# Simulation of A Soap Film Catenoid 

Pornchanit Subvilaia ${ }^{a, b}$<br>${ }^{a}$ Graduate School of Natural Science and Technology, Kanazawa University, Kakuma, Kanazawa 920-1192 Japan<br>${ }^{b}$ Department of Mathematics, Faculty of Science, Chulalongkorn University, Pathumwan, Bangkok 10330 Thailand<br>E-mail: hippoarale@gmail.com


#### Abstract

There are many interesting phenomena concerning soap film. One of them is the soap film catenoid. The catenoid is the equilibrium shape of the soap film that is stretched between two circular rings. When the two rings move farther apart, the radius of the neck of the soap film will decrease until it reaches zero and the soap film is split. In our simulation, we show the evolution of the soap film when the rings move apart before the film splits. We use the BMO algorithm for the evolution of a surface accelerated by the mean curvature.


Keywords: soap film catenoid, minimal surface, hyperbolic mean curvature flow, BMO algorithm

## 1. Introduction

The phenomena that concern soap bubble and soap films are very interesting. For example when soap bubbles are blown with any shape of bubble blowers, the soap bubbles will be round to be a minimal surface that is the minimized surface area. One of them that we are interested in is a soap film catenoid. The catenoid is the minimal surface and the equilibrium shape of the soap film stretched between two circular rings.

In the observation of the behaviour of the soap bubble catenoid [3], if two rings move farther apart, the radius of the neck of the soap film will decrease until it reaches zero. Then the soap film is split and a small bubble will appear.

In this simulation, the soap film catenoid was simulated when the rings move apart before the film split by adapting BMO algorithm for the evolution of a surface accelerated by the mean curvature.

## 2. Derivation of the equation

When the force act to the interface in the normal direction, the equation of interfacial dynamics [2] is given by

$$
\begin{equation*}
A=-\kappa n, \tag{1}
\end{equation*}
$$

where $A$ is the interfacial normal acceleration, $\kappa$ is the mean curvature, and $n$ is the unit normal vector to the interface.

The differential equation for the evolution (1) is

$$
\begin{cases}\alpha_{t t}(t, x)=-\kappa(x) & \text { in }(0, T) \times \Omega \\ \alpha_{t}(t=0, x)=v_{0}(x) & \text { in } \Omega \\ \alpha(t=0, x)=\gamma(x) & \text { in } \Omega,\end{cases}
$$

where $\alpha$ is a position function, $v_{0}$ is an initial velocity and $\gamma$ is a parametrized curve.

## 3. BMO algorithm for the evolution of a surface accelerated by the mean curvature

The original BMO was proposed by Bence, Merriman, and Osher in [1]. The BMO is useful for computing motion of a soap film by mean curvature. In this paper, we use the idea of BMO for the evolution of a surface accelerated by the mean curvature [2] because the the soap film catenoid has the force that act on the interface to reduce its area.
At a given time $T, h=T / M$, where $0<h \ll 1$ and $M$ is a positive number. For $k=0,1, \ldots, M$
Let $\Gamma_{k}$ be a smooth curve and $E_{k}$ be region enclosed by $\Gamma_{k}$ and boundary at time $k$.
We use two signed distance functions for initial condition.

$$
\begin{gathered}
d_{0}(x)= \begin{cases}\inf _{y \in \Gamma_{0}}\|x-y\| & \text { if } x \in E_{0} \\
-\inf \\
y \in \Gamma_{0} \\
\inf ^{2}-y \| & \text { otherwise, },\end{cases} \\
d_{-1}(x)=\left\{\begin{array}{ll}
y \in \inf _{0}-v_{0} h & \|x-y\| \\
\text {-iff } x \in E_{-1} \\
y \in \Gamma_{0}-v_{0} h
\end{array}\|x-y\|\right. \\
\text { otherwise. }
\end{gathered}
$$

1) Set $u_{0}(x)=2 d_{k}-d_{k-1}$, where

$$
d_{k}(x)=\left\{\begin{array}{ll}
\inf _{y \in \Gamma_{k}} & \|x-y\| \\
\inf _{y \in \Gamma_{k}} & \|x-y\|
\end{array} \quad \text { if } x \in E_{k} . \quad \text { otherwise } .\right.
$$

2) Solve the wave equation with zero initial velocity and initial condition $u_{0}$ at a time $h$

$$
\begin{cases}u_{t t}=\triangle u & \text { in }(0, h) \times \Omega \\ \frac{\partial u}{\partial v}=0 & \text { on }(0, h) \times \partial \Omega \\ u(t=0, x)=u_{0} & \text { in } \Omega \\ u_{t}(t=0, x)=0 & \text { in } \Omega .\end{cases}
$$

3) Update the interface and the region.
4) Calculate $d_{k+1}$.

## 4. Model of a soap bubble catenoid

### 4.1 Derivation of the equation concerning a soap bubble catenoid

For the soap bubble catenoid, there is the force that act on the interface to reduce its area because the catenoid is the minimal surface. The force act on the interface that the curvature is $1 / r$, where $r$ is the radius of the circle in the catenoid for each grid point. Adding $1 / r$ in the equation (1), we get

$$
\begin{equation*}
A=(-\kappa+1 / r) n \tag{2}
\end{equation*}
$$



Figure 1. The half of catenoid
The differential equation concerning the soap bubble catenoid (2) is

$$
\begin{cases}\alpha^{\prime \prime}(t, x)=-\kappa(x)+1 / r(x) & \text { in }(0, T) \times \Omega \\ \alpha^{\prime}(t=0, x)=0 & \text { in } \Omega \\ \alpha(t=0, x)=\gamma(x) & \text { in } \Omega,\end{cases}
$$

where $\gamma(x)$ is a closed curve that be parametrised and initial velocity is zero.
4.2 BMO algorithm for a soap film catenoid

At a given time $T, h=T / M$, where $0<h \ll 1$ and $M$ is a positive number. For $k=0,1, \ldots, M$

1) Set $u_{0}(x)=d_{k}$.
2) Compute the radius in the catenoid for each grid point.
3) Solve the equation with zero initial velocity and initial condition $u_{0}$ at a time $h$

$$
\begin{cases}u_{t t}=\Delta u+1 / r & \text { in }(0, h) \times \Omega  \tag{5}\\ u=u_{0} & \text { on }(0, h) \times \partial \Omega \\ u(t=0, x)=u_{0} & \text { in } \Omega \\ u_{t}(t=0, x)=0 & \text { in } \Omega\end{cases}
$$

4) Update the interface and the region.
5) Calculate $d_{k+1}$.

## 5. The equation solving

We solve the equation (5) by using finite difference method(central difference method for the second derivative). Let $u_{i, j}^{k}$ be an approximate solution at $\left(t_{k}, x_{i}, y_{j}\right), \Delta t=t_{k}-t_{k-1}, \quad \Delta x=x_{i}-x_{i-1}$, $\Delta y=y_{j}-y_{j-1}$ and $\Delta x=\Delta y$. Then the approximate solution is

$$
u_{i, j}^{k+1}=2 u_{i, j}^{k}-u_{i, j}^{k-1}+\left(\frac{\Delta t}{\Delta x}\right)^{2}\left(u_{i+1, j}^{k}+u_{i-1, j}^{k}+u_{i, j+1}^{k}+u_{i, j-1}^{k}-4 u_{i, j}^{k}\right)+\Delta t^{2} / r_{i, j}^{k}
$$

## 6. Numerical results

This simulation that we show is only some part of soap film catenoid.
$400 \times 400$ grids, time discretisation is $0.001 / 10$.


## Acknowledgment

This work was supposed by Japan student services organization (JASSO). I would like to express my gratitude to Prof. Seiro Omata and Prof. Norbert Pozar for their many suggestions. I also acknowledge my friends for all their support.

## References

1. B. Merriman, J. Bence, S. Osher (1994), Motion of Multiple Junctions: A Level Set Approach, J. Comp. Phys. 112, 334-363.
2. E. Ginder and K. Svadlenka (2014), On an algorithm for curvature-dependent interfacial acceleration, Proceedings of Computational Engineering Conference, vol. 19
3. M. Ito and T. Sato (2010), In situ observation of a soap-film catenoid-a simple educational physics experiment, Eur. J. Phys., 31, 357-365.
