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Revised paper

Title page

X-ray induced defects as obstacles to dislocation motion in alkali halide crystals

Y. Kohzuki^{1,*} and T. Ohgaku²

¹*Department of Mechanical Engineering, Saitama Institute of Technology, Saitama, Japan E-mail: kohzuki@sit.ac.jp*

²*Graduate School of Natural Science and Technology, Kanazawa University, Kanazawa, Japan E-mail: ogk@po3.nsknet.or.jp*

Address correspondence to E-mail: kohzuki@sit.ac.jp (telephone number: +81-48-585-6907)

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Y. Kohzuki^{1,*} and T. Ohgaku²

¹*Department of Mechanical Engineering, Saitama Institute of Technology, Saitama, Japan*

²*Graduate School of Natural Science and Technology, Kanazawa University, Kanazawa, Japan*

ABSTRACT

NaBr and KBr single crystals were exposed to X-ray (W-target, 30 kV, 20 mA) at room temperature and strain-rate cycling tests associated with ultrasonic oscillation were carried out for the crystals at 77 to 293 K. The stress decrement ($\Delta\tau$) due to oscillation and the stress change due to strain-rate cycling have been measured during plastic deformation. The relative curve of $\Delta\tau$ and strain-rate sensitivity (λ) of flow stress has a stair-like shape for the two kinds of crystals. That is to say, the curve has two bending points and is divided into three regions: two plateau regions and the region between the two bending points, where λ decreases gradually with increasing $\Delta\tau$. The first region is a plateau within the small $\Delta\tau$. This implies that X-ray induced defects have weak interaction with dislocation and act as obstacles to dislocation motion. Furthermore, the dependence of stress decrement (τ_p) at the first bending point on the activation volume (V) obtained from the difference between λ in the first and second plateau regions reflects the interaction between dislocation and defects induced by the X-irradiation. The activation energy for the break-away of a dislocation from the defect could be obtained on the basis of τ_p - V curve fitting the Barnett model to the experimental results. Then, the values of activation energy are 0.76 and 0.81 eV for NaBr and for KBr, respectively.

Introduction

Stress-strain curve is altered drastically by electron irradiation. The effect generally increases with increasing electron dosage, decreasing irradiation temperature and decreasing test temperature. The yield stress and the strain-rate sensitivity of flow stress increase with the irradiation dose [1]. It is well known that alkali halide crystals harden when they are exposed to X-ray irradiation [2–5]. This is because lattice defects induced by X-ray irradiation obstruct dislocation motion. The defect is electron-center

or hole-center such as F-center, V_{2-} or V_{3-} -center. The interaction between dislocation and the defect is not clearly established yet. Interaction between dislocation and point defects has been widely investigated by the yield stress measurements (e.g., ref. [6, 7]), the direct observations of dislocation (e.g., ref. [8, 9]), the internal friction measurements (e.g., ref. [10, 11]), or the stress relaxation tests (e.g., ref. [12, 13]) so far. However, these widely well-known methods have some weak points as follows. Yield stress depends on dislocation velocity, dislocation density, and multiplication of dislocations [14]. With regard to the direct observations, electron microscopy is useful for a thin specimen but not for bulk, and also light scattering method provides the information on the a transparent specimen. Internal friction measurements concern the motion of the dislocation which breaks away from the weak obstacles such as impurities between two forest dislocations with vibration [15]. This is not for dislocation motion during plastic deformation. As for stress relaxation tests, it is assumed that internal structure of crystals does not change, i.e. dislocation density and internal stress are constant. These are the reasons why it is difficult to study dislocation-point defects interaction in bulk during plastic deformation. But the original method (the strain-rate cycling tests during the Blaha effect) is considered to break through the weak points in the widely known methods and has the merit different from these methods. The Blaha effect is the phenomenon that the static flow stress is decreased when ultrasonic oscillations are superimposed during plastic deformation. It has been reported so far that the measurement of strain-rate sensitivity under application of ultrasonic oscillatory stress during plastic deformation gives useful information on the interaction between a mobile dislocation and point defects in alkali halide crystals [16–18]. We report the interaction between dislocation and defects induced by X-ray irradiation in NaBr and

KBr single crystals through the use of data obtained by the strain-rate cycling tests associated with ultrasonic oscillation.

Experimental procedure

The specimens, which were prepared by cleaving out of single crystalline ingots to the size of $5 \times 5 \times 15 \text{ mm}^3$, were kept at 973 K for 20 h and were cooled to room temperature at the rate of 40 K/h. The specimens were exposed to X-ray for 3 h on each of the pair wide surfaces at room temperature. The total exposure time is 6 h. The irradiation was carried out with X-rays from W-target operated at 30 kV and 20 mA by a Shimadzu XD-610. Fig. 1 shows the absorption spectrum of the X-ray irradiated NaBr and KBr, which was measured with a Hitachi U-3500 spectrophotometer. Seitz [19, 20] has reported the color center of alkali halide. The peak of F-band (2.3 eV) is seen in the spectra of X-ray irradiated NaBr. The peak above 6 eV is attributed to OH [21]. In the spectra of the X-ray irradiated KBr, the peak of F-(2.0 eV) band is seen and the peak of 4.7 eV is attributed to V_2 -band. The average concentration of F-centers, which was estimated from the Smakula formula [22], is $15 \times 10^{16} \text{ cm}^{-3}$ for NaBr and $13 \times 10^{16} \text{ cm}^{-3}$ for KBr.

Figure 2 A schematic illustration of the apparatus is shown in Fig. 2. A resonator composed of a vibrator and a horn with the resonant frequency of 20 kHz was attached to a Shimadzu DSS-500 testing machine. The specimens were lightly fixed on a piezoelectric transducer and then cooled down to a test temperature. The specimens were deformed by compression along the $\langle 100 \rangle$ axis and the ultrasonic oscillatory stress was intermittently superimposed for one or two minutes by the resonator in the

same direction as the compression. The amplitude of the oscillatory stress was monitored by the output voltage from the piezoelectric transducer. The strain of the specimen is considered to be homogeneous, because the wavelength, which is 167 mm for NaBr and 152 mm for KBr by calculating from the data of ref. [23], is about 10 ~ 11 times as long as the length of the specimen.

Figure 3

The strain-rate cycling test during the Blaha effect measurement is illustrated in Fig. 3. Superposition of oscillatory stress (τ_v) during plastic deformation causes a stress drop ($\Delta\tau$). Keeping the stress amplitude of τ_v constant, strain-rate cycling between the strain rates of $\dot{\epsilon}_1$ ($1.1 \times 10^{-5} \text{ s}^{-1}$) and $\dot{\epsilon}_2$ ($5.6 \times 10^{-5} \text{ s}^{-1}$) was carried out at 77 K to room temperature. Then, the stress change due to the strain-rate cycling is $\Delta\tau'$. For the tests at 77 K, the specimen was immersed in the liquid nitrogen. The other temperature measurements were made by thermocouples of Ni-55%Cu vs. Cu. Each specimen was held at the test temperature for 30 min prior to the compression test and the stability of temperature during the test was kept within 2 K. The strain-rate sensitivity ($\Delta\tau'/\Delta\ln\dot{\epsilon}$) of the flow stress, which is given by $\Delta\tau'/1.609$, was used as a measure of the strain-rate sensitivity ($\lambda = \Delta\tau'/\Delta\ln\dot{\epsilon}$).

Results and discussion

Relation between $\Delta\tau$ and λ

Figure 4a,b

Fig. 4a for X-ray irradiated NaBr at 173 K shows that $\Delta\tau$ does not change significantly with shear strain but increases with stress amplitude at a given temperature and shear strain. Fig. 4b shows that λ decreases with stress amplitude and the variation of it with τ_v is small at low and high amplitudes. The λ increases with shear strain due to the

increase of the forest dislocation density. The values of λ and $\Delta\tau$ obtained from

Figure 5

Fig. 4a and b at shear strains of 15 and 20 % are plotted in Fig. 5 as the relation between

λ and $\Delta\tau$ for the given internal structure. The variation of λ with $\Delta\tau$ is stair-like.

That is to say, there are two bending points and two plateau regions: the first plateau

region ranges below the first bending point at low stress decrement and the second one

extends from the second bending point at high stress decrement. The λ decreases

gradually with increasing $\Delta\tau$ between the two bending points. The value of $\Delta\tau$ at the

first bending point is referred to as τ_p in Fig. 5. As given in the below description the

relative curve between $\Delta\tau$ and λ reflects the effect of ultrasonic oscillation on the

dislocation motion on the slip plane containing many weak obstacles and a few forest

dislocations during plastic deformation, which was detailed in the paper [18]. Within the

first plateau region of $\Delta\tau$ versus λ curve in low $\Delta\tau$, the application of oscillation with

low stress amplitude cannot influence the average length of dislocation segments (\bar{l})

and \bar{l} is considered to remain constant. All weak obstacles act as impedimenta to the

dislocation motion in the region. Applying oscillation with higher stress amplitude, the

dislocation begins to break-away from the weak ones between the forest dislocations

and \bar{l} begins to increase. Then, the λ of flow stress starts to decrease at the first

bending point τ_p , since λ is inversely proportional to \bar{l} [24]. Some weak obstacles stop

Figure 6

acting as impedimenta there. λ_p denoted in the figure is introduced later. Fig. 6 shows

the influence of temperature on the relationship between λ and $\Delta\tau$ for the KBr crystals.

τ_p tends to be lower as the temperature is larger and disappears at room temperature.

Figure 7

Fig. 7 corresponds to the case for non-irradiated KBr single crystals at 193 K. This

figure represents the influence of shear strain on the relation between λ and $\Delta\tau$. The

relative curve shifts upward with increasing shear strain. This phenomenon is caused by the part of the λ which depends on dislocation cuttings. Furthermore, there is only one bending point on each curve in contrast to the irradiated KBr crystals, which is considered to correspond to second bending point as shown in Figs 5 and 6. The relative curve for the irradiated KBr crystal has first bending point at near 193 K as shown in Fig. 6. Similar result is obtained also for the NaBr crystals at low temperature. This implies that the first bending point becomes to appear on the relative curve at low temperatures by irradiating NaBr and KBr. So far, the first bending point, τ_p , at low stress decrement has been considered to represent the effective stress due to the weak obstacles such as the dopant ions which lie on the dislocation when a dislocation begins to break-away from the weak obstacles with the help of oscillation during plastic deformation of alkali halide doped with monovalent or divalent ions [16–18]. The weak obstacles are considered to be the X-ray induced defects here. Therefore, τ_p in this study corresponds to the effective stress due to the defects. The difference between λ at the first plateau place and that at the second one, λ_p defined in Fig. 5, has been regarded as a component of λ due to point defects [25]. And λ_p is related to the inverse of the average spacing, l_p , of point defects along dislocation (i.e. activation volume, V , for the thermally activated dislocation motion) at a given temperature, T , as given in the following equation:

$$\lambda_p = \frac{kT}{bl_p d} \quad (1)$$

where k is the Boltzmann constant, b the magnitude of Burgers vector, and d the

activation distance.

Nadeau [3, 26] reported on several alkali halide crystals subjected to γ -ray irradiation as below: the flow stress increase due to the irradiation correlates with F-center concentration and further the hardening is not directly due to F-centers but due to other factors like interstitials which have the same concentration as F-centers. Sirdeshmukh *et al.* [27] also suggested that the radiation hardening is due to the role of radiation-induced defects other than F-centers. F-centers are so weak to interact with dislocation that they do not act as obstacles to dislocation motion [28]. This would be because F-center has isotropic defect. Hole-center such as V_2 -center, which is observed in the spectra of the X-ray irradiated KBr in Fig. 1, has not isotropic but is considered to have tetragonal distortion. With regard to the spectra of NaBr irradiated with the X-ray, the peak of it is not clear within the measurement. The irradiation-induced lattice defects act as weak obstacles to dislocation motion. This could lead to the appearance of τ_p on the relative curve of λ and $\Delta\tau$ at low temperature for the X-irradiated KBr and NaBr.

Interaction between dislocation and the radiation-induced defects

The force-distance profile between screw dislocation and tetragonal defect was reported by Barnett and Nix [29]. On the assumption that the force-distance profile derived by Barnett and Nix can be applied to the X-irradiated KBr and NaBr, the interaction between dislocation and the radiation-induced defects is studied as below. Fig. 8a and b shows the relation between τ_p and temperature for the crystals. The value of τ_p tends to decrease with increasing temperature and appears to approach to zero at the critical

Figure 8 a,b

Figure 9 a,b

temperature, T_c , for both the crystals. Fig. 9a and b shows the dependence of the

strain-rate sensitivity of flow stress due to point defects on temperature. The solid circles correspond to λ_p for the crystals. Fig. 10a and b shows the relation between τ_p and activation volume V (i.e. kT/λ_p) for the X-ray irradiated NaBr and KBr. Figs 8 to 10 reflect the interaction between dislocation and defects. The solid curves in these figures are obtained as follows. When a dislocation breaks-away from weak obstacles dispersed in the lattice with the aid of thermal fluctuations under the effective stress (τ_e), the activation energy (G) at temperature T is expressed by Eq. (2) [24]:

$$G = \frac{T}{T_c} G_0 \quad (2)$$

where T_c is the critical temperature at which τ_e becomes zero and the dislocation breaks-away from the dopants only with the help of thermal activation, and G_0 the Gibbs free energy for overcoming of the dopant by dislocation at absolute zero. As a function of distance $\xi = x/b$ in units of b from the position of defect on the slip plane b away from it, the force f acted on the dislocation is given by Eq. (3) [29, 30]:

$$f = f_0 \alpha^{3/2} \frac{\alpha \xi^2 + 2\sqrt{2}\xi - 1}{(\alpha \xi^2 + 1)^2} \quad (3)$$

where $\alpha = (C_{11} - C_{12})/2C_{44}$ the anisotropy factor, $f_0 = \Omega \Delta \varepsilon C_{44} / 4\pi b$, Ω the volume of the defect and $\Delta \varepsilon = \varepsilon_{11} - \varepsilon_{22}$ the tetragonality. When the dislocation moves under τ_e at T , the force f_e is acted on the defect. The value of f_e is numerically calculated

using Eqs. (2) and (3) with parameters of appropriate T_c and G_0 so that the area surrounded by the force-distance curve of Eq. (3) (i.e. Barnett-Nix relation) and the straight line of $f = f_e$ satisfies G . τ_e is obtained from the Friedel relation [31]:

$$\tau_e = \tau_0 \left(\frac{f_e}{f_m} \right)^{3/2} \quad (4)$$

where τ_0 is the effective stress at absolute zero and replaced by τ_{p0} (value of τ_p at absolute zero) and f_m the maximum of f . τ_e at different strain-rate (i.e. different T_c in Eq. (2)) is also calculated similarly and $\lambda (= \Delta\tau_e / \Delta \ln \dot{\epsilon})$ is estimated from the difference ($\Delta\tau_e$) between the two τ_e values. V is obtained on the basis of $\Delta\tau_e / \Delta \ln \dot{\epsilon}$.

In this way τ_e and λ at a given temperature and further V vs. τ_e are represented as the solid curves in Figs 8 to 10, where the values of G_0 , τ_0 and T_c are determined as to be fitted to the experimental results (i.e. solid circles in these figures) using the least squares method. These curves agree with the measured data analysed on the basis of $\Delta\tau$ versus λ curves for the crystals. The parameters (τ_{p0} , T_c and G_0) used for

Table 1

calculation are listed in Table 1. The value of G_0 for the irradiated KBr is somewhat larger than for the NaBr. This may suggest that the tetragonality around the X-ray induced defect in the KBr is somewhat larger than in the NaBr.

Conclusion

The strain-rate cycling tests associated with ultrasonic oscillation were carried out for the purpose of investigation on the interaction between dislocation and defects during

plastic deformation of the NaBr and KBr single crystals irradiated with X-ray at room temperature. On the basis of data analyzed in terms of λ versus $\Delta\tau$ due to the oscillation, it is considered that the X-ray induced defects have weak interaction with dislocation and act as obstacles to dislocation motion. Then, the activation energy for the break-away of dislocation from the defect is found to be 0.76eV for NaBr and 0.81eV for KBr. Furthermore, the values of τ_{p0} and T_c are also obtained through the study for the two kinds of crystals.

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Compliance with ethical standards

Conflict of interest The authors declare that they have no conflict of interest.

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Table 1 Parameters used for calculation

Crystal	τ_{p0} (MPa)	T_c (K)	G_0 (eV)
X-Irr. NaBr	1.99	369	0.76
X-Irr. KBr	0.67	660	0.81

Figure captions

Figure 1 Absorption spectrum of X-ray irradiated NaBr and KBr single crystals.

Figure 2 Schematic illustration of the apparatus.

Figure 3 Variation of applied shear stress, τ_a , when the strain-rate cycling between the strain rates, $\dot{\epsilon}_1 (1.1 \times 10^{-5} \text{ s}^{-1})$ and $\dot{\epsilon}_2 (5.6 \times 10^{-5} \text{ s}^{-1})$, is carried out under superposition of ultrasonic oscillatory shear stress, τ_v .

Figure 4 Variation of (a) the stress decrement ($\Delta\tau$) and (b) the λ ($\Delta\tau'/\Delta \ln \dot{\epsilon}$) of flow stress with shear strain at 173 K and various stress amplitude for X-ray irradiated NaBr single crystal. The ultrasonic oscillatory stress (τ_v) increases with its value (arb. units).

Figure 5 Relation between the λ and the stress decrement ($\Delta\tau$) for X-ray irradiated NaBr single crystal at 173 K and the strains of 15 % and 20 %. The plotted points are obtained from Figure 4.

Figure 6 Relation between the λ and the stress decrement ($\Delta\tau$) for X-ray irradiated KBr single crystal at various temperatures and strains.

Figure 7 Relation between the λ and the stress decrement ($\Delta\tau$) for a non-irradiated KBr single crystal at 193 K and the strains of 8-20 %.

Figure 8 Relation between τ_p and temperature for X-ray irradiated (a) NaBr and (b) KBr single crystals. The solid curves are given by numerical calculation.

Figure 9 Dependence of λ_p on temperature for X-ray irradiated (a) NaBr and (b) KBr single crystals. The solid curves are given by numerical calculation.

Figure 10 Relation between τ_p and activation volume for X-ray irradiated (a) NaBr and (b) KBr single crystals. The solid curves are given by numerical calculation.

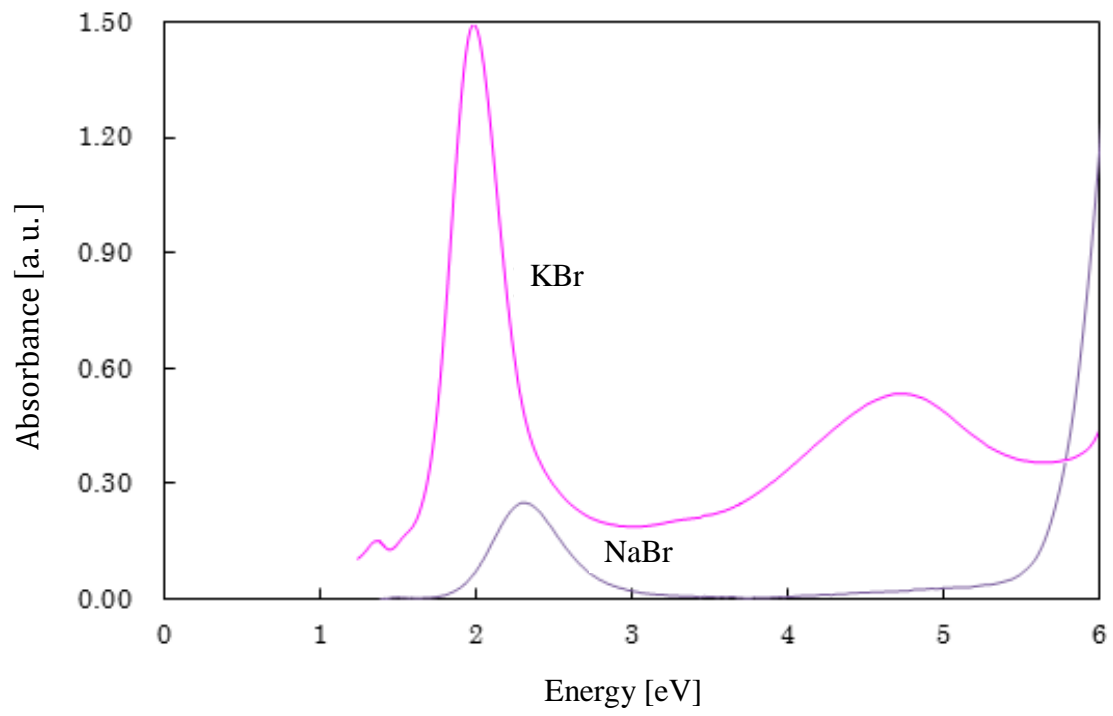


Figure 1

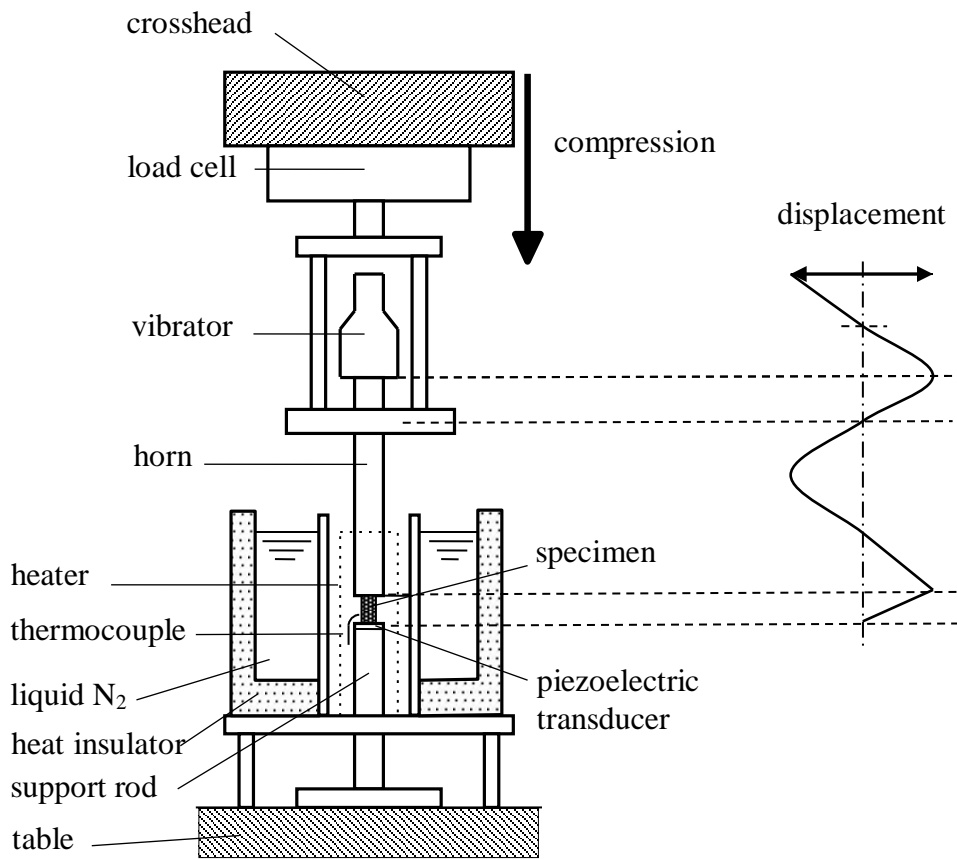


Figure 2

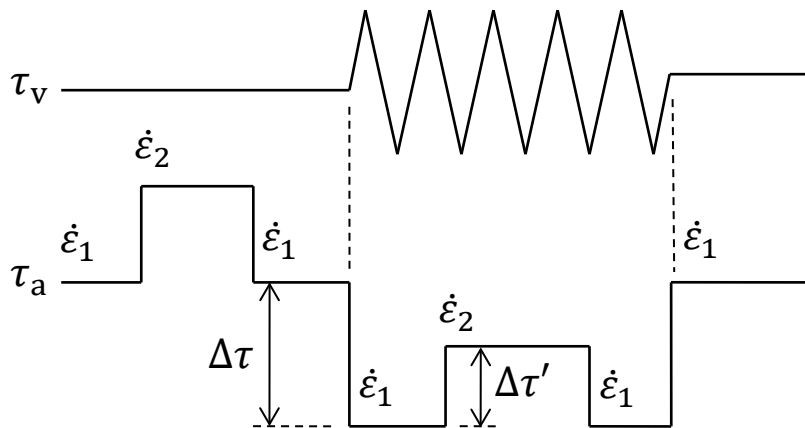


Figure 3

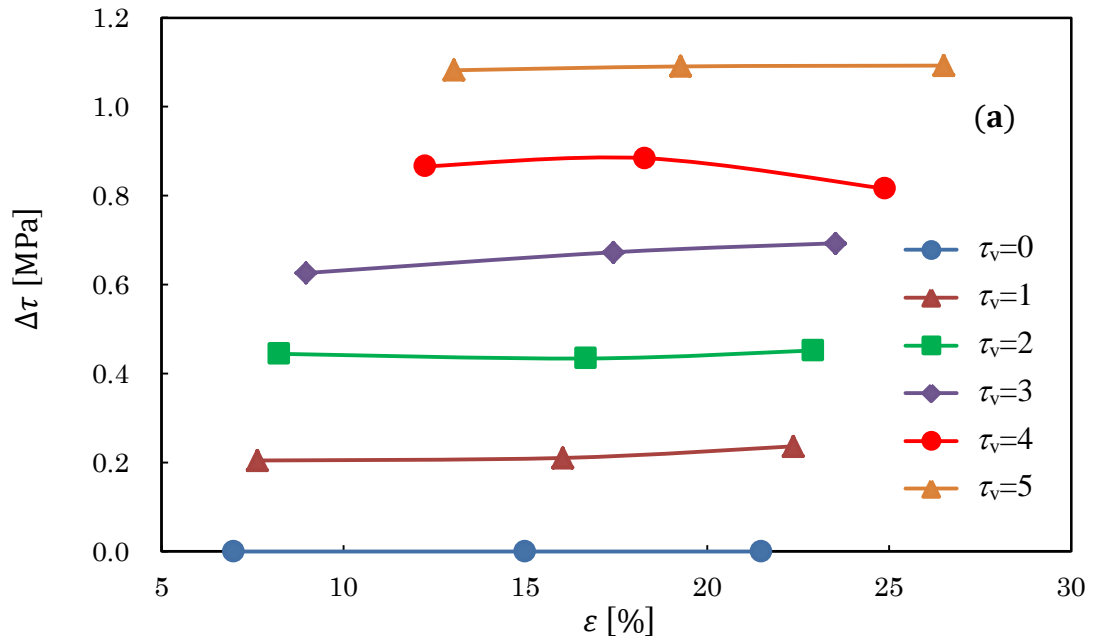


Figure 4 (a)

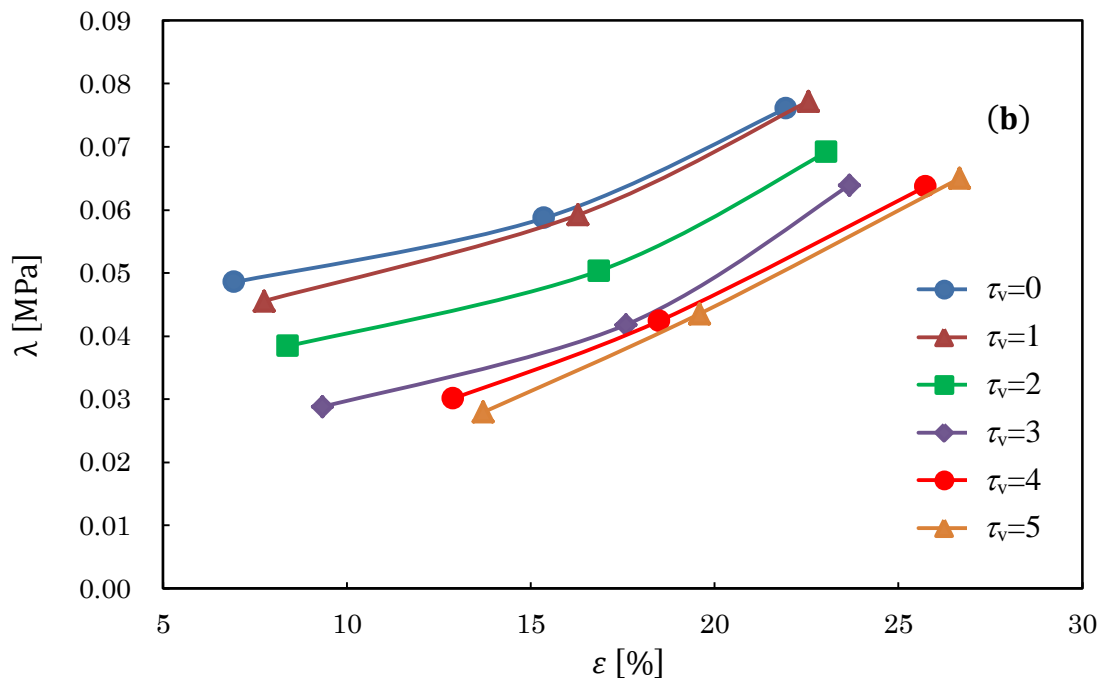


Figure 4 (b)

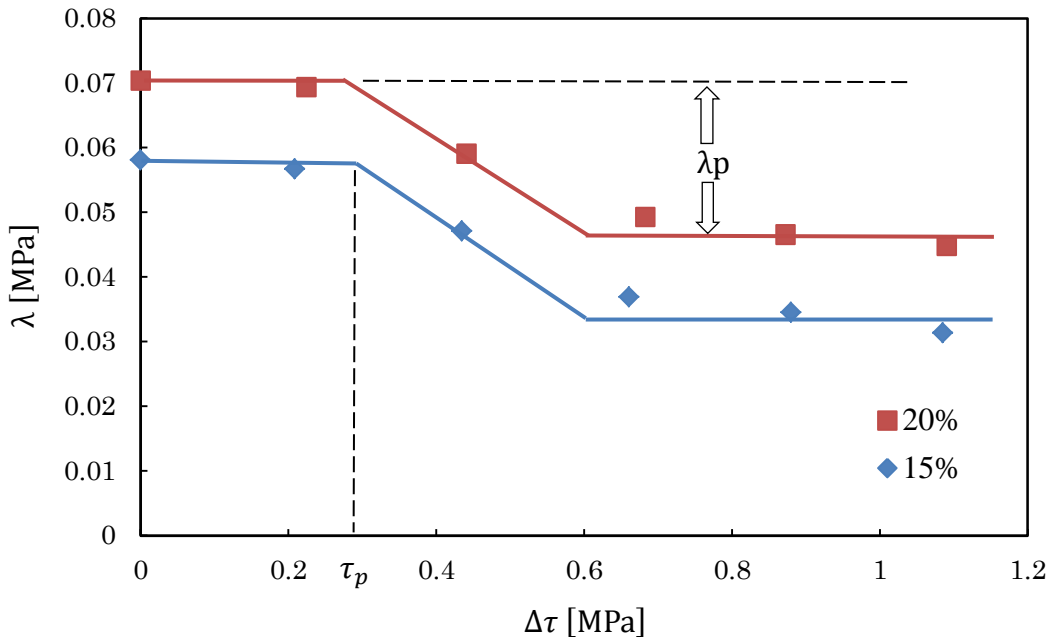


Figure 5

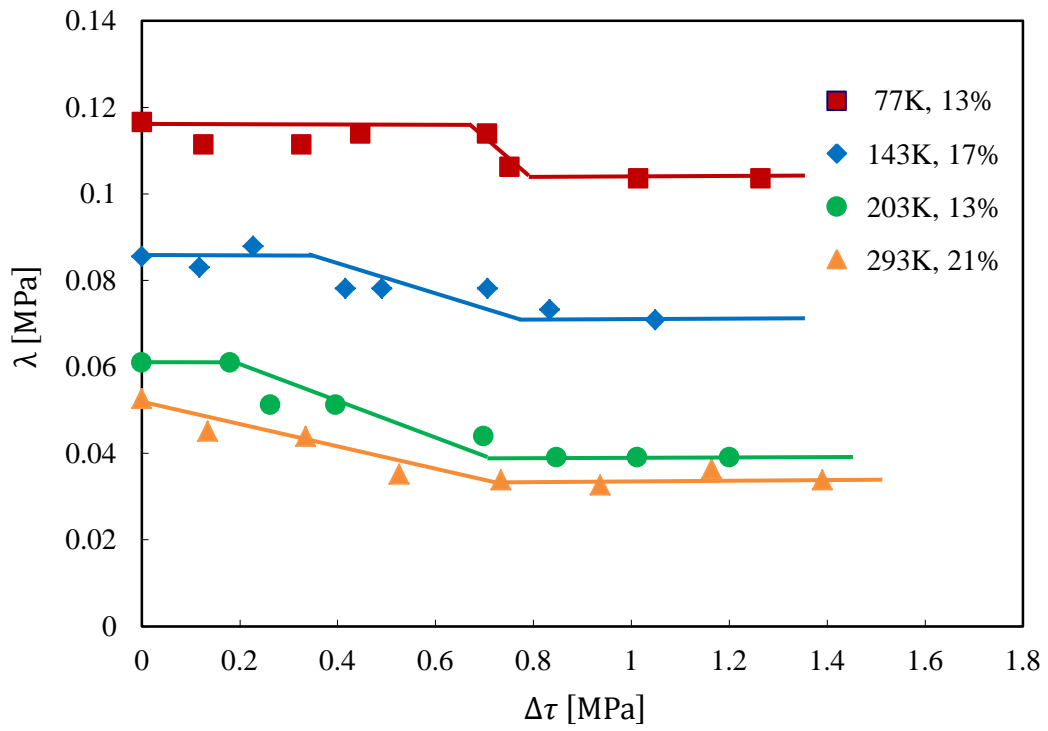


Figure 6

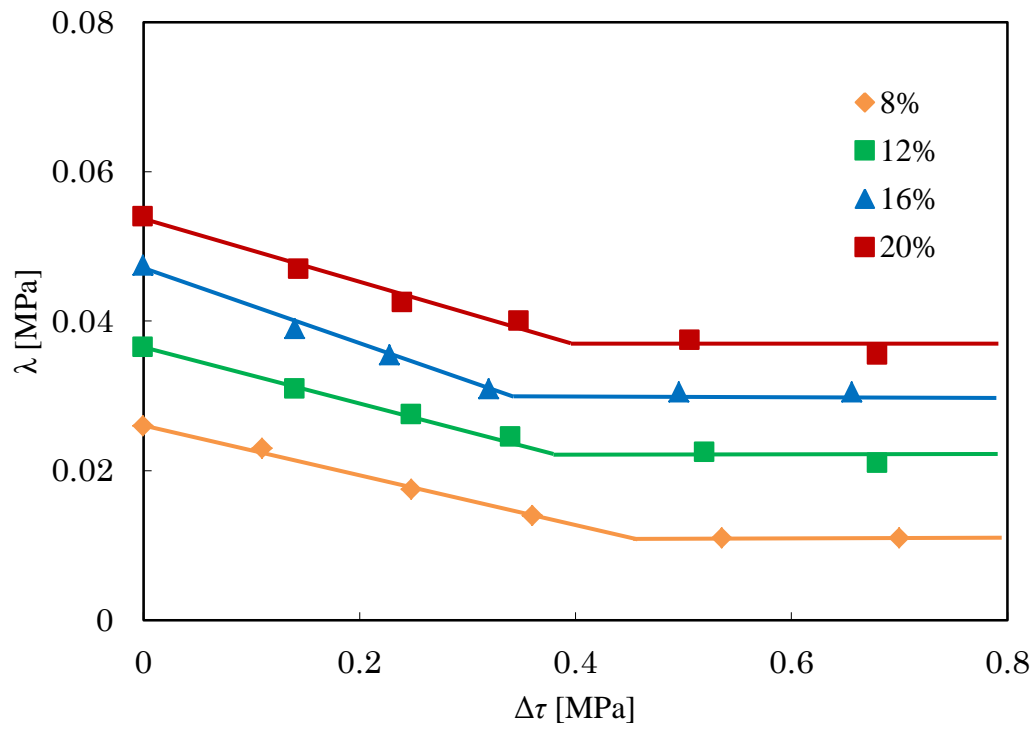


Figure 7

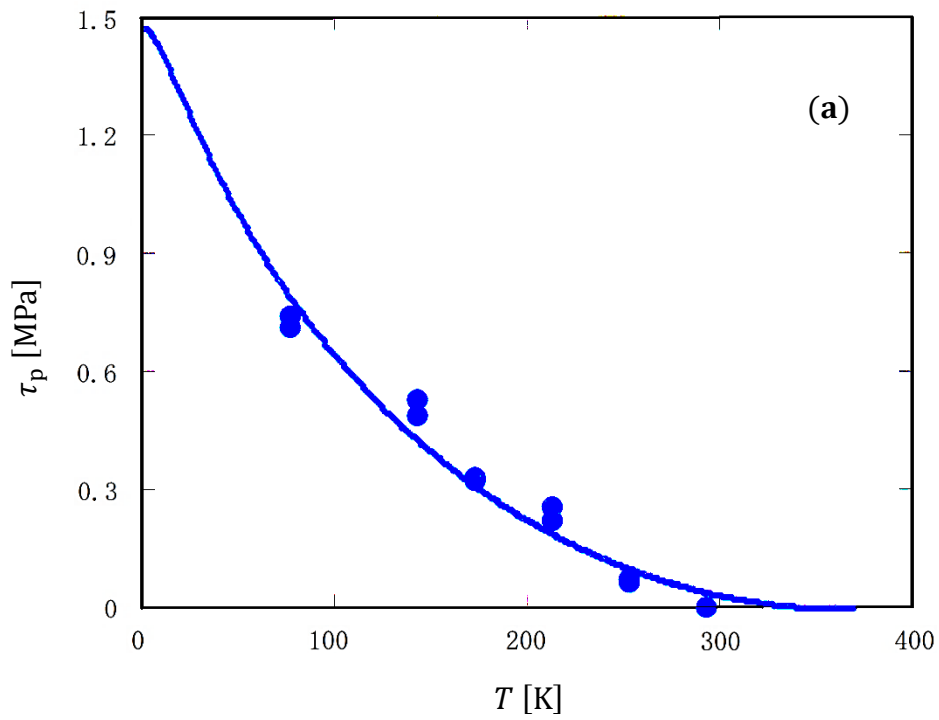


Figure 8 (a)

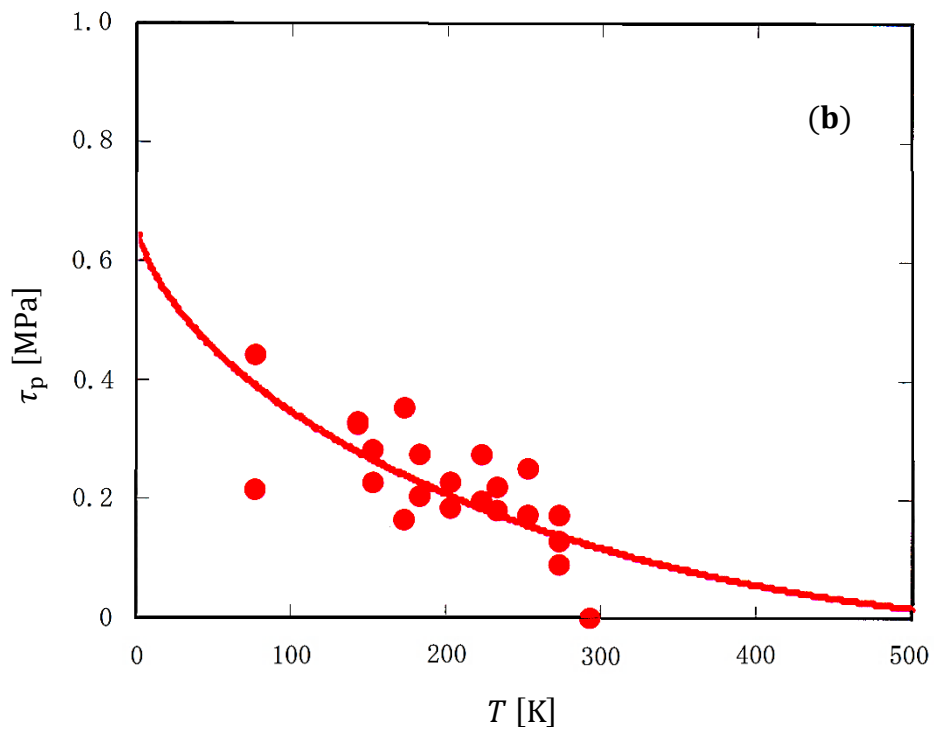


Figure 8 (b)

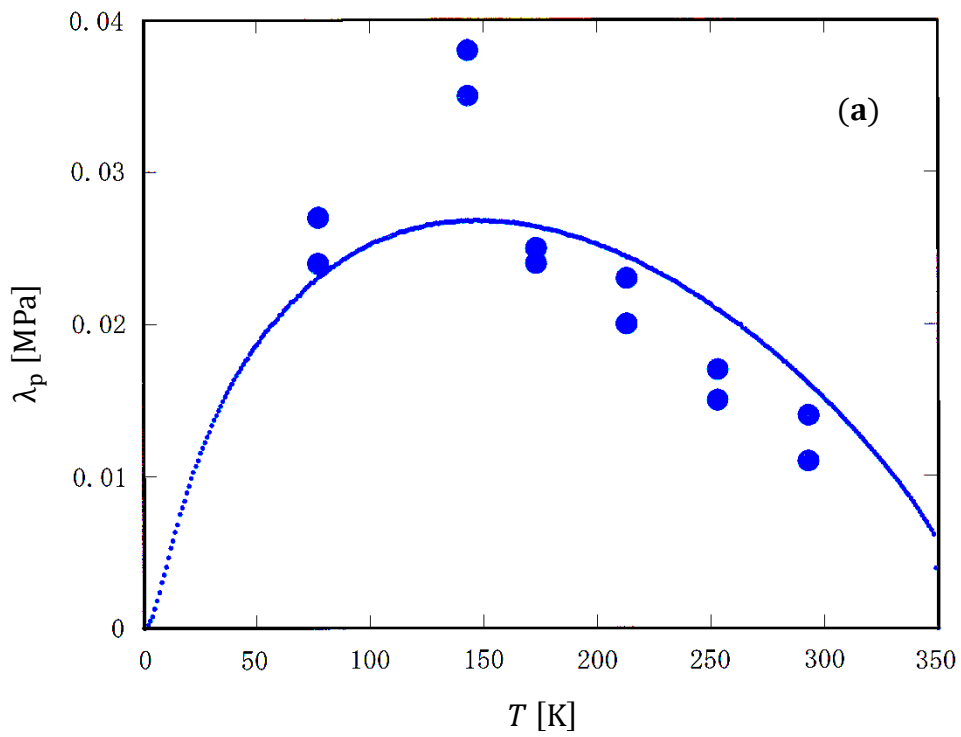


Figure 9 (a)

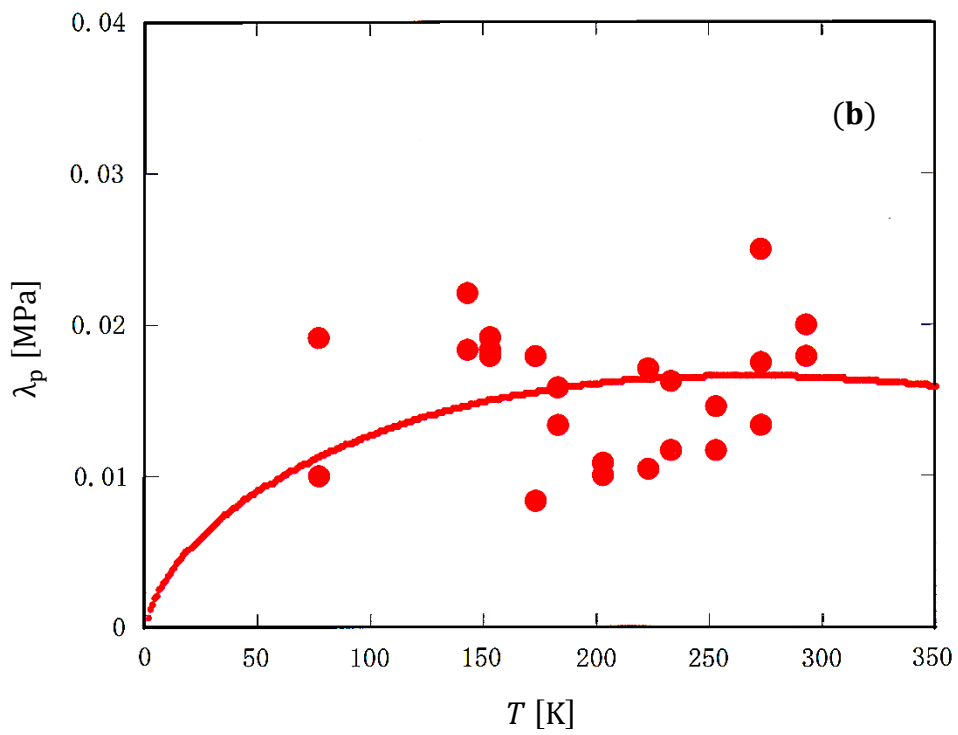


Figure 9 (b)

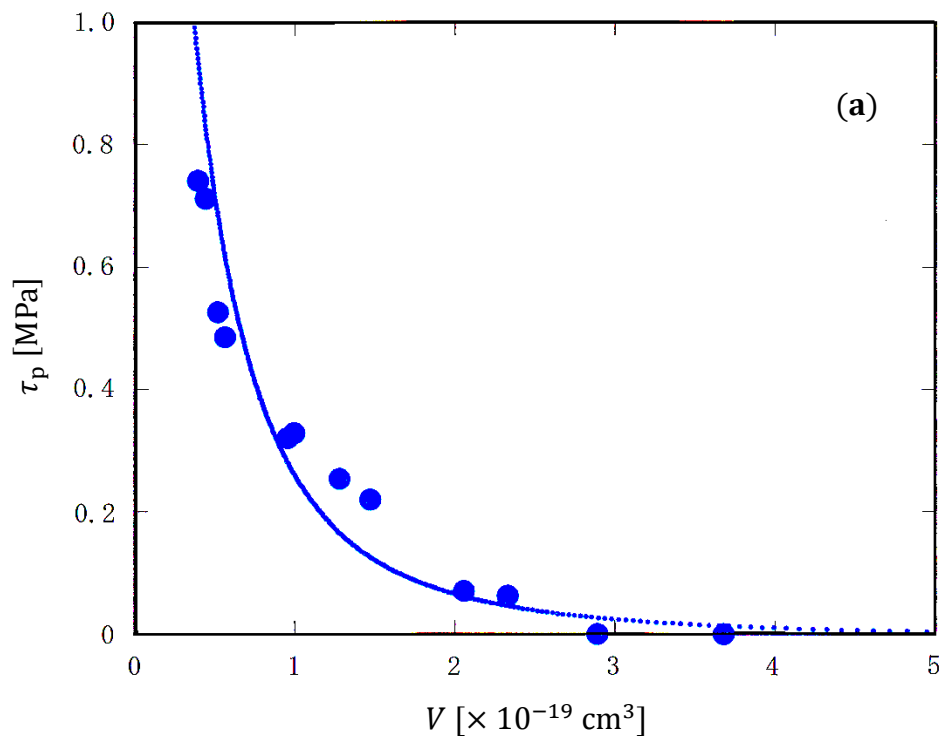


Figure 10 (a)

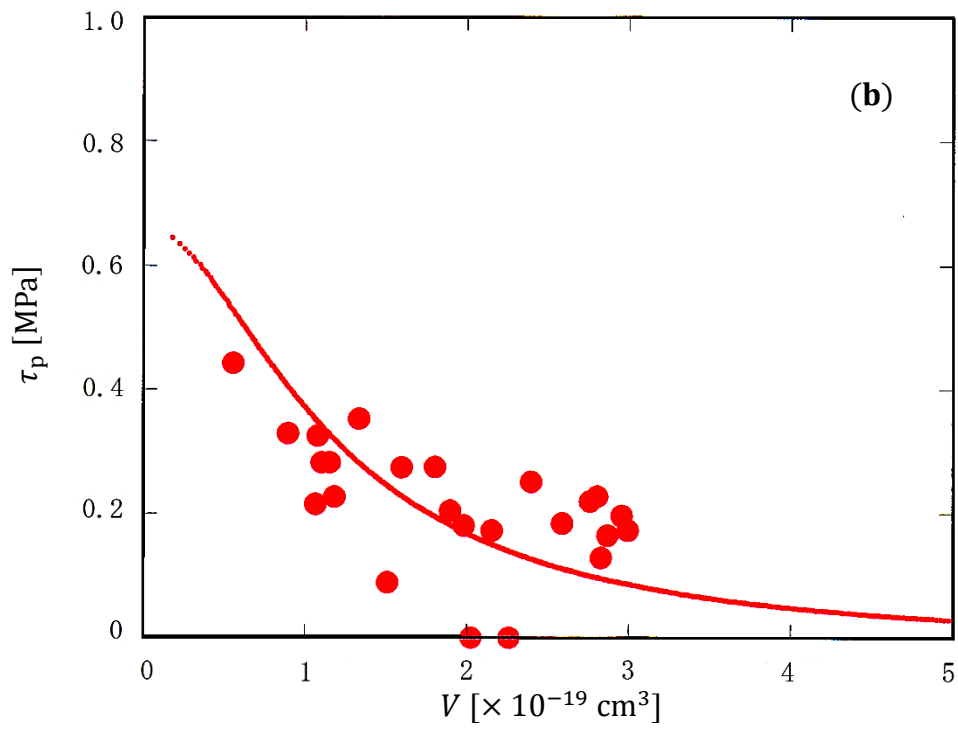


Figure 10 (b)