, \quad (3.12)$$

$$\begin{aligned}
|\hat{e}(x)|_{\ell^q}^q &= \frac{1 + |ix_3 - x_4|^q + |ix_1 - x_2|^q}{(1+|x|^2)^{2q}} = \frac{1 + (x_3^2 + x_4^2)^{q/2} + (x_1^2 + x_2^2)^{q/2}}{(1+|x|^2)^{2q}} \\
&\geq \frac{(1+|x|^2)^{q/2}}{(1+|x|^2)^{2q}} = \left( \frac{1}{1+|x|^2} \right)^{3q/2} \quad (1 \leq q \leq 2). \quad (3.13)
\end{aligned}$$

For each positive integer  $n$ , put  $f_n(x) = \rho_n(|x|)\hat{e}(x)$ , where  $\rho_n(r)$  is the same nonnegative cutoff function in  $C_0^\infty(\mathbb{R})$  as before such that  $\rho_n(r) = 1$  ( $r \leq n$ );  $= 0$  ( $r \geq n+2$ ), and further  $|\rho_n'(r)| \equiv |(d/dr)\rho_n(r)| \leq 1$  for all  $r \geq 0$ . Then it is evident that  $f_n$  belongs to  $C_0^\infty(\mathbb{R}^4; \mathbb{C}^4)$ .

We are going to see that inequality (1.59), corresponding to (1.16) in the case for  $\beta \cdot p$ , does not hold with any constant  $C > 0$  for  $p = 1$ ,  $q = \frac{4}{3}$  and hence  $\frac{p}{q} = \frac{3}{4}$ . Indeed, there exists no constant  $C$  such that, for all  $n$ ,

$$\|f_n\|_{4/3} \leq C \|(\alpha \cdot p)f_n\|_1^{3/4} \|f_n\|_{B_{\infty,\infty}^{-3}}^{1/4}. \quad (3.14)$$

First, we show that the sequence  $\{(\beta \cdot p)f_n\}_{n=1}^\infty$  is uniformly bounded in  $L^1$ . Indeed, since

$$\begin{aligned} (\beta \cdot p)f_n(x) &= \rho_n(|x|)(\beta \cdot p)\hat{e}(x) + ((\beta \cdot p)\rho_n(|x|))\hat{e}(x) \\ &= \rho_n(|x|) \frac{4}{1+|x|^2} \hat{e}(x) - i\rho_n'(|x|) \frac{\beta \cdot x}{|x|} \hat{e}(x) \\ &= \frac{4\rho_n(|x|)}{(1+|x|^2)^3} \begin{pmatrix} 1 \\ 0 \\ ix_3 - x_4 \\ ix_1 - x_2 \end{pmatrix} + \frac{\rho_n'(|x|)}{|x|(1+|x|^2)^2} \begin{pmatrix} |x|^2 \\ 0 \\ -ix_3 + x_4 \\ -ix_1 + x_2 \end{pmatrix}, \end{aligned}$$

we can estimate the  $L^1$  norm of  $(\beta \cdot p)f_n$ , noting  $\rho_n'(|x|) = 0$  for  $|x| \leq n$  and  $|x| \geq n+2$  and using polar coordinates, to get

$$\begin{aligned} \|(\beta \cdot p)f_n\|_1 &\leq \int_{\{x \in \mathbb{R}^4; |x| \leq n+2\}} \frac{4(1+|ix_3-x_4|+|ix_1-x_2|)}{(1+|x|^2)^3} dx \\ &\quad + \int_{\{x \in \mathbb{R}^4; n \leq |x| \leq n+2\}} \frac{|x|^2 + |-ix_3+x_4| + |-ix_1+x_2|}{|x|(1+|x|^2)^2} dx \\ &\leq \int_{|x| \leq n+2} \frac{4(1+\sqrt{2}|x|)}{(1+|x|^2)^3} dx + \int_{n \leq |x| \leq n+2} \frac{|x|^2 + \sqrt{2}|x|}{|x|(1+|x|^2)^2} dx \\ &= \int_0^{n+2} \frac{4(1+\sqrt{2}r)2\pi^2 r^3 dr}{(1+r^2)^3} + \int_n^{n+2} \frac{(r+\sqrt{2})2\pi^2 r^3 dr}{(1+r^2)^2} \\ &\leq 8\pi^2 \int_0^{n+2} \frac{2}{1+r^2} dr + 2\pi^2 \int_n^{n+2} 2 dr \\ &= 16\pi^2 \tan^{-1}(n+2) + 8\pi^2 \leq 16\pi^2 \cdot \frac{\pi}{2} + 8\pi^2, \end{aligned} \quad (3.15)$$

where in the second inequality we have used that  $(r+\sqrt{2})r^3 \leq 2(1+r^2)^2$  and  $(1+\sqrt{2}r)r^3 \leq 2(1+r^2)^2$  for all  $r \geq 0$ . Thus we have shown the sequence  $\{(\alpha \cdot p)f_n\}_1$  is uniformly bounded.

Next, we study how  $\{f_n\}_{n=1}^\infty$  behaves in the norm of  $B_{\infty,\infty}^{-3}(\mathbb{R}^4; \mathbb{C}^4)$  for large  $n$ . In fact, we show

$$\|f_n\|_{B_{\infty,\infty}^{-3}} = O(\log n). \quad (3.16)$$

Here we note that  $\frac{p}{p-q} = -\frac{1}{\frac{4}{3}-1} = -3$ . Indeed, we have with (3.12)

$$\begin{aligned}
\|f_n\|_{B_{\infty,\infty}^{-3}} &= \sup_{t>0} t^{3/2} \|P_t f_n\|_{\infty} = \sup_{t>0} t^{3/2} \sup_x \int_{\mathbb{R}^4} \frac{1}{(4\pi t)^2} e^{-\frac{|x-y|^2}{4t}} \rho_n(|y|) |\hat{e}(y)|_{\ell^\infty} dy \\
&\leq \frac{2}{(4\pi)^2} \sup_{t>0} \sup_x \int \left( \frac{|x-y|^2}{4t} \right)^{1/2} e^{-\frac{|x-y|^2}{4t}} \frac{\rho_n(|y|)}{|x-y|(1+|y|^2)^{3/2}} dy \\
&\leq \frac{2}{(4\pi)^2} (2e)^{-1/2} \sup_x \int_{|y|\leq n+2} \frac{1}{|x-y|(1+|y|^2)^{3/2}} dy,
\end{aligned}$$

where the last inequality is due to the fact that  $s^{1/2}e^{-s} \leq (2e)^{-1/2}$  for all  $s > 0$ . Then we use polar coordinates and  $\sin^2\theta \leq \sin\theta$  ( $0 \leq \theta \leq \pi$ ) to get

$$\begin{aligned}
\|f_n\|_{\mathbb{B}_{\infty,\infty}^{-3}} &\leq \frac{2}{(4\pi)^2} (2e)^{-1/2} \sup_x \int_0^{n+2} \frac{r^3}{(1+r^2)^{3/2}} dr \int_0^\pi \frac{4\pi \sin^2\theta d\theta}{(|x|^2 + r^2 - 2|x|r\cos\theta)^{1/2}} \\
&\leq \frac{2}{(4\pi)^2} (2e)^{-1/2} \sup_x \int_0^{n+2} \frac{r^3}{(1+r^2)^{3/2}} dr \int_0^\pi \frac{4\pi \sin\theta d\theta}{(|x|^2 + r^2 - 2|x|r\cos\theta)^{1/2}} \\
&= \frac{4\pi}{(4\pi)^2} (2e)^{-1/2} \sup_x \int_0^{n+2} \frac{r^3}{(1+r^2)^{3/2}} dr \left[ \frac{(|x|^2 + r^2 - 2|x|r\cos\theta)^{1/2}}{|x|r} \right]_{\theta=0}^{\theta=\pi} \\
&= \frac{1}{4\pi(2e)^{1/2}} \sup_x \frac{1}{|x|} \int_0^{n+2} \frac{r^2((|x|+r) - ||x|-r|)}{(1+r^2)^{3/2}} dr \\
&= \left[ \sup_{|x|\geq n+2} \vee \sup_{|x|\leq n+2} \right] \frac{1}{4\pi(2e)^{1/2}} \frac{1}{|x|} \int_0^{n+2} \frac{r^2((|x|+r) - ||x|-r|)}{(1+r^2)^{3/2}} dr \\
&=: V_{\beta,1} + V_{\beta,2}.
\end{aligned}$$

Then we can conclude (3.16), noting

$$\begin{aligned}
4\pi(2e)^{1/2}V_{\beta,1} &= \sup_{|x|\geq n+2} \frac{1}{|x|} \int_0^{n+2} \frac{2r^3}{(1+r^2)^{3/2}} dr \leq 2, \\
4\pi(2e)^{1/2}V_{\beta,2} &= \sup_{|x|\leq n+2} \frac{1}{|x|} \left\{ \int_0^{|x|} \frac{2r^3}{(1+r^2)^{3/2}} dr + \int_{|x|}^{n+2} \frac{2|x|r^2}{(1+r^2)^{3/2}} dr \right\} \\
&\leq 2 + \log[(n+2) + (1+(n+2)^2)^{1/2}] = O(\log n).
\end{aligned}$$

Thus, by (3.16) and since, as already seen above, the sequence  $\{\|(\beta \cdot p)f_n\|_1\}$  is uniformly bounded, we see the sequence  $\{\|(\beta \cdot p)f_n\|_1^{3/4} \|f_n\|_{\mathbb{B}_{\infty,\infty}^{-3}}^{1/4}\}$  on the right-hand side of (3.14) is of order  $O((\log n)^{1/4})$ , while, for the left-hand side, we have by (3.13) with  $q = \frac{4}{3}$

$$\begin{aligned}
\|f_n\|_{4/3} &\geq \left( \int_{|x|\leq n} |\hat{e}(x)|_{\ell^{4/3}}^{4/3} dx \right)^{3/4} \geq \left( \int_{|x|\leq n} \left( \frac{1}{1+|x|^2} \right)^{(3/2)\cdot(4/3)} dx \right)^{3/4} \\
&= \left( \int_0^n \frac{2\pi^2 r^3 dr}{(1+r^2)^2} \right)^{3/4} = O((\log n)^{3/4})
\end{aligned} \tag{3.17}$$

for large  $n$ . This means that inequality (3.14) or (1.59) with  $p = 1$ ,  $q = \frac{4}{3}$  does not hold.

## 4 Proof of Theorem 1.1

*Proof of Theorem 1.1 (i).* We follow the lucid arguments used in Ledoux [Le]. The proof is divided into three steps. In step I, we mention the weak-type inequality (1.14) given by [BES] with the idea of [Le] to sketch its proof, for the paper to be somehow self-contained. In step II we show the inequality (1.28) in the special case under the condition  $f \in L^q(\mathbf{R}^3; \mathbf{C}^4)$  and then the general case in step III.

I. So we begin with a sketch of proof of inequality (1.14).

To do so, assume that  $f$  satisfies  $M_{\alpha \cdot p}(f) < \infty$ . Note that this implies with (1.22) that  $\|(\alpha \cdot p)f\|_p < \infty$ . And further assume that our  $f$  satisfies  $\|f\|_{B_{\infty, \infty}^{p/(p-q)}} < \infty$ . We may suppose by our convention (1.7) of notations and by homogeneity that

$$\|f\|_{B_{\infty, \infty}^{p/(p-q)}} = \sup_{t>0} t^{-p/2(p-q)} \|P_t f\|_{\infty} \leq 1. \quad (4.1)$$

Therefore  $|P_t f|_{\ell^\infty} \leq t^{p/2(p-q)}$  pointwise. For  $u > 0$ , put  $t = t_u \equiv u^{2(p-q)/p}$ , so that  $|P_{t_u} f|_{\ell^\infty} \leq u$ . Hence that  $|f|_{\ell^\infty} \geq 2u$  pointwise implies that  $|f - P_{t_u} f|_{\ell^\infty} \geq |f|_{\ell^\infty} - |P_{t_u} f|_{\ell^\infty} \geq u$  pointwise. Then

$$\begin{aligned} u^q |\{|f|_{\ell^\infty} \geq 2u\}| &\leq u^q |\{|f - P_{t_u} f|_{\ell^\infty} \geq u\}| \\ &\leq u^q \int \frac{|f - P_{t_u} f|_{\ell^\infty}^p}{u^p} dx = u^q \int \bigvee_{k=1}^4 \frac{|f_k - e^{t_u \Delta} f_k|^p}{u^p} dx \\ &\leq u^{q-p} \int \sum_{k=1}^4 |f_k - e^{t_u \Delta} f_k|^p dx \\ &= u^{q-p} \int |f - P_{t_u} f|_{\ell^p}^p dx = u^{q-p} \|f - P_{t_u} f\|_p^p. \end{aligned}$$

In [BES], it is shown that

$$\|f - P_{t_u} f\|_p \leq c_0 t_u^{1/2} \|(\alpha \cdot p)f\|_p. \quad (4.2)$$

with a positive constant  $c_0$  depending only on  $p$ . Then by (4.2) and since  $q - p + p(p - q)/p = 0$ , we have

$$u^q |\{|f|_{\ell^\infty} \geq 2u\}| \leq c_0 u^{q-p} t_u^{p/2} \int |(\alpha \cdot p)f|_{\ell^p}^p dx = c_0 \int |(\alpha \cdot p)f|_{\ell^p}^p dx.$$

This yields the weak type inequality (1.14), taking account of definition of  $\|f\|_{q, \infty}$  in (1.15).

II. Next we want to replace the weak  $L^q$  norm on the left-hand side of (1.14) by the strong  $L^q$  norm. Here we note with (1.22) that (1.14) holds also with  $M_{\alpha \cdot p}(f)$  in place of  $\|(\alpha \cdot p)f\|_p$ . We show inequality (1.28) for  $f$  which satisfies  $M_{\alpha \cdot p}(f) < \infty$  and (4.1), i.e.  $\|f\|_{B_{\infty, \infty}^{p/(p-q)}} \leq 1$ , as in step I, and the extra condition  $f \in L^q(\mathbf{R}^3; \mathbf{C}^4)$ . In step III below, we shall remove this latter condition.

Then what we need to show is that there exists a constant  $C$  (depending only on  $q$  and  $p$ ) such that

$$\int |f|_{\ell^q}^q dx \leq C M_{\alpha \cdot p}(f)^p, \quad (4.3)$$



which amounts to our goal inequality (1.28), if only  $f$  replaced by  $f/\|f\|_{B_{\infty,\infty}^{p/(p-q)}}$  in (4.3).

Now, for  $u > 0$ , let  $t = t_u = u^{2(p-q)/p}$  again. Let  $c \geq 5$  (depending on  $q$  and  $p$ ) to be specified later.

Note the ‘layer cake’ representation [LLo, p.26, Theorem 1.13] for any nonnegative measurable function  $\psi(x)$ :

$$\psi(x) = \int_0^\infty \chi_{\{\psi > s\}}(x) ds. \quad (4.4)$$

In particular, we have

$$|f(x)|_{\ell^q}^q = \int_0^\infty \chi_{\{|f|_{\ell^q}^q > s\}}(x) ds = \int_0^\infty \chi_{\{|f|_{\ell^q} > u\}}(x) d(u^q),$$

so that by Fubini’s theorem

$$\begin{aligned} \frac{1}{20^q} \|f\|_q^q &= \frac{1}{20^q} \int |f(x)|_{\ell^q}^q dx = \frac{1}{20^q} \int dx \int_0^\infty \chi_{\{|f|_{\ell^q} > u\}}(x) d(u^q) \\ &= \frac{1}{20^q} \int_0^\infty d(u^q) \int \chi_{\{|f|_{\ell^q} \geq u\}}(x) dx = \int_0^\infty |\{|f|_{\ell^q} \geq 20u\}| d(u^q). \end{aligned} \quad (4.5)$$

For every  $u > 0$  and for  $f(x) = {}^t(f_1(x), f_2(x), f_3(x), f_4(x))$ , let

$$\begin{aligned} f_u(x) &= {}^t(f_{u,1}(x), f_{u,2}(x), f_{u,3}(x), f_{u,4}(x)), \\ f_{u,k}(x) &:= (f_k(x) - u)^+ \wedge ((c-1)u) + (f_k(x) + u)^- \vee (-(c-1)u), \quad k = 1, 2, 3, 4, \end{aligned} \quad (4.6)$$

for any  $c > 1$ . Here, as in (1.7),  $a \vee b$  denotes  $\max\{a, b\}$ , while  $a \wedge b$  denotes  $\min\{a, b\}$ .

Notice that  $f_u$  also satisfies the same condition as  $f$ . Each  $f_{u,k}(x)$  satisfies  $0 \leq |f_{u,k}(x)| \leq (c-1)u$ . It vanishes when  $|f_k(x)| \leq u$  and is equal to  $(c-1)u$  when  $f_k(x) \geq cu$ , and to  $-(c-1)u$  when  $f_k(x) \leq -cu$ .

We see that, since on the set  $\{|f_k| \geq 5u\}$ , we have  $|f_{u,k}| \geq 4u$  for each fixed  $k$ , and that on the set  $\{|f|_{\ell^\infty} \geq 5u\}$ , we have  $|f_u|_{\ell^\infty} \geq 4u$ . We have

$$|f_{u,k}| \leq |f_{u,k} - e^{t_u \Delta} f_{u,k}| + e^{t_u \Delta} |f_{u,k} - f_k| + |e^{t_u \Delta} f_k|, \quad k = 1, 2, 3, 4. \quad (4.7)$$

By noting the notation (1.7) of the  $\ell^p/\ell^\infty$  norm of a four-vector we have

$$\begin{aligned} \int_0^\infty |\{|f|_{\ell^q} \geq 20u\}| d(u^q) &\leq \int_0^\infty |\{|f|_{\ell^\infty} \geq 5u\}| d(u^q) \\ &\leq \int_0^\infty |\{|f_u|_\infty \geq 4u\}| d(u^q) = \int_0^\infty |\{\vee_{k=1}^4 |f_{u,k}| \geq 4u\}| d(u^q) \\ &\leq \int_0^\infty |\{\vee_{k=1}^4 |f_{u,k} - e^{t_u \Delta} f_{u,k}| \geq u\}| d(u^q) \\ &\quad + \int_0^\infty |\{\vee_{k=1}^4 e^{t_u \Delta} |f_{u,k} - f_k| \geq 2u\}| d(u^q) \\ &=: J_1 + J_2, \end{aligned} \quad (4.8)$$

where we have used the fact that  $|P_{t_u}(f)|_{\ell^\infty} \leq u$ , which holds by our choice of  $f$  in (4.1).

We shall estimate the last member  $J_1 + J_2$  of (4.8). First, to treat the second term  $J_2$ , we confirm that

$$|f_{u,k} - f_k| = |f_{u,k} - f_k| \chi_{\{|f_k| \leq cu\}} + |f_{u,k} - f_k| \chi_{\{|f_k| > cu\}} \leq u + |f_k| \chi_{\{|f_k| > cu\}}. \quad (4.9)$$

This is checked with (4.6) as follows. Indeed, we see (4.6) imply that

$$f_{u,k}(x) - f_k(x) = \begin{cases} (-u) \wedge (-f_k(x) + (c-1)u), & \text{if } f_k(x) \geq u, \\ u \vee (-f_k(x) - (c-1)u), & \text{if } f_k(x) \leq -u. \end{cases}$$

This further implies on the one hand that

$$f_{u,k}(x) - f_k(x) = \begin{cases} -u, & \text{if } u \leq f_k(x) \leq cu, \\ u, & \text{if } -u \geq f_k(x) \geq -cu, \end{cases}$$

so that  $|f_{u,k}(x) - f_k(x)| = u$ , if  $u \leq |f_k(x)| \leq cu$ , and on the other hand that

$$f_{u,k}(x) - f_k(x) = \begin{cases} -f_k(x) + (c-1)u \geq -f_k(x), & \text{if } f_k(x) \geq cu, \\ -f_k(x) - (c-1)u \leq -f_k(x), & \text{if } f_k(x) \leq -cu, \end{cases}$$

so that  $|f_{u,k}(x) - f_k(x)| \leq |f_k(x)|$ , if  $|f_k(x)| \geq cu$ . This yields (4.9).

Then, since  $e^{tu\Delta}$  is positivity-preserving, it follows that

$$\begin{aligned} J_2 &= \int_0^\infty |\{\vee_{k=1}^4 e^{tu\Delta} |f_{u,k} - f_k| \geq 2u\}| d(u^q) \\ &\leq \int_0^\infty |\{\vee_{k=1}^4 e^{tu\Delta} |f_k| \chi_{\{|f_k| > cu\}} \geq u\}| d(u^q) \\ &\leq \int_0^\infty \left( \int \vee_{k=1}^4 \frac{e^{tu\Delta} |f_k| \chi_{\{|f_k| > cu\}}}{u} dx \right) d(u^q) \\ &= \int_0^\infty \frac{1}{u} \left( \int \vee_{k=1}^4 |f_k| \chi_{\{|f_k| > cu\}} dx \right) d(u^q) \leq \int_0^\infty \frac{1}{u} \left( \int |f(x)|_{\ell^\infty} \chi_{\{|f|_{\ell^\infty} > cu\}} dx \right) d(u^q) \\ &\leq \frac{q}{q-1} \int |f|_{\ell^q} \left( \int_0^\infty \chi_{\{|f|_{\ell^q} > cu\}} d(u^{q-1}) \right) dx = \frac{q}{q-1} \frac{1}{c^{q-1}} \|f\|_q^q. \end{aligned} \quad (4.10)$$

Here the last fourth equality is due to that

$$\begin{aligned} \int e^{tu\Delta} |f_k| \chi_{\{|f_k| > cu\}} dx &= \int \left( \int (e^{tu\Delta}(x-y) |f_k(y)| \chi_{\{|f_k(y)| > cu\}} dy) \right) dx \\ &= \int |f_k(y)| \chi_{\{|f_k(y)| > cu\}} dy, \end{aligned}$$

because the heat kernel  $e^{tu\Delta}(x-y)$  satisfies  $\int e^{tu\Delta}(x) dx = 1$  for  $t_u > 0$ , and the last second inequality is due to that  $\vee_{k=1}^4 |f_k(x)| \leq |f(x)|_{\ell^\infty} \leq |f(x)|_{\ell^q}$  by (1.7).

Next, as for the first term  $J_1$  of the last member of (4.8), we have by (4.2)

$$\begin{aligned} |\{\vee_{k=1}^4 |f_{u,k} - e^{tu\Delta} f_{u,k}| \geq u\}| &\leq \int \vee_{k=1}^4 \frac{|f_{u,k} - e^{tu\Delta} f_{u,k}|^p}{u^p} dx \\ &\leq u^{-p} \int \sum_{k=1}^4 |f_{u,k} - e^{tu\Delta} f_{u,k}|^p dx \\ &= u^{-p} \int |f_u - P_{t_u}(f_u)|_{\ell^p}^p dx \\ &\leq c_0 u^{-p} t_u^{p/2} \int |(\alpha \cdot p) f_u|_{\ell^p}^p dx = c_0 u^{-q} \|(\alpha \cdot p) f_u\|_p^p \\ &\leq C_0 u^{-q} M_{\alpha \cdot p; p}(f_u)^p, \end{aligned}$$

with  $C_0 := 2^{1-(1/p)}c_0$ , where the last inequality is due to (1.22), so that

$$J_1 \leq C_0 \int_0^\infty d(u^q)u^{-q}M_{\alpha;p;p}(f_u)^p. \quad (4.11)$$

For (4.11), we want to show the following lemma.

**Lemma 4.1.** *Let  $f = {}^t(f_1, f_2, f_3, f_4)$  satisfy  $M_{\alpha;p;p}(f) < \infty$  and  $\|f\|_{B_{\infty,\infty}^{p/(p-q)}} \leq 1$ . Let  $f_u = {}^t(f_{u,1}, f_{u,2}, f_{u,3}, f_{u,4})$  as in (4.6). Then*

$$\int_0^\infty d(u^q)u^{-q}M_{\alpha;p;p}(f_u)^p = q(\log c)M_{\alpha;p;p}(f)^p. \quad (4.12)$$

*Proof.* For  $f_u$  in (4.6) instead of  $f$ , we have by (1.21)

$$\begin{aligned} & M_{\alpha;p;p}(f_u)^p \\ &= \int (|(\partial_1 + i\partial_2)f_{u,1}|^p + |\partial_3 f_{u,1}|^p)dx + \int (|(\partial_1 - i\partial_2)f_{u,2}|^p + |\partial_3 f_{u,2}|^p)dx \\ & \quad + \int (|(\partial_1 + i\partial_2)f_{u,3}|^p + |\partial_3 f_{u,3}|^p)dx + \int (|(\partial_1 - i\partial_2)f_{u,4}|^p + |\partial_3 f_{u,4}|^p)dx \\ &=: F_1(u) + F_2(u) + F_3(u) + F_4(u). \end{aligned} \quad (4.13)$$

Therefore

$$\int_0^\infty d(u^q)u^{-q}M_{\alpha;p;p}(f_u)^p = \sum_{k=1}^4 \int_0^\infty d(u^q)u^{-q}[F_k(u) + F_{k+1}(u) + F_{k+2}(u) + F_{k+3}(u)].$$

We compute the integral of the first term on the right-hand side concerning  $F_1(u)$ . Before that, we note that

$$\begin{aligned} F_1(u) &= \int_{u \leq |f_1(x)| \leq cu} (|(\partial_1 + i\partial_2)f_{u,1}|^p + |\partial_3 f_{u,1}|^p)dx \\ &= \int_{u \leq |f_1(x)| \leq cu} (|(\partial_1 + i\partial_2)f_1|^p + |\partial_3 f_1|^p)dx, \end{aligned} \quad (4.14)$$

as the  $x$ -integration in the third member of (4.14) may be done only on the set  $\{x; u \leq |f_1(x)| \leq cu\}$  because  $f_{u,1}(x) = 0$  when  $|f_1(x)| \leq u$ , and  $f_{u,1}(x)$  is constant (with  $|f_{u,1}(x)| = (c-1)u$  when  $|f_1(x)| \geq cu$ ). Further, the last equality in (4.14) is due to the fact that  $\partial_j f_{u,1}(x) = \partial_j f_1(x)$ ,  $j = 1, 2, 3$ , on the set  $\{x; u \leq |f_1(x)| \leq cu\}$ .

Thus, through (4.14) we have

$$\begin{aligned} \int_0^\infty d(u^q)u^{-q}F_1(u) &= \int_0^\infty d(u^q)u^{-q} \int_{u \leq |f_1(x)| \leq cu} (|(\partial_1 + i\partial_2)f_1|^p + |\partial_3 f_1|^p)dx \\ &= q \int dx (|(\partial_1 + i\partial_2)f_1|^p + |\partial_3 f_1|^p) \int_{\frac{|f_1(x)|}{c}}^{|f_1(x)|} \frac{du}{u} \\ &= q(\log c) \int (|(\partial_1 + i\partial_2)f_1|^p + |\partial_3 f_1|^p)dx. \end{aligned} \quad (4.15)$$

In the same way for  $F_2(u)$ ,  $F_3(u)$ ,  $F_4(u)$  in (4.13), we can get

$$\begin{aligned}\int_0^\infty d(u^q)u^{-q}F_2(u) &= q(\log c) \int (|(\partial_1 - i\partial_2)f_2|^p + |\partial_3 f_2|^p)dx, \\ \int_0^\infty d(u^q)u^{-q}F_3(u) &= q(\log c) \int (|(\partial_1 + i\partial_2)f_3|^p + |\partial_3 f_3|^p)dx, \\ \int_0^\infty d(u^q)u^{-q}F_4(u) &= q(\log c) \int (|(\partial_1 - i\partial_2)f_4|^p + |\partial_3 f_4|^p)dx.\end{aligned}$$

So we obtain

$$\begin{aligned}\int_0^\infty d(u^q)u^{-q}M_{\alpha \cdot p; p}(f_u)^p &= q(\log c) [\|(\alpha \cdot p)P_{13}f\|_p^p + \|(\alpha \cdot p)P_{24}f\|_p^p] \\ &= q(\log c)M_{\alpha \cdot p; p}(f)^p,\end{aligned}$$

establishing (4.12) of the lemma.  $\square$

Then, noting (4.5)/(4.8) to put together (4.10) and (4.11) with Lemma 4.1, we get

$$\frac{1}{20^q}\|f\|_q^q \leq C_0 q(\log c)M_{\alpha \cdot p; p}(f)^p + \frac{q}{q-1} \frac{1}{c^{q-1}}\|f\|_q^q. \quad (4.16)$$

Thus, since  $\|f\|_q$  is finite by assumption, taking  $c$  sufficiently large in (4.16) and putting  $C = \frac{C_0 q(\log c)}{\frac{1}{20^q} - \frac{q}{q-1} \frac{1}{c^{q-1}}}$ , we have shown the desired inequality (4.3) in step II. In the whole arguments in Step II we need the condition  $f \in L^q(\mathbb{R}^3; \mathbb{C}^4)$ , i.e. that  $\|f\|_q < \infty$ , only here in (4.16) so that we can obtain inequality (4.3) from (4.16).

III. Finally we show that if  $M_{\alpha \cdot p; p}(f) < \infty$  and  $\|f\|_{B^{p/(p-q)}(\mathbb{R}^3; \mathbb{C}^4)} \leq 1$ , then  $f \in L^q(\mathbb{R}^3; \mathbb{C}^4)$ , and that  $\|f\|_q \leq C M_{\alpha \cdot p; p}(f)$  with a constant  $C$  independent of  $f$ .

We already know by the weak type inequality (1.14) that  $\|f\|_{q, \infty} < \infty$ . Therefore, in view of the second member of (4.8), we may consider, for every  $0 < \varepsilon < 1$ ,

$$N_\varepsilon(f) := \int_{u=\varepsilon}^{u=1/\varepsilon} |\{ |f|_{\ell^\infty} \geq 5u \}| d(u^q) < \infty. \quad (4.17)$$

Note that

$$\frac{1}{20^q}\|f\|_q^q \leq \lim_{\varepsilon \rightarrow 0} N_\varepsilon(f). \quad (4.18)$$

By modifying the arguments in (4.8)–(4.10) and (4.16), we obtain

$$\begin{aligned}N_\varepsilon(f) &\leq C_0 q(\log c)M_{\alpha \cdot p; p}(f)^p + \int_{u=\varepsilon}^{u=1/\varepsilon} \frac{1}{u} \left( \int |f(x)|_{\ell^\infty} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) dx \right) d(u^q) \\ &=: I_1 + I_2.\end{aligned} \quad (4.19)$$

The layer cake representation (4.4) leads the second term  $I_2$  on the right-hand side to

$$\begin{aligned}I_2 &= \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \frac{1}{u} d(u^q) \int \chi_{\{|f|_{\ell^\infty} > s\}}(x) \chi_{\{|f|_{\ell^\infty} > cu\}}(x) ds \\ &= \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \frac{1}{u} d(u^q) \left[ \int_0^{cu} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) ds + \int_{cu}^\infty \chi_{\{|f|_{\ell^\infty} > s\}}(x) ds \right] \\ &= c \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) d(u^q) + \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} qu^{q-2} du \int_{cu}^\infty \chi_{\{|f|_{\ell^\infty} > s\}}(x) ds.\end{aligned}$$

Then by integration by parts we have

$$\begin{aligned}
I_2 &= c \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) d(u^q) \\
&\quad + \int dx \left[ \frac{q}{q-1} u^{q-1} \int_{cu}^{\infty} \chi_{\{|f|_{\ell^\infty} > s\}}(x) ds \right]_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \\
&\quad + \int dx \int_{\varepsilon}^{\frac{1}{\varepsilon}} \frac{q}{q-1} u^{q-1} c \chi_{\{|f|_{\ell^\infty} > cu\}}(x) du \\
&= c \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) d(u^q) + \frac{c}{q-1} \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) d(u^q) \\
&\quad + \frac{q}{q-1} \left[ \frac{1}{\varepsilon^{q-1}} \int_{c\frac{1}{\varepsilon}}^{\infty} \chi_{\{|f|_{\ell^\infty} > s\}}(x) ds - \varepsilon^{q-1} \int_{c\varepsilon}^{\infty} \chi_{\{|f|_{\ell^\infty} > s\}}(x) ds \right] \\
&\leq \frac{cq}{q-1} \int dx \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} \chi_{\{|f|_{\ell^\infty} \geq cu\}}(x) d(u^q) + \frac{cq}{q-1} \frac{1}{\varepsilon^{q-1}} \int dx \int_{\frac{1}{\varepsilon}}^{\infty} \chi_{\{|f|_{\ell^\infty} \geq cu\}}(x) du \\
&= \frac{cq}{q-1} \int_{\varepsilon}^{\frac{1}{\varepsilon}} |\{|f|_{\ell^\infty} \geq cu\}| d(u^q) + \frac{cq}{q-1} \frac{1}{\varepsilon^{q-1}} \int_{\frac{1}{\varepsilon}}^{\infty} |\{|f|_{\ell^\infty} \geq cu\}| du \\
&=: I_{21} + I_{22}, \tag{4.20}
\end{aligned}$$

where the last equality is due to Fubini's theorem, so that  $I_2 \leq I_{21} + I_{22}$ . Changing, in  $I_{21}$  and  $I_{22}$ , the variable  $cu = 5s$  and writing  $u$  for  $s$  again, we see by (4.17) and by the definition (1.15) of weak  $L^q$  norm,

$$\begin{aligned}
I_{21} &= \frac{cq}{q-1} \int_{\varepsilon}^{\frac{1}{\varepsilon}} |\{|f|_{\ell^\infty} \geq cu\}| d(u^q) = \frac{cq}{q-1} \left(\frac{5}{c}\right)^q \int_{u=\frac{c}{5}\varepsilon}^{u=\frac{c}{5}\frac{1}{\varepsilon}} |\{|f|_{\ell^\infty} \geq 5u\}| d(u^q) \\
&= \frac{q}{q-1} \frac{5^q}{c^{q-1}} \left( \int_{u=\varepsilon}^{u=\frac{1}{\varepsilon}} + \int_{u=\frac{1}{\varepsilon}}^{u=\frac{c}{5}\frac{1}{\varepsilon}} - \int_{u=\varepsilon}^{u=\frac{c}{5}\varepsilon} \right) |\{|f|_{\ell^\infty} \geq 5u\}| d(u^q) \\
&\leq \frac{q}{q-1} \frac{5^q}{c^{q-1}} \left\{ N_\varepsilon(f) + \int_{u=\frac{1}{\varepsilon}}^{u=\frac{c}{5}\frac{1}{\varepsilon}} (5u)^{-q} (5u)^q |\{|f|_{\ell^\infty} \geq 5u\}| d(u^q) \right\} \\
&\leq \frac{q}{q-1} \frac{5^q}{c^{q-1}} \left\{ N_\varepsilon(f) + \|f\|_{q,\infty}^q \int_{u=\frac{1}{\varepsilon}}^{u=\frac{c}{5}\frac{1}{\varepsilon}} (5u)^{-q} d(u^q) \right\} \\
&= \frac{q}{q-1} \frac{5^q}{c^{q-1}} \left\{ N_\varepsilon(f) + \|f\|_{q,\infty}^q \frac{q}{5^q} \left[ \log u \right]_{\frac{1}{\varepsilon}}^{\frac{c}{5}\frac{1}{\varepsilon}} \right\} \\
&= \frac{q}{q-1} \frac{5^q}{c^{q-1}} N_\varepsilon(f) + \frac{q}{q-1} \frac{\log \frac{c}{5}}{c^{q-1}} \|f\|_{q,\infty}^q. \tag{4.21}
\end{aligned}$$

For  $I_{22}$  we have

$$\begin{aligned}
I_{22} &= \frac{cq}{q-1} \frac{1}{\varepsilon^{q-1}} \int_{\frac{1}{\varepsilon}}^{\infty} (cu)^{-q} [(cu)^q |\{|f|_{\ell^\infty} \geq cu\}|] du \\
&\leq \frac{cq}{q-1} \frac{1}{\varepsilon^{q-1}} \|f\|_{q,\infty}^q \int_{\frac{1}{\varepsilon}}^{\infty} (cu)^{-q} du = \frac{q}{(q-1)^2} \frac{1}{c^{q-1}} \|f\|_{q,\infty}^q. \tag{4.22}
\end{aligned}$$

Then

$$I_2 \leq I_{21} + I_{22} \leq \frac{q}{q-1} \frac{5^q}{c^{q-1}} N_\varepsilon(f) + \frac{q}{q-1} \frac{1}{c^{q-1}} \|f\|_{q,\infty}^q \left( \frac{1}{q-1} + \log \frac{c}{5} \right).$$

Therefore from (4.19)

$$\begin{aligned}
N_\varepsilon(f) &\leq I_1 + I_2 \\
&\leq C_0 q (\log c) M_{\alpha \cdot p; p}(f)^p + \frac{q}{q-1} \frac{5^q}{c^{q-1}} N_\varepsilon(f) \\
&\quad + \frac{q}{q-1} \frac{1}{c^{q-1}} \|f\|_{q, \infty}^q \left( \frac{1}{q-1} + \log \frac{c}{5} \right) \\
&\leq \frac{q}{q-1} \frac{5^q}{c^{q-1}} N_\varepsilon(f) + \left[ C_0 q (\log c) + \frac{q}{q-1} \frac{1}{c^{q-1}} \left( \frac{1}{q-1} + \log \frac{c}{5} \right) \right] M_{\alpha \cdot p; p}(f)^p,
\end{aligned}$$

where the last inequality is due to the fact that by (1.14) and (1.22)  $\|f\|_{q, \infty} \leq \|(\alpha \cdot p)f\|_p \leq M_{\alpha \cdot p; p}(f)$ . Then take  $c$  large (if necessary, larger than the  $c$  chosen once already at the end of step II) such that  $1 - \frac{q}{q-1} \frac{5^q}{c^{q-1}} < \frac{1}{2}$ , and we have with (4.18)

$$\|f\|_q^q \leq 2 \cdot 20^q \left[ C_0 q (\log c) + \frac{q}{q-1} \frac{1}{c^{q-1}} \left( \frac{1}{q-1} + \log \frac{c}{5} \right) \right] M_{\alpha \cdot p; p}(f)^p. \quad (4.23)$$

Thus, taking  $C := 2^{1/q} 20 \left[ C_0 q (\log c) + \frac{q}{q-1} \frac{1}{c^{q-1}} \left( \frac{1}{q-1} + \log \frac{c}{5} \right) \right]^{1/q}$  and noting homogeneity, we have shown the desired inequality (1.28), ending the proof of Theorem 1.1 (i).  $\square$

*Proof of Theorem 1.1 (ii).* In case  $p > 1$ , in our previous paper [IS] we have shown that  $H_{\alpha \cdot p, 0}^{1, p}(\mathbb{R}^3; \mathbb{C}^4) = H_0^{1, p}(\mathbb{R}^3; \mathbb{C}^4)$ , so that the norms  $\|f\|_{M_{\alpha \cdot p, 1, p}} := (\|f\|_p^p + M_{\alpha \cdot p; p}(f)^p)^{1/p}$  and  $\|f\|_{\alpha \cdot p, 1, p} := (\|f\|_p^p + \|(\alpha \cdot p)f\|_p^p)^{1/p}$  are equivalent to the norm  $\|f\|_{1, p} := (\|f\|_p^p + \|\nabla f\|_p^p)^{1/p}$ . But this may not be sufficient to derive (1.30).

To show the assertion, we need show that for  $p > 1$  the two semi-norms  $\|(\alpha \cdot p)f\|_p$  and  $\|\nabla f\|_p$  are equivalent. However, noting the two inequalities (1.22), we have only to show the following lemma.

**Lemma 4.2.** *For  $1 < p < \infty$ , there exists a positive constant  $C$  such that*

$$\|\nabla f\|_p \leq C \|(\alpha \cdot p)f\|_p \quad (4.24)$$

for every  $f \in C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$ .

*Proof.* We give two proofs.

(i) (A first proof with functional analysis) In the proof of [IS, Proposition 3.1], we had already seen this fact of the lemma. Here let us briefly sketch the argument.

Let  $f = {}^t(f_1, f_2, f_3, f_4) \in C_0^\infty(\mathbb{R}^3; \mathbb{C}^4)$  so that  $(\alpha \cdot p)f \in L^p(\mathbb{R}^3; \mathbb{C}^4)$ , and

$$g = {}^t(g_1, g_2, g_3, g_4) := (\alpha \cdot p)f = -i[\alpha_1 \partial_1 f + \alpha_2 \partial_2 f + \alpha_3 \partial_3 f],$$

belongs to  $L^p(\mathbb{R}^3; \mathbb{C}^4)$ .

Since  $-\Delta f = (\alpha \cdot p)^2 f = (\alpha \cdot p)g$ , we have  $\Delta(\partial_j f) = i[\alpha_1 \partial_1 + \alpha_2 \partial_2 + \alpha_3 \partial_3] \partial_j g$ , ( $j = 1, 2, 3$ ), where the derivatives are taken in distribution sense. Then we can show for each  $j = 1, 2, 3$ ,  $k = 1, 2, 3, 4$ , that there exist constants  $C_{j, kl}$ ,  $k, l = 1, 2, 3, 4$ , such that

$$\begin{aligned}
|\langle \partial_j f_k, \Delta \phi \rangle| &\leq [(C_{j, k1} \|g_1\|_p + C_{j, k2} \|g_2\|_p + C_{j, k3} \|g_3\|_p + C_{j, k4} \|g_4\|_p] \|\Delta \phi\|_{p'} \\
&\leq C \left( \sum_{l=1}^4 \|g_l\|_p^p \right)^{1/p} \|\Delta \phi\|_{p'} = C \|g\|_p \|\psi\|_{p'}
\end{aligned}$$

for all  $\phi \in C_0^\infty(\mathbb{R}^3)$  with  $C := (\sum_{l=1}^4 C_{j,kl}^{p'})^{1/p'}$ , where the last second inequality is due to Hölder's inequality with  $\frac{1}{p} + \frac{1}{p'} = 1$ . Hence  $|\langle \partial_j f_k, \psi \rangle| \leq C \|g\|_p \|\psi\|_{p'}$  for all  $\psi \in L^{p'}(\mathbb{R}^3)$ , since for  $p > 1$  the space  $\Delta(C_0^\infty(\mathbb{R}^3))$  is dense in  $L^{p'}(\mathbb{R}^3)$ , so that  $\partial_j f_k$  belongs to  $L^p(\mathbb{R}^3)$  for  $j = 1, 2, 3$ ,  $k = 1, 2, 3, 4$ , and

$$\|\partial_j f_k\|_p \leq C \|g\|_p = C \|(\alpha \cdot p) f\|_p.$$

This proves the desired inequality (4.24).

(ii) (A second proof with pseudodifferential calculus)

To show the assertion, we have only to show that for  $j = 1, 2, 3$ ,  $-i\partial_j/(\alpha \cdot p)$  is a bounded operator on  $L^p(\mathbb{R}^3; \mathbb{C}^2)$ . To see it, since  $(\alpha \cdot p)^2 = -\Delta$ , we note that

$$\frac{-i\partial_j}{\alpha \cdot p} = \frac{-i\partial_j}{-\Delta}(\alpha \cdot p) = \frac{-i\partial_j}{(-\Delta)^{1/2}} \sum_{k=1}^3 \frac{-i\alpha_k \partial_k}{(-\Delta)^{1/2}} = - \sum_{k=1}^3 \alpha_k \cdot R_j R_k,$$

where  $R_k = \frac{-i\partial_k}{(-\Delta)^{1/2}}$ ,  $k = 1, 2, 3$ , is the Riesz transform which is a pseudo-differential operator having symbol  $i\xi_k/|\xi|$ , and if  $1 < p < \infty$ , we have  $\|R_k g\|_p \leq C \|g\|_p$  with a constant  $C > 0$ , e.g. by the Calderon-Zygmund theorem [e.g. S, 4.2, Theorem 3, p.29] or by Fefferman's theorem [Fe, Theorem, a, p.414]. Therefore we obtain for each  $j = 1, 2, 3$ ,

$$\|[-i\partial_j/(\alpha \cdot p)]f\|_p \leq 3C^2 \|f\|_p.$$

This proves (4.24), again showing the lemma.  $\square$

Thus we have proved Theorem 1.1 (ii), completing the proof of Theorem 1.1.  $\square$

## 5 Proof of Corollaries 1.2, 1.3, Theorem 1.4 , Corollary 1.5 and Theorem 1.6

*Proof of Corollary 1.2.* Let  $h := {}^t(h_1, h_2)$  be a  $\mathbb{C}^2$ -valued function and put  $f = {}^t(f_1, f_2, f_3, f_4)$  with  $f_1 = h_1$ ,  $f_2 = h_2$ ,  $f_3 = f_4 = 0$ . Then (1.33) is nothing but (1.28). This proves Corollary 1.2 (i). (ii) can be seen as in the proof of Theorem 1.1 (ii).  $\square$

*Proof of Corollary 1.3.* Let  $\psi$  be a  $\mathbb{C}$ -valued function and put  $f = {}^t(f_1, f_2, f_3, f_4)$  with  $f_2 = \psi$ ,  $f_1 = f_3 = f_4 = 0$ . Then (1.37) is nothing but (1.28). This proves Corollary 1.3 (i). (ii) can be seen as in the proof of Theorem 1.1 (ii).  $\square$

*Proof of Theorem 1.4.* The proof is divided into two parts (a) and (b). First in (a), we show (1.41) for the operator  $(\sigma \cdot p)^{(a)}$  in (1.40a), and then in (b) for the other two  $(\sigma \cdot p)^{(b)}$ ,  $(\sigma \cdot p)^{(c)}$  in (1.40bc).

(a) The case for  $(\sigma \cdot p)^{(a)}$  in (1.40a): First we are going to show (4.16) with  $\alpha \cdot p$  replaced by  $(\sigma \cdot p)^{(a)}$  in (1.40a), and then the proof proceeds to use almost the same arguments as in steps I, II, III of the proof of Theorem 1.1 (i). In step II we shall not need introduce some other semi-norm like  $M_{(\alpha \cdot p)^{(a)}, p}(f)$  than  $\|(\alpha \cdot p)^{(a)} f\|_p$ , and have only to go with the semi-norm  $\|(\sigma \cdot p)^{(a)} f\|_p$  for  $\mathbb{C}^2$ -valued functions  $f(x) = {}^t(f_1(x), f_2(x))$  on  $\mathbb{R}^2$ .

I. In the same way as before, we can show an inequality corresponding to (1.14), i.e. that there exists a constant  $C$  such that  $\|f\|_{q,\infty} \leq C\|(\sigma \cdot p)^{(a)} f\|_p^{p/q} \|f\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)}$  for every  $f = {}^t(f_1, f_2)$  which satisfies  $(\sigma \cdot p)^{(a)} f \in L^p(\mathbb{R}^2; \mathbb{C}^2)$  and belongs to  $B_{\infty,\infty}^{p/(p-q)}(\mathbb{R}^2; \mathbb{C}^2)$ .

II. This step contains a slight improvement in its own. We want to replace the weak  $L^q$  norm by the strong  $L^q$  norm. Under the same hypothesis as in step I above but with  $\|f\|_{B_{\infty,\infty}^{p/(p-q)}} \leq 1$ , we are going to show the following inequality:

$$\int |f|_{\ell^q}^q dx \leq C\|(\sigma \cdot p)^{(a)} f\|_p^p \quad (5.1)$$

with a constant  $C$  independent of  $f$ , a sharper inequality than the previous (4.3), assuming the extra condition  $f \in L^q(\mathbb{R}^2; \mathbb{C}^2)$ , which will turn out to be unnecessary in step III below.

To this end, we can proceed as in II of the proof of Theorem 1.1 (i), Section 4, to obtain an analogous version of (4.8) :

$$\begin{aligned} \frac{1}{20^q} \|f\|_q^q &= \int_0^\infty |\{ |f|_{\ell^q} \geq 20u \}| d(u^q) \\ &\leq \int_0^\infty |\{ \vee_{k=1}^2 |f_{u,k} - e^{t_u \Delta} f_{u,k}| \geq u \}| d(u^q) \\ &\quad + \int_0^\infty |\{ \vee_{k=1}^2 \{ e^{t_u \Delta} |f_{u,k} - f_k| \geq 2u \} \}| d(u^q) \\ &=: J'_1 + J'_2, \end{aligned}$$

where  $\Delta$  is the Laplacian in  $\mathbb{R}^2$ ,  $f(x) := {}^t(f_1(x), f_2(x)) \in L^q(\mathbb{R}^2; \mathbb{C}^2)$  with  $\|f\|_{B_{\infty,\infty}^{p/(p-q)}} \leq 1$  and  $f_u(x) := {}^t(f_{u,1}(x), f_{u,2}(x))$  is given by (4.6) with the subscription moving over  $\{1, 2\}$ , not  $\{1, 2, 3, 4\}$ . By the same arguments used before to get (4.10) and (4.11), respectively, we have  $J'_2 \leq \frac{q}{q-1} \frac{1}{c^{q-1}} \|f\|_q^q$  and

$$\begin{aligned} J'_1 &\leq C_0 \int_0^\infty d(u^q) u^{-q} \|(\sigma \cdot p)^{(a)} f_u\|_p^p \\ &= C_0 \int_0^\infty d(u^q) u^{-q} [\|(\partial_1 + i\partial_2) f_{u,1}\|_p^p + \|(\partial_1 - i\partial_2) f_{u,2}\|_p^p]. \end{aligned}$$

Noting that

$$\begin{aligned} \int_0^\infty d(u^q) u^{-q} \|(\partial_1 + i\partial_2) f_{u,1}\|_p^p &= \int d(u^q) u^{-q} \int_{u \leq |f_1(x)| \leq cu} |(\partial_1 + i\partial_2) f_{u,1}(x)|^p dx \\ &= \int_0^\infty d(u^q) u^{-q} \int_{u \leq |f_1(x)| \leq cu} |(\partial_1 + i\partial_2) f_1(x)|^p dx \\ &= q(\log c) \int |(\partial_1 + i\partial_2) f_1(x)|^p dx \\ &= q(\log c) \|(\partial_1 + i\partial_2) f_1\|_p^p, \end{aligned}$$

and in the same way

$$\int_0^\infty d(u^q) u^{-q} \|(\partial_1 - i\partial_2) f_{u,2}\|_p^p = q(\log c) \|(\partial_1 - i\partial_2) f_2\|_p^p,$$



we have  $J'_1 \leq C_0 q (\log c) \|(\sigma \cdot p)^{(a)} f\|_p^p$ . Thus

$$\frac{1}{20^q} \|f\|_q^q \leq J'_1 + J'_2 \leq C_0 q (\log c) \|(\sigma \cdot p)^{(a)} f\|_p^p + \frac{q}{q-1} \frac{1}{c^{q-1}} \|f\|_q^q,$$

whence we get the desired inequality (4.1), taking  $C = \frac{C_0 q (\log c)}{\frac{1}{20^q} - \frac{q}{q-1} \frac{1}{c^{q-1}}}$  for  $c$  sufficiently large.

III. Finally we remove the condition that  $f \in L^q(\mathbb{R}^2; \mathbb{C}^2)$  assumed in step II. In fact, we show that if  $\|(\sigma \cdot p)^{(a)} f\| < \infty$  and  $\|f\|_{B_{\infty, \infty}^{p/(p-q)}} \leq 1$ , then  $f \in L^q(\mathbb{R}^2; \mathbb{C}^2)$ .

The proof proceeds in the same way as in III of the proof of Theorem 1.1 (i), Section 4. Indeed, with the corresponding  $N_\varepsilon(f)$  as in (4.17) and (4.18), we can show, instead of (4.19),

$$N_\varepsilon(f) \leq C_0 q (\log c) \|(\sigma \cdot p)^{(a)} f\|_p^p + \int_{u=\varepsilon}^{u=1/\varepsilon} \frac{1}{u} \left( \int |f(x)|_{\ell^\infty} \chi_{\{|f|_{\ell^\infty} > cu\}}(x) dx \right) d(u^q).$$

Estimating, in the same way as before, the two terms on the right-hand side, we can obtain the desired inequality  $\|f\|_q^q \leq C \|(\sigma \cdot p)^{(a)} f\|_p^p$ . This shows (1.41) in Theorem 1.4 for  $(\sigma \cdot p)^{(a)}$  in (1.40a).

(b) The other cases for  $(\sigma \cdot p)^{(b)}$  and  $(\sigma \cdot p)^{(c)}$  in (1.40bc): Each of these two cases is reduced to the case (a) for  $(\sigma \cdot p)^{(a)}$  by a linear transformation. The idea is based on the following lemma.

**Lemma 5.1.** *The three 2-dimensional Weyl–Dirac (or Pauli) operators  $(\sigma \cdot p)^{(a)}$ ,  $(\sigma \cdot p)^{(b)}$ ,  $(\sigma \cdot p)^{(c)}$  in (1.40abc) are unitarily equivalent. In fact, there exist unitary  $2 \times 2$ -matrices  $N$ ,  $N'$  such that for  $f = {}^t(f_1, f_2)$  and  $h = {}^t(h_1, h_2) := Nf$ ,  $h = {}^t(h_1, h_2) := N'f$ ,*

$$(\sigma \cdot p)^{(a)} h = (\sigma \cdot p)^{(a)} Nf = N(\sigma \cdot p)^{(b)} f, \text{ with } h = {}^t(h_1, h_2) = Nf, \quad (5.2)$$

$$(\sigma \cdot p)^{(a)} h = (\sigma \cdot p)^{(a)} N'f = N'(\sigma \cdot p)^{(c)} f, \text{ with } h = {}^t(h_1, h_2) = N'f. \quad (5.3)$$

*Proof.* Take matrices  $N := \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i \\ 1 & i \end{pmatrix}$ ,  $N' := \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix}$ , which are unitary. We have  $N^{-1} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ i & -i \end{pmatrix}$ ,  $(N')^{-1} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$ , and

$$\begin{aligned} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i \\ 1 & i \end{pmatrix} \begin{pmatrix} \partial_1 & \partial_2 \\ \partial_2 & -\partial_1 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ i & -i \end{pmatrix} &= \begin{pmatrix} 0 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & 0 \end{pmatrix}, \\ \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} \partial_1 & -i\partial_2 \\ i\partial_2 & -\partial_1 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} &= \begin{pmatrix} 0 & \partial_1 - i\partial_2 \\ \partial_1 + i\partial_2 & 0 \end{pmatrix}. \end{aligned}$$

Taking into account the definition (1.40abc) of  $(\sigma \cdot p)^{(a)}$ ,  $(\sigma \cdot p)^{(b)}$ ,  $(\sigma \cdot p)^{(c)}$  yields (5.2) and (5.3), showing Lemma 5.1.  $\square$

Now we continue the proof (b) of Theorem 1.4. Take the same matrices  $N$  and  $N'$  as in Lemma 5.1, which we see reduce the cases  $(\sigma \cdot p)^{(b)}$  and  $(\sigma \cdot p)^{(c)}$  to the case  $(\sigma \cdot p)^{(a)}$ .

Note the bounds of the matrix norms of them and their inverses that for  $1 \leq r \leq \infty$ ,

$$\begin{aligned} \|N\|_{\ell^r \rightarrow \ell^r} &\leq \sqrt{2}, & \|N^{-1}\|_{\ell^r \rightarrow \ell^r} &\leq \sqrt{2}; \\ \|N'\|_{\ell^r \rightarrow \ell^r} &\leq \sqrt{2}, & \|(N')^{-1}\|_{\ell^r \rightarrow \ell^r} &\leq \sqrt{2}. \end{aligned} \quad (5.4)$$

It follows that if  $h = Nf$  or  $h = N'f$ , then

$$\|f\|_r \leq \sqrt{2}\|h\|_r, \quad \|h\|_r \leq \sqrt{2}\|f\|_r. \quad (5.5)$$

First, we treat the case  $(\sigma \cdot p)^{(b)}$  with  $N$ . We have by (5.2) in Lemma 5.1 and (5.4)

$$\|(\sigma \cdot p)^{(a)}h\|_p = \|N(\sigma \cdot p)^{(b)}f\|_p \leq \|N\|_{\ell^p \rightarrow \ell^p} \|(\sigma \cdot p)^{(b)}f\|_p \leq \sqrt{2}\|(\sigma \cdot p)^{(b)}f\|_p. \quad (5.6)$$

We note that  $P_t$  commutes with  $N$  to get

$$\begin{aligned} |(P_th)(x)|_{\ell^\infty} &= |(P_tNf)(x)|_{\ell^\infty} = |(NP_tf)(x)|_{\ell^\infty} \\ &\leq \|N\|_{\ell^\infty \rightarrow \ell^\infty} |(P_tf)(x)|_{\ell^\infty} \leq \sqrt{2}|(P_tf)(x)|_{\ell^\infty}, \end{aligned}$$

whence

$$\begin{aligned} \|h\|_{B_{\infty,\infty}^{p/(p-q)}} &= \sup_{t>0} \|P_th\|_\infty = \sup_{t>0} \sup_x |(P_th)(x)|_{\ell^\infty} \\ &\leq \sqrt{2} \sup_{t>0} \sup_x |(P_tf)(x)|_{\ell^\infty} = \sqrt{2} \sup_{t>0} \|P_tf\|_\infty = \sqrt{2}\|f\|_{B_{\infty,\infty}^{p/(p-q)}}. \end{aligned} \quad (5.7)$$

Then, since we already (1.41) holds for  $(\sigma \cdot p)^{(a)}$  with  $h$  in place of  $f$ , we combine it with (5.5), (5.6), (5.7) to get

$$\|f\|_q \leq \sqrt{2}\|h\|_q \leq \sqrt{2}C\|(\sigma \cdot p)^{(a)}h\|_p^{p/q} \|h\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)} \leq 2C\|(\sigma \cdot p)^{(b)}f\|_p^{p/q} \|f\|_{B_{\infty,\infty}^{p/(p-q)}}^{1-(p/q)},$$

which yields the desired inequality (1.41) for  $(\sigma \cdot p)^{(b)}$ .

Next, as for the other last case  $(\sigma \cdot p)^{(c)}$ , exactly the same arguments apply to it as those just made in the case  $(\sigma \cdot p)^{(b)}$  above, with the matrix  $N$ , relation (5.2) replaced by the matrix  $N'$ , relation (5.3).

This completes the proof of Theorem 1.4.  $\square$

*Proof of Corollary 1.5.* (1.42) follows from Corollary 1.3 (1.39) because our function  $\psi(x) = \psi(x_1, x_2)$  here is independent of  $x_3$ , or from Theorem 1.4 (1.41) for  $h = {}^t(h_1, h_2)$  with  $h_1 = 0$ ,  $h_2 = \psi$ .  $\square$

*Proof of Theorem 1.6.* The proof is done by analogous arguments used to prove Theorem 1.1. We only note that Lemma 4.1 is replaced by the following lemma, which can be shown in the same way as before.

**Lemma 5.2.** For  $f = {}^t(f_1, f_2, f_3, f_4)$ , one has

$$\int_0^\infty d(u^q)u^{-q}M_{\beta \cdot p; p}(f_u)^p = q(\log c)M_{\beta \cdot p; p}(f)^p. \quad (5.8)$$

Here we only note with (1.50) that the proof turns out to deal, instead of (4.13), with

$$\begin{aligned} &M_{\beta \cdot p; p}(f_u)^p \\ &= \int (|(\partial_1 + i\partial_2)f_{u,1}|^p + |(\partial_3 + i\partial_4)f_{u,1}|^p)dx + \int (|(\partial_1 - i\partial_2)f_{u,2}|^p + |(\partial_3 - i\partial_4)f_{u,2}|^p)dx \\ &\quad + \int (|(\partial_1 + i\partial_2)f_{u,3}|^p + |(\partial_3 - i\partial_4)f_{u,3}|^p)dx + \int (|(\partial_1 - i\partial_2)f_{u,4}|^p + |(\partial_3 + i\partial_4)f_{u,4}|^p)dx. \end{aligned}$$

$\square$

## 6 Concluding Comments

We have originated a version of improved Sobolev embedding theorem for vector-valued functions involved with for the three-dimensional Dirac operator  $D = \alpha \cdot p$ , the three-dimensional Weyl–Dirac (or Pauli) operator  $D = \sigma \cdot p$ , and the four-dimensional Euclidian Dirac operator  $D = \beta \cdot p$ . To this end we have introduced in Section 1 the corresponding first-order-derivative semi-norms  $M_{\alpha \cdot p; p}(f)$ ,  $M_{\sigma \cdot p; p}(h)$  and  $M_{\beta \cdot p; p}(f)$  by decomposing them into two parts:  $D = D_1 + D_2$ . Although the used decomposition looked to be artificial, it turns out there are other meaningful decompositions which give the same semi-norms as thus defined. In fact, we have characterized, in Proposition 1.0 for  $\alpha \cdot p$  and its counterpart for  $\sigma \cdot p$  and  $\beta \cdot p$ , which kind of decompositions are fit for our semi-norms at all. It turns out that they should be those which satisfy the condition that *each row of the matrices of both the parts  $D_1$  and  $D_2$  contains only one nonzero entry*. Why one needs this condition is simply because our proof given in Section 4 needs it.

In this section we will make some further comments and observe that after all this semi-norm is of reasonably good and optimal choice, having intrinsic and universal character and being an intermediate one in strength lying between both the semi-norm  $\|(\alpha \cdot p)f\|_p$ ,  $\|(\sigma \cdot p)h\|_p$  or  $\|(\beta \cdot p)f\|_p$  and the seminorm  $\|\nabla f\|_p$ ,  $\|\nabla h\|_p$  or  $\|\nabla f\|_p$ , respectively. We describe only with the 4-dimensional Euclidian Dirac operator, as we can deal with the other two operators just in the same way.

So consider the 4-dimensional Euclidian Dirac operator  $D := \beta \cdot p$  in (1.46) and its decomposition into the sum of its two parts :  $D = D_1 + D_2$ . Ignoring the order of the pair  $(D_1, D_2)$ , we regard the two decomposition  $(D_1, D_2)$  and  $(D_2, D_1)$  as the same. Then there are totally  $\frac{1}{2} \cdot 2^7 = 64$  decompositions including the trivial decomposition with  $(D_1, D_2) = (D, 0)$  or  $(D_1, D_2) = (0, D)$ . The set of all decompositions of  $D = \beta \cdot p$  is denoted by  $\text{Decom}(D)$ . Let  $\text{Decom}_1(D)$  be the subset of all  $(D_1, D_2)$  in  $\text{Decom}(D)$  which satisfy the condition that *each row of  $D_1$  and  $D_2$  contains only one nonzero entry*. It is seen that  $\text{Decom}_1(D)$  consists of  $\frac{1}{2} \cdot 2^4 = 8$  decompositions of  $D$ . The decompositions (1.47), (1.53), (1.54) and (1.55) are examples of elements of  $\text{Decom}_1(D)$ . With the decomposition (1.47), i.e.  $((\beta \cdot p)P_{13}, (\beta \cdot p)P_{24})$ , we have defined the semi-norm  $M_{\beta \cdot p; 1}(f)$  by (1.48), that is,

$$M_{\beta \cdot p; p}(f) := [\|(\beta \cdot p)P_{13}f\|_p + \|(\beta \cdot p)P_{24}f\|_p]^{1/p}. \quad (6.1)$$

We have shown Theorem 1.6, a version of improved Sobolev embedding theorem for vector-valued functions, that inequality (1.57) holds with this semi-norm  $M_{\beta \cdot p; 1}(f)$  for  $1 \leq p < \infty$ , and also seen in Section 3 that in case of  $p = 1$  one cannot replace the semi-norm  $M_{\beta \cdot p; p}(f)$  on the right by a weaker one  $\|(\beta \cdot p)f\|_p$ , though one can for  $1 < p < \infty$ . Actually we have

$$(D_1, D_2) \in \text{Decom}_1(D) \Rightarrow M_{\beta \cdot p; p}(f) := M_{D_1 \vee D_2; p}(f) = [\|D_1 f\|_p + \|D_2 f\|_p]^{1/p}. \quad (6.2)$$

Thus our semi-norm  $M_{\beta \cdot p; p}(f)$  is characterized as the one associated with  $\text{Decom}_1(D)$ . At this point also notice that this semi-norm has the *very* expression (1.50) with symmetric arrangement of eight terms in its last member. Inequality (1.51) shows that  $M_{\beta \cdot p; p}(f)$  is lying in strength between the semi-norms  $\|(\beta \cdot p)f\|_p$  and  $\|\nabla f\|_p$ . Notice that the condition that *each row of  $D_1$  and  $D_2$  contains only one nonzero entry* is satisfied by neither the 3-dimensional Dirac operator (1.17), 3-dimensional Weyl–Dirac (or

Pauli) operator (1.31) nor 4-dimensional Euclidian Dirac operator (1.46) *themselves*. Otherwise, our proof could establish for  $p = 1$  inequality (1.28) of Theorem 1.1, (1.35) of Corollary 1.2 and (1.57) of Theorem 1.6 with the semi-norm  $\|(\alpha \cdot p)f\|_1$ ,  $\|(\sigma \cdot p)h\|_1$  and  $\|(\beta \cdot p)f\|_1$  in place of  $M_{\alpha \cdot p;1}(f)$ ,  $M_{\sigma \cdot p;1}(h)$  and  $M_{\beta \cdot p;1}(f)$  on the right-hand side. But this is not in general possible because we have counterexamples as given in Section 3.

In Kähler Geometry and/or Spin Geometry (e.g. [Fr], [LawM]), the four-dimensional Euclidian Dirac operator  $D$  appears as an operator acting on the Clifford algebra  $CL(\mathbb{R}^4)$ , which is canonically isomorphic to the exterior algebra  $\Lambda^*(\mathbb{R}^4) \equiv \Lambda^*(T^*(\mathbb{R}^4))$ . On this  $\Lambda^*(\mathbb{R}^4)$ , in turn, there act two canonical first-order differential operators, namely, the exterior derivative  $d : \Lambda^*(\mathbb{R}^4) \rightarrow \Lambda^*(\mathbb{R}^4)$  and its formal adjoint  $d^* : \Lambda^*(\mathbb{R}^4) \rightarrow \Lambda^*(\mathbb{R}^4)$ , which satisfy  $d^2 = d^{*2} = 0$ . Then the fact is that the Dirac operator  $D$  is considered to decompose into their sum:  $D \cong d + d^*$ . In passing, it is conversely along with such a decomposition that the Dirac operator of *even* infinite dimension is defined on a Fock space in [A1, 2].

In this connection, notice that  $\text{Decom}_1(D)$  contains two pairs  $(D_1, D_2)$ , (1.53) and (1.54), which satisfy the one condition  $D_2 = D_1^*$ , but neither of the elements of  $\text{Decom}_1(D)$  satisfy the other condition  $D_1^2 = D_2^2 = 0$ . We ask: how about the inequality

$$\|f\|_q \leq CM_{D_1 \vee D_2; p}(f)^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)} \quad (6.3)$$

like (1.57) for the decompositions *not* belonging to  $\text{Decom}_1(D)$ , to hold with a fixed constant  $C > 0$  for all functions  $f(x) = {}^t(f_1(x), f_2(x), f_3(x), f_4(x))$  on  $\mathbb{R}^4$ ? To answer it, consider the following three decompositions  $D \equiv \beta \cdot p = D_1 + D_2$  in  $\text{Decom}(D) \setminus \text{Decom}_1(D)$  which are typical in some sense :

$M_{\beta}^{(4)}$

$$\begin{aligned} \beta \cdot p &= (\beta \cdot p)P_{12} + (\beta \cdot p)P_{34} \\ &= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ p_3 + i p_4 & p_1 - i p_2 & 0 & 0 \\ p_1 + i p_2 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & p_3 - i p_4 & p_1 - i p_2 \\ 0 & 0 & p_1 + i p_2 & -(p_3 + i p_4) \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \end{aligned} \quad (6.4a)$$

$$\begin{aligned} M_{\beta \cdot p;1}^{(4)}(f) &:= \|(\beta \cdot p)P_{12}f\|_1 + \|(\beta \cdot p)P_{34}f\|_1 \\ &= \|(\partial_1 + i\partial_2)f_1 - (\partial_3 - i\partial_4)f_2\|_1 + \|(\partial_1 - i\partial_2)f_2 + (\partial_3 + i\partial_4)f_1\|_1 \\ &\quad + \|(\partial_1 + i\partial_2)f_3 - (\partial_3 + i\partial_4)f_4\|_1 + \|(\partial_1 - i\partial_2)f_4 + (\partial_3 - i\partial_4)f_3\|_1; \end{aligned} \quad (6.4b)$$

$$\underline{M_\beta^{(5)}}$$

$$\begin{aligned} \beta \cdot p &= \begin{pmatrix} 0 & 0 & p_3 - i p_4 & p_1 - i p_2 \\ 0 & 0 & 0 & 0 \\ 0 & p_1 - i p_2 & 0 & 0 \\ 0 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & p_1 + i p_2 & -(p_3 + i p_4) \\ p_3 + i p_4 & 0 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} \\ &=: (\beta \cdot p)_5 + (\beta \cdot p)_6, \end{aligned} \quad (6.5a)$$

$$\begin{aligned} M_{\beta \cdot p; 1}^{(5)}(f) &:= \|(\beta \cdot p)_5 f\|_1 + \|(\beta \cdot p)_6 f\|_1 \\ &= \|(\partial_1 + i\partial_2)f_1\|_1 + \|(\partial_3 + i\partial_4)f_1\|_1 + \|(\partial_1 - i\partial_2)f_2\|_1 + \|(\partial_3 - i\partial_4)f_2\|_1 \\ &\quad + \|(\partial_1 + i\partial_2)f_3 - (\partial_3 + i\partial_4)f_4\|_1 + \|(\partial_1 - i\partial_2)f_4 + (\partial_3 - i\partial_4)f_3\|_1; \end{aligned} \quad (6.5b)$$

$$\underline{M_\beta^{(6)}}$$

$$\begin{aligned} \beta \cdot p &= \begin{pmatrix} 0 & 0 & p_3 - i p_4 & p_1 - i p_2 \\ 0 & 0 & 0 & -(p_3 + i p_4) \\ 0 & p_1 - i p_2 & 0 & 0 \\ 0 & -(p_3 - i p_4) & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & p_1 + i p_2 & 0 \\ p_3 + i p_4 & 0 & 0 & 0 \\ p_1 + i p_2 & 0 & 0 & 0 \end{pmatrix} \\ &=: (\beta \cdot p)_7 + (\beta \cdot p)_8, \end{aligned} \quad (6.6a)$$

$$\begin{aligned} M_{\beta \cdot p; 1}^{(6)}(f) &:= \|(\beta \cdot p)_7 f\|_1 + \|(\beta \cdot p)_8 f\|_1 \\ &= \|(\partial_1 + i\partial_2)f_1\|_1 + \|(\partial_3 + i\partial_4)f_1\|_1 + \|(\partial_1 - i\partial_2)f_2\|_1 + \|(\partial_3 - i\partial_4)f_2\|_1 \\ &\quad + \|(\partial_1 + i\partial_2)f_3\|_1 + \|(\partial_3 + i\partial_4)f_4\|_1 + \|(\partial_1 - i\partial_2)f_4 + (\partial_3 - i\partial_4)f_3\|_1. \end{aligned} \quad (6.6b)$$

Here the first decomposition (6.4a) and the second (6.5a) enjoy the same property as the Dirac operator  $D$  mentioned above in connection with Kähler Geometry and/or Spin Geometry. Further, the former (6.4a), which we have already referred to in Section 1 below Proof of Proposition 1.0 and also below equations (1.56a, b, c), has a beauty of symmetry. The latter (6.4a) has another beauty that each nonzero entry of  $(\beta \cdot p)_6$  is either of the two Cauchy–Riemann operators in the variables  $(x_1, x_2)$  and  $(x_3, x_4)$ , while that of  $(\beta \cdot p)_5$  either of their adjoints. The third decomposition (6.5a), which is a slight modification of (6.4a), looks artificial, lacking in beauty of symmetry and satisfying neither  $(\beta \cdot p)_8 = (\beta \cdot p)_7^*$  nor  $(\beta \cdot p)_7^2 = (\beta \cdot p)_8^2 = 0$ .

Our answer from the present paper is affirmative for  $1 < p < \infty$ , as already shown in Theorem 1.6 (ii), because, for any decomposition  $(D_1, D_2) \in \text{Decom}(D)$ , the semi-norm  $M_{D_1 \vee D_2; p}(f)$  is equivalent to the semi-norms  $\|(\beta \cdot p)f\|_p$  and  $\|\nabla f\|_p$  as seen in (1.51). However, as for  $p = 1$ , it will be negative, so long as one requires that  $D_1^2 = D_2^2 = 0$ . Thus the problem is when  $p = 1$ .

Comparing with the semi-norm  $M_{\beta \cdot p; 1}(f)$  in (1.50) for  $p = 1$ , we note (cf. (1.51))

$$\|(\beta \cdot p)f\|_1 = M_{\beta \cdot p; 1}^{(4)}(f) \leq M_{\beta \cdot p; 1}^{(5)}(f) \leq M_{\beta \cdot p; 1}^{(6)}(f) \leq M_{\beta \cdot p; 1}(f) \leq \|\nabla f\|_1, \quad (6.7)$$

where these three semi-norms are not equivalent to one another. Hence we also realize that  $M_{\beta \cdot p; 1}^{(6)}(f)$  is *next* weaker than  $M_{\beta \cdot p; 1}(f)$ , and  $M_{\beta \cdot p; 1}^{(5)}(f)$  is *next* weaker than  $M_{\beta \cdot p; 1}^{(6)}(f)$ .

**Proposition 6.1.** *For  $p = 1$ , inequality (6.3) does not hold with the semi-norm  $M_{D_1 \vee D_2; 1}(f)$  replaced by  $M_{\beta \cdot p; 1}^{(4)}(f)$  in (6.4b) and  $M_{\beta \cdot p; 1}^{(5)}(f)$  in (6.5b) corresponding to the decompositions (6.4a) and (6.5a), respectively.*

The proof of Proposition 6.1 is omitted. We give only some notes here. As to  $M_{\beta \cdot p; 1}^{(4)}(f)$  in (6.4b), the assertion is clear, because the last member of this semi-norm is the same as (1.49), namely,  $M_{\beta \cdot p; 1}^{(4)}(f) = \|(\beta \cdot p)f\|_1$ . As to  $M_{\beta \cdot p; 1}^{(5)}(f)$  in (6.5b), we can show the same sequence  $\{f_n\}_{n=1}^\infty$  used to construct the counterexample in Section 3 violates inequality (6.3) for  $p = 1$ ,  $q = \frac{4}{3}$ , so that  $\frac{p}{p-q} = -3$ .

It should be probably appropriate to mention here whether the present work has any connection with those of [BoBr] and [LanSt]. They proved an inequality of the form

$$\|u\|_{n/(n-1)} \leq C (\|du\|_1 + \|d^*u\|_1)$$

holds with a constant  $C > 0$  for all smooth  $m$ -forms  $u$  on  $\mathbb{R}^n$ , when  $m$  is neither 1 nor  $n - 1$ . For  $m = 1$ , it holds with  $\|d^*u\|_1$  replaced by  $\|d^*u\|_{H^1}$ , and for  $m = n - 1$ , with  $\|du\|_1$  replaced by  $\|du\|_{H^1}$ , where  $H^1$  is the real Hardy space. This looks a little similar since (1.57) implies that  $\|f\|_q \leq C_1 M_{\beta \cdot p; 1}(f) + C_2 \|f\|_{B_{\infty, \infty}^{1/(1-q)}}$  with constants  $C_1, C_2 > 0$ . But we don't know whether it is related to our results, partly because, though it will be the case  $n = 4$ ,  $m = 1$  and  $q = \frac{4}{3}$ , so that if our paper should have a relation, as Proposition 6.1 above says, inequality (6.3) fails to hold for the semi-norms  $M_{\beta \cdot p; 1}^{(4)}(f)$  in (6.4b) and  $M_{\beta \cdot p; 1}^{(5)}(f)$  in (6.5b) in place of  $M_{D_1, D_2; 1}(f)$ .

Finally, as for the third semi-norm  $M_{\beta \cdot p; 1}^{(6)}(f)$  in (6.6b) associated with the decomposition (6.6a), it is not clear whether or not (6.3) holds, although we learn in Theorem 1.6 that it holds for its *next* stronger semi-norm  $M_{\beta \cdot p; 1}(f)$ , but in Proposition 6.1 above that it does not for its *next weaker* semi-norm  $M_{\beta \cdot p; 1}^{(5)}(f)$ . However, it should be probably noted here that the sequence  $\{f_n\}$  used to construct the counterexample in Section 3 *does not violate but keeps* inequality (1.57) with semi-norm  $M_{\beta \cdot p; 1}^{(6)}(f)$  in place of  $M_{\beta \cdot p; 1}(f)$ . Needless to say, this sequence  $\{f_n\}$  of course keeps inequality (1.57) safe, though.

## 7 Summary

In this work we have extended the improved Sobolev embedding theorem (1.1), which originally is for single-valued functions, to a vector-valued version, (1.28) and (1.30), which are connected with the three-dimensional massless Dirac operator  $\alpha \cdot p$  in (1.4)/(1.17):

$$1 \leq p < q < \infty : \quad \|f\|_q \leq C M_{\alpha \cdot p; p}(f)^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.28)$$

$$1 < p < q < \infty : \quad \|f\|_q \leq C \|(\alpha \cdot p)f\|_p^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.30)$$

where  $f(x) = {}^t(f_1(x), f_2(x), f_3(x), f_4(x))$  are  $\mathbb{C}^4$ -valued functions on  $\mathbb{R}^3$ . The first-order-derivative semi-norm  $M_{\alpha \cdot p; p}(f)$  on the right of (1.28) is at first defined by (1.19) with the rather artificial decomposition (1.18) of  $\alpha \cdot p$  into the sum of its two parts, but then can be seen, through its explicit expression (1.21), to coincide with the ones to be defined with the other decompositions like (1.24), (1.25) and (1.26), just as clarified in Proposition 1.0. This will reveal the semi-norm  $M_{\alpha \cdot p; p}(f)$  to have an intrinsic meaning.

When  $1 < p < q < \infty$ , the semi-norm  $M_{\alpha \cdot p; p}(f)$  is equivalent to the semi-norm  $\|(\alpha \cdot p)f\|_p$  as well as  $\|\nabla f\|_p$ . Therefore, in this case it is no wonder that inequality (1.30) holds, because (1.28) is reduced to (1.30) which is also equivalent to (1.13). It also is an improvement of the (1.14) that has the weak  $L^q$  norm on the left-hand side.

But when  $p = 1$ , these three first-order-derivative semi-norms are not equivalent to one another, cf. (1.22). In this case, (1.16)/(1.30) does not hold in general. A counterexample is given in Section 3. Further, for  $p = 1$  two inequalities (1.28) and (1.14) cannot be compared so as to say which of them is sharper.

Analogous improved Sobolev embedding theorems are also given for the three-dimensional Weyl–Dirac (or Pauli) operator  $\sigma \cdot p$  in (1.31), the Cauchy–Riemann operator  $\frac{1}{2}(\partial_1 + i\partial_2)$  and the four-dimensional Euclidian Dirac operator  $\beta \cdot p$  in (1.46). Here, for the last one  $\beta \cdot p$ , in the same way as for  $\alpha \cdot p$ , the semi-norm  $M_{\beta \cdot p; p}(f)$ , which is defined at first by (1.48) with the rather artificial decomposition (1.47), turns out to coincide with the ones to be defined with the other decompositions like (1.53), (1.54) and (1.55), and so to be meaningful. Noted is in Section 2, 5<sup>o</sup> that all the results are also valid for the other representations of the three-dimensional massless and the four-dimensional Euclidian Dirac operators.

However, exceptionally for the *two*-dimensional Weyl–Dirac (or Pauli) operator  $(\sigma \cdot p)^{(2)}$  in (1.40abc), we have proved an inequality which is just expected as (1.16) for all  $1 \leq p < q < \infty$ :

$$\|f\|_q \leq C \|(\sigma \cdot p)^{(2)} f\|_p^{p/q} \|f\|_{B_{\infty, \infty}^{p/(p-q)}}^{1-(p/q)}, \quad (1.41)$$

for  $\mathbb{C}^2$ -valued functions  $f(x) = {}^t(f_1(x), f_2(x))$  on  $\mathbb{R}^2$ , which might be said to be a *true* extension of the single-valued (1.1) to the vector-valued version.

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